

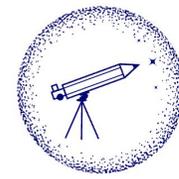
Fitting all spectra and light curves of gamma-ray pulsars with an austere synchro-curvature model

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Co-authors: Diego F Torres, Daniele Viganò

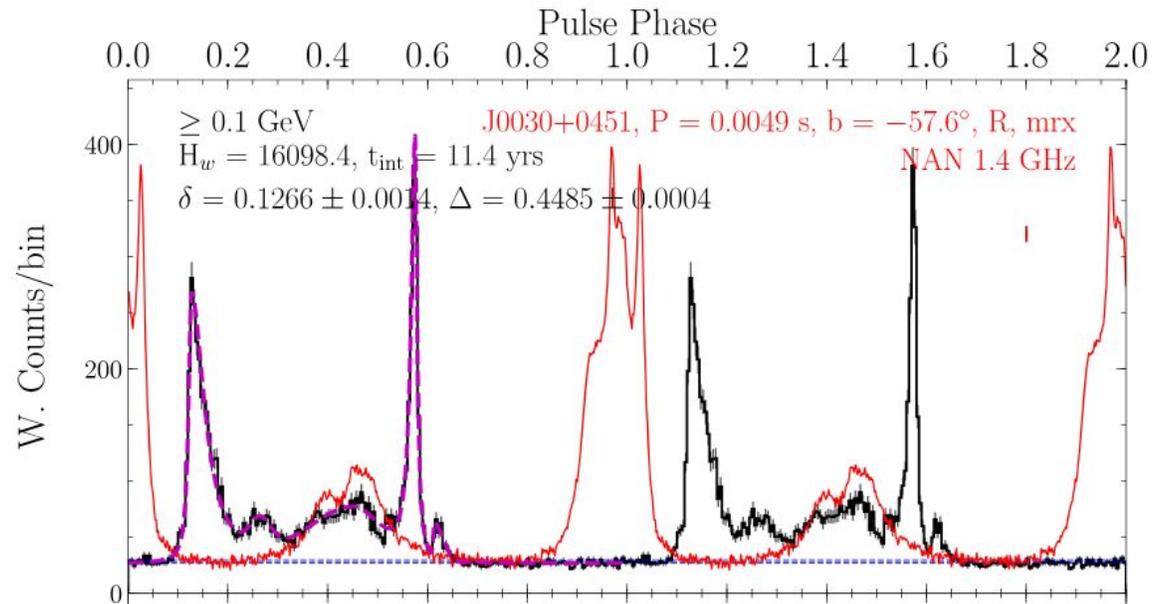
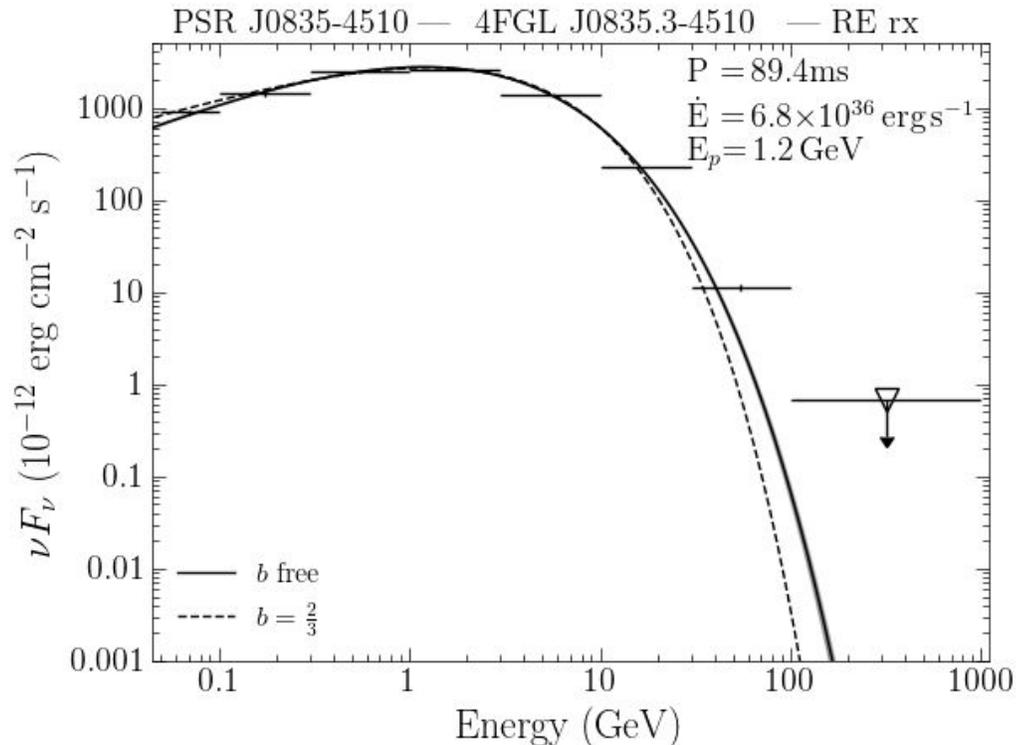
Feeling the pull and pulse of relativistic magnetospheres Workshop

6-11th April 2025 - Les Houches, France



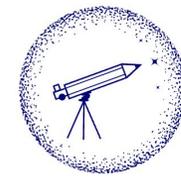
Motivation of the project

- Models of high-energy emission of pulsars can be tested with e.g. the data of the gamma-ray emitting pulsars released in the 3PC (Smith et al. (2023))



- Our goal is to reproduce observational data with a model that contains simple but realistic physics and is computationally affordable

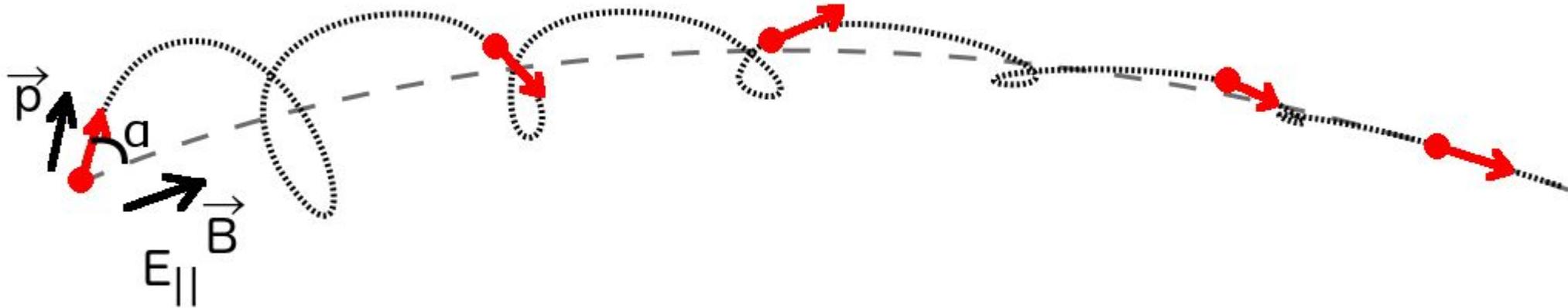
Spectral model: particle dynamics



- We follow the dynamics of the emitting particles, ruled by electric acceleration and synchro-curvature losses and with two free parameters involved: E_{\parallel} , b

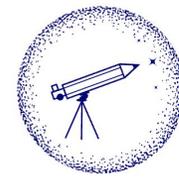
$$\frac{d\mathbf{p}}{dt} = Ze E_{\parallel} \hat{b} - \frac{P_{sc}}{v} \hat{p} \quad B(x) = B_{\star} \left(\frac{R_{\star}}{x} \right)^b \quad r_c(x) = R_{lc} \left(\frac{x}{R_{lc}} \right)^{\eta}$$

- Solving the equation of motion gives the evolution of the relativistic momentum and of the Lorentz factor Γ and pitch angle α



[Viganò et al. 2015, MNRAS, 447, 1164]



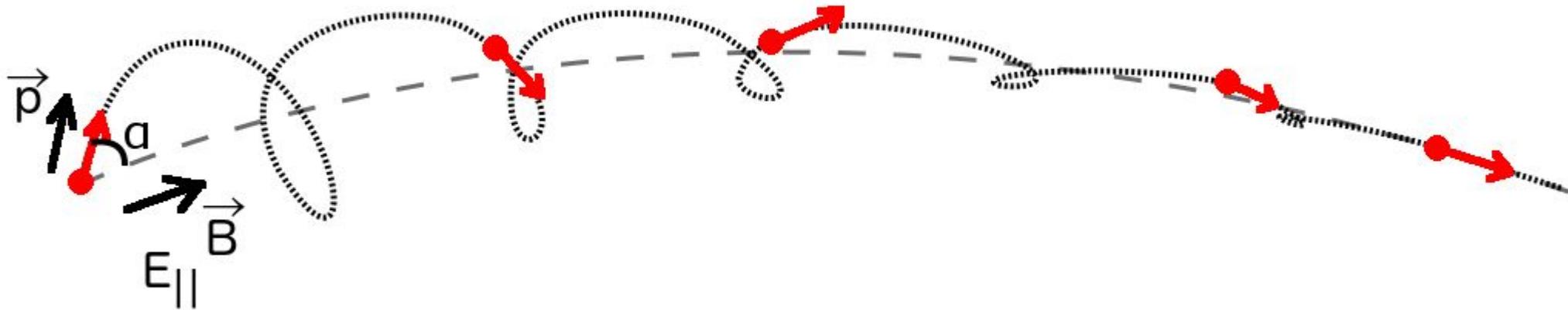


Spectral model: emission

- Synchro-curvature formulae gives the emission of the particles all along the trajectory, which convolved with an effective particle distribution gives the total emission from the region

$$\frac{dP_{sc}}{dE} = \frac{\sqrt{3}e^2\Gamma y}{4\pi\hbar r_{eff}} [(1+z)F(y) - (1-z)K_{2/3}(y)] \quad \frac{dN_e}{dx} = \left[\frac{N_0}{x_0(1 - e^{-(x_{out}-x_{in})/x_0})} \right] e^{-(x-x_{in})/x_0}$$

- We produce theoretical spectra with just three free parameter ($E_{||}$, b , x_0) and a normalization factor $\Gamma \sim 10^3$ $\Gamma \sim 10^7$

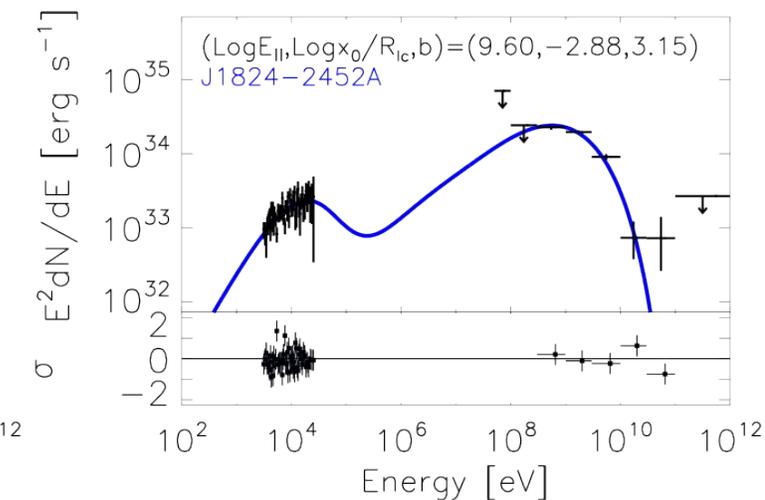
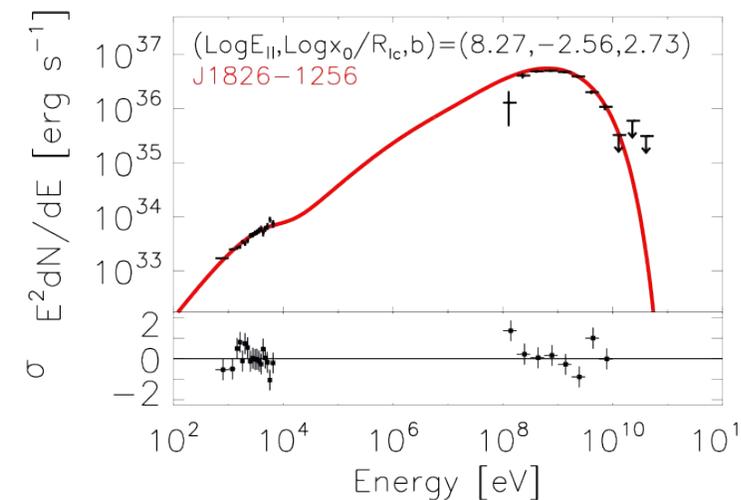
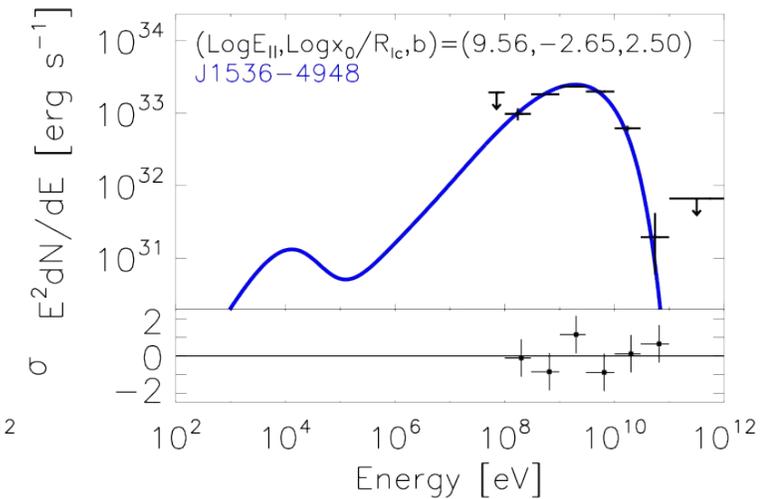
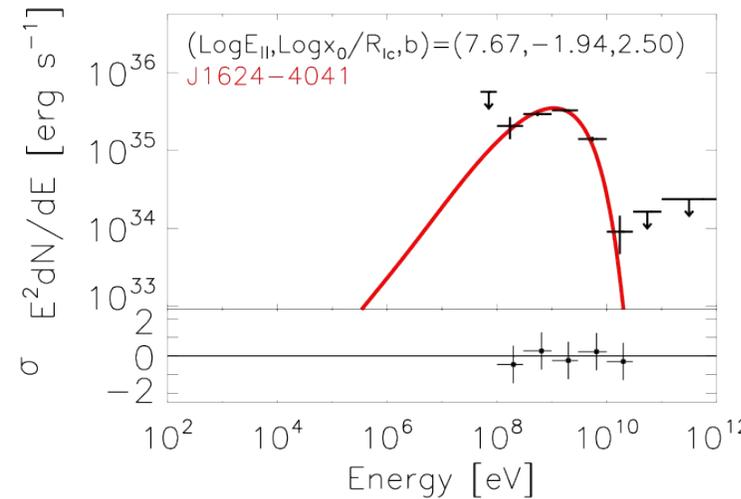
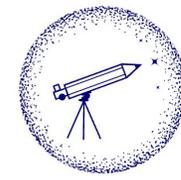


[Cheng & Zhang (1996), Viganò et al. 2015, MNRAS, 447, 1164]

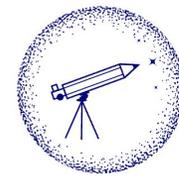
Spectral fitting

- The model successfully fitted the 117 gamma ray-pulsars on the 2PC (Abdo et al. (2013)) and the ~300 gamma-ray pulsars on the 3PC (Smith et al. (2023)). It can be extended to resemble the X-ray regime too, doing it in a majority of the ~40 high-energy pulsars

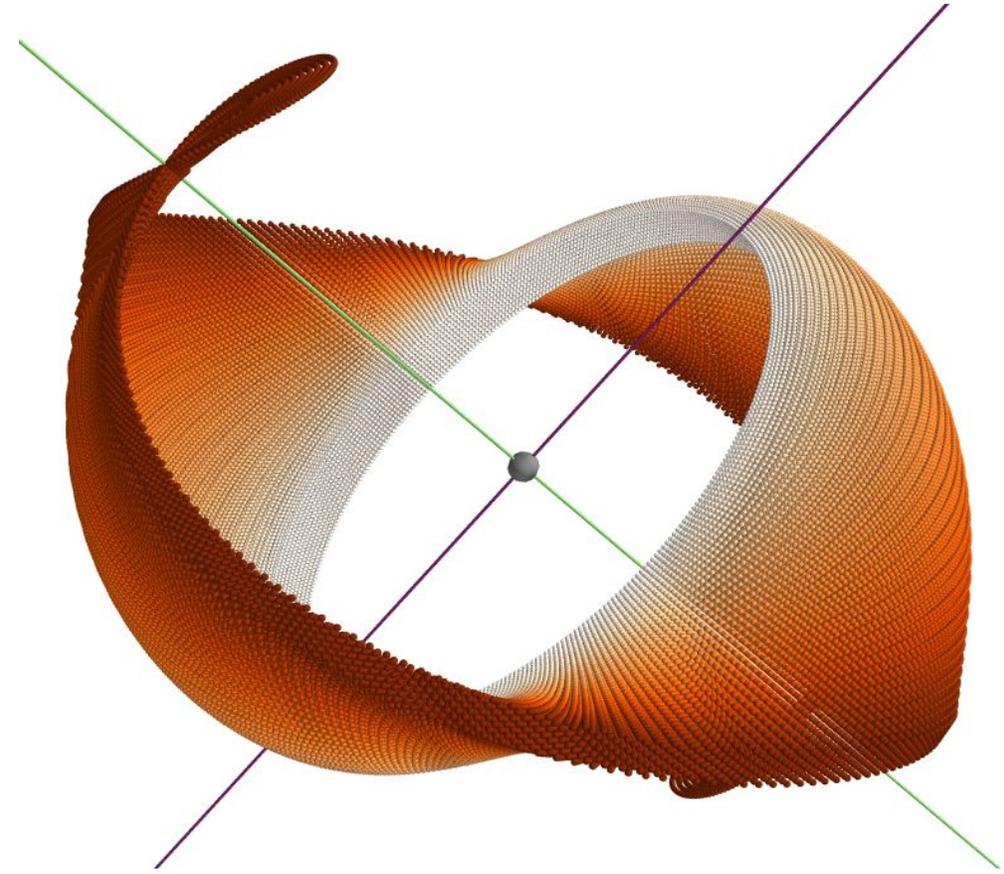
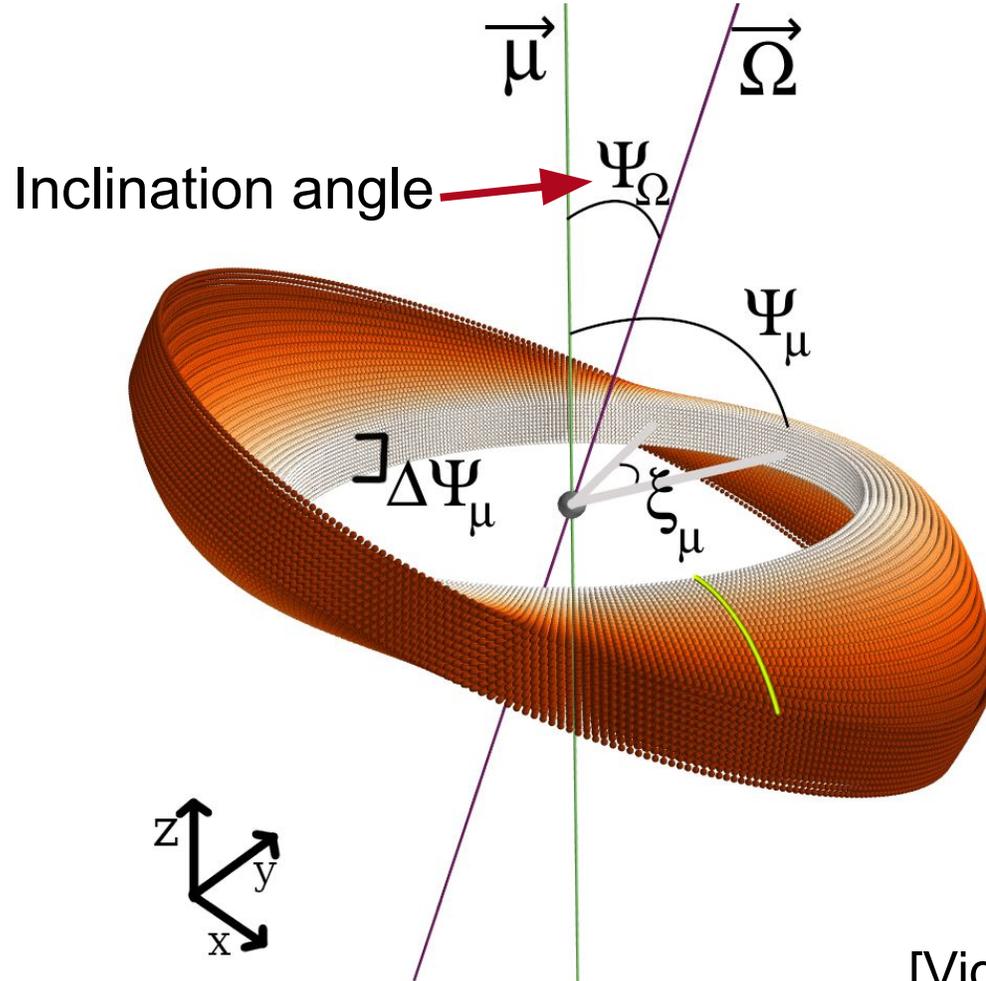
[Viganò et al. 2015, MNRAS, 453, 2599;
Torres 2018, Nat. Astron., 2, 247;
Torres et al. 2019, MNRAS, 489, 5494;
Íñiguez-Pascual, Viganò & Torres 2022, MNRAS, 516, 2475;
Íñiguez-Pascual, Torres & Viganò 2025, (submitted)]



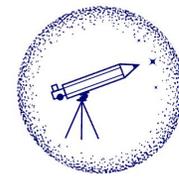
Geometrical model



- The inclination angle defines the geometry of the emission region



[Viganò & Torres (2019); Íñiguez-Pascual, Torres & Viganò (2024)]



Geometrical model: geometry formulae

- Frenet-Serret equations allow to geometrically describe a curved trajectory with torsion

$$\frac{d\hat{t}}{d\lambda} = \frac{1}{r_c} \hat{n} \quad \frac{d\hat{n}}{d\lambda} = -\frac{1}{r_c} \hat{t} + \tau \hat{b} \quad \frac{d\hat{b}}{d\lambda} = -\tau \hat{n}$$

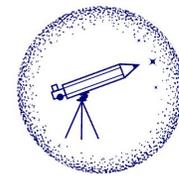
- The emission region is build around a centered value of the magnetic colatitude:

$$\Psi_{\mu}^c(\xi_{\mu}, R, \Psi_{\Omega}) = \frac{\pi}{2} + A(R, \Psi_{\Omega}) \sin(\xi_{\mu} - \pi/2) \quad \text{with} \quad A(R, \psi_{\Omega}) = K \Psi_{\Omega} (R/R_{lc} - R^0/R_{lc})^2$$

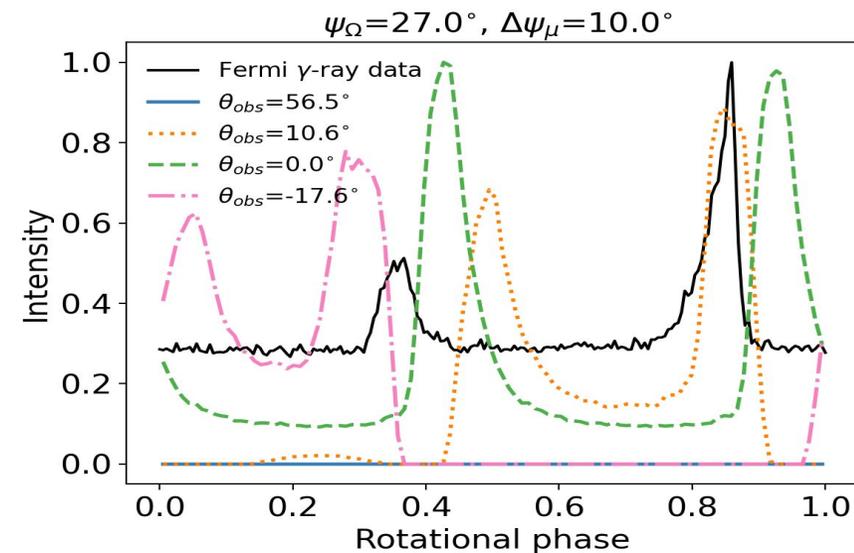
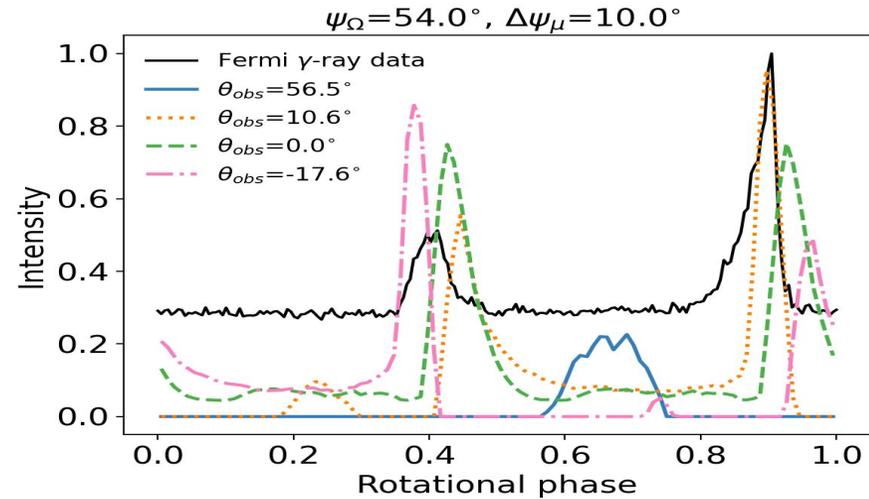
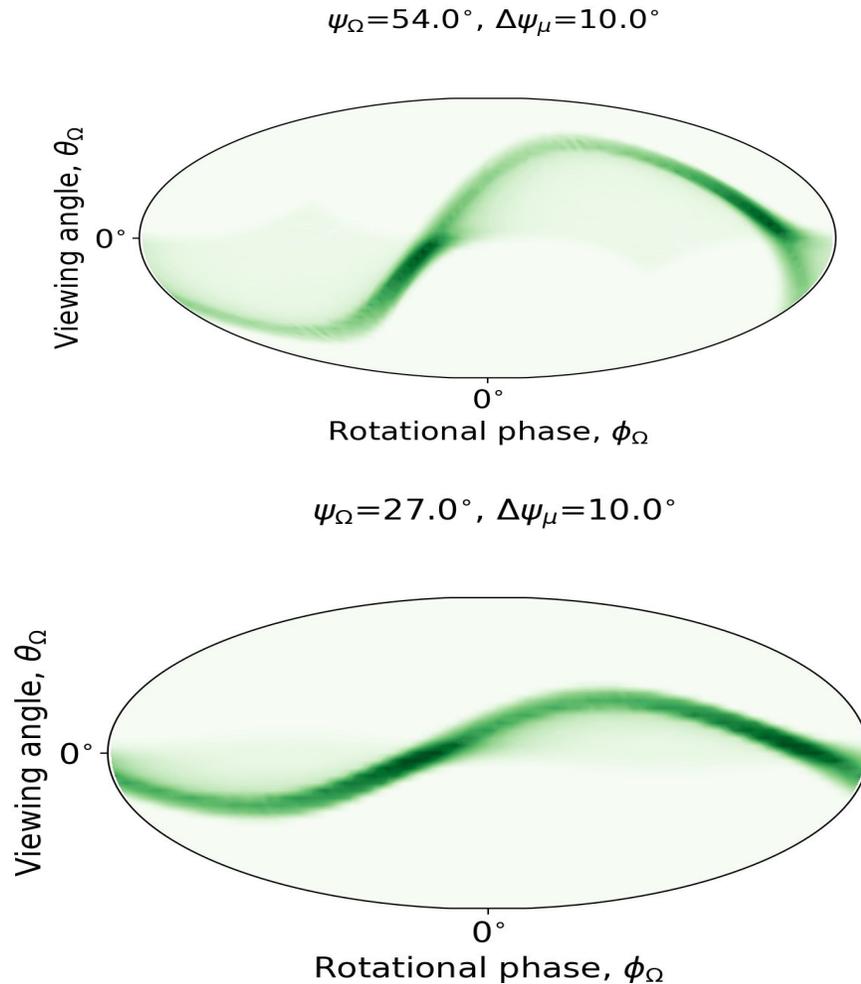
- Synchro-curvature emission maps are generated:

$$M_E(\theta_{\text{obs}}, \phi_{\Omega}, E) = \int_{R^0}^{R^0 + \Delta R} \int_{\lambda_{\text{in}}}^{\lambda_{\text{out}}} \int_0^{2\pi} \int_{\Delta\Psi_{\mu}} \frac{dD(\lambda)}{d\Omega_{\Omega}} \left\langle \frac{dP_{\text{sc}}(\lambda)}{dE} \right\rangle \frac{dN(\lambda)}{d\lambda} d\Psi_{\mu} d\xi_{\mu} d\lambda dx_{\text{in}}$$

Emission maps and light curves

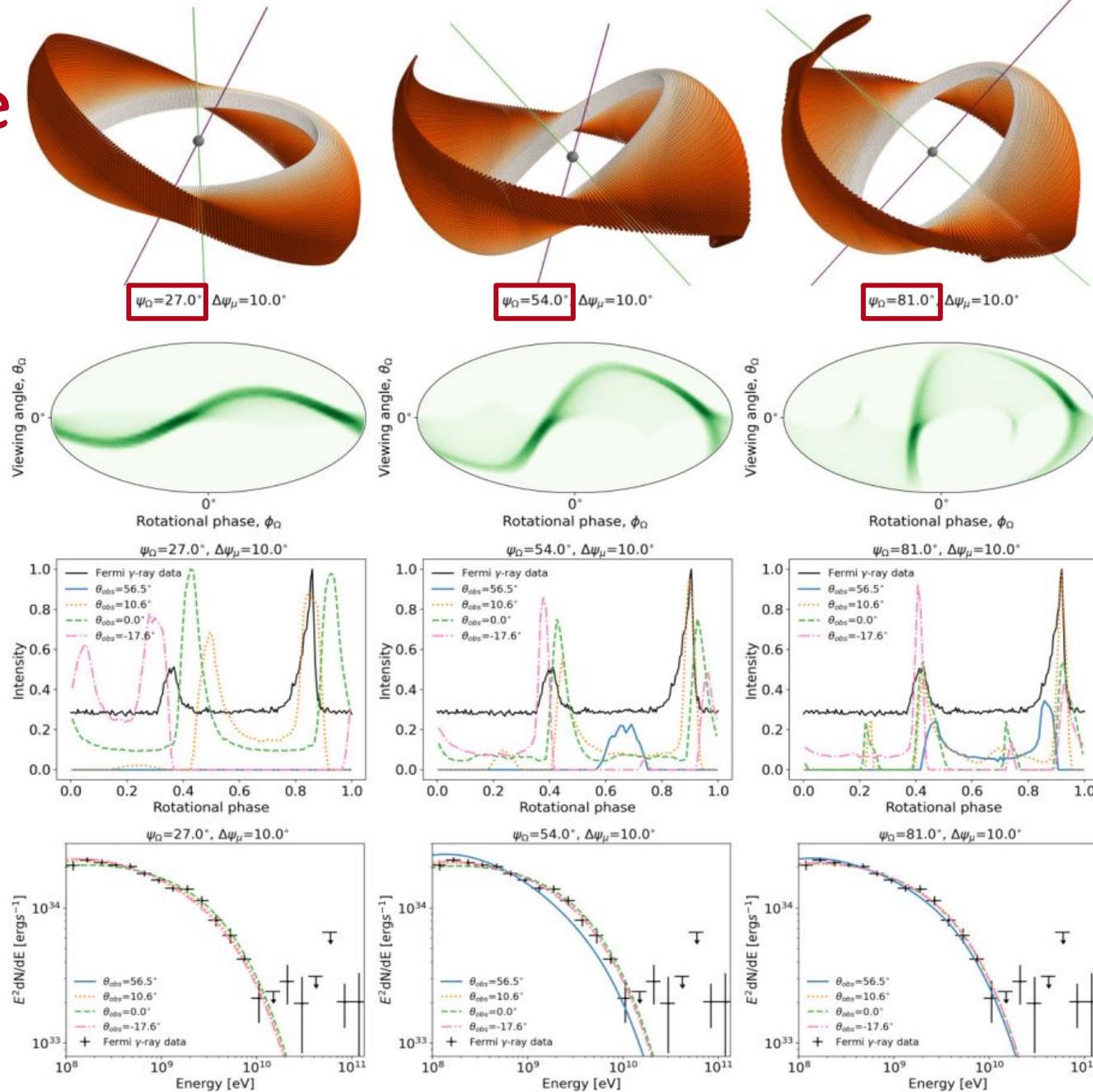
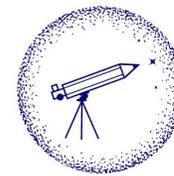


- We can build emission maps (skymaps), from which light curves are obtained



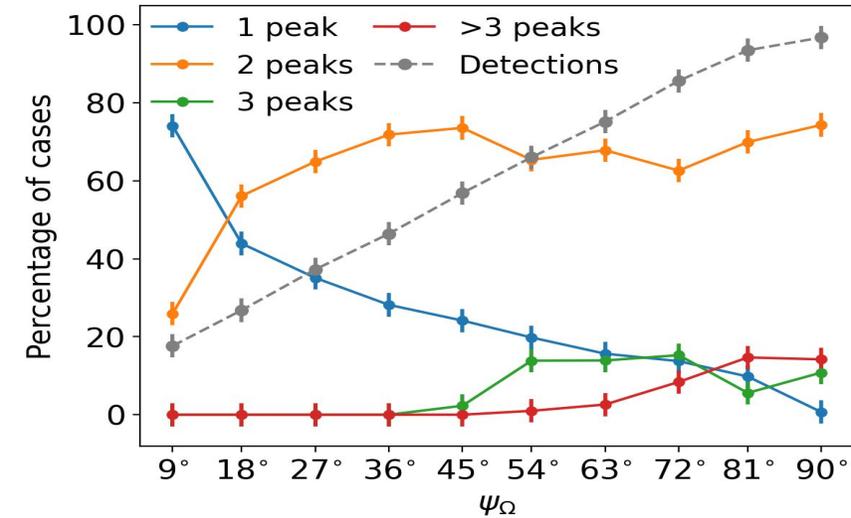
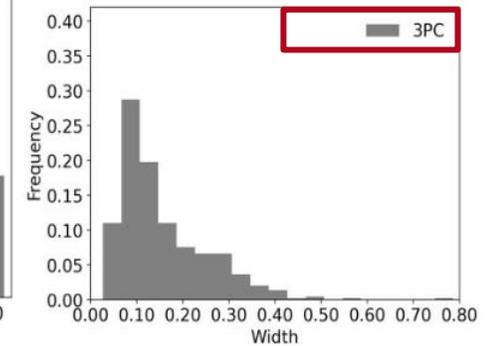
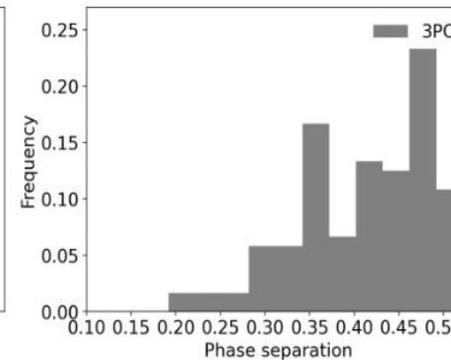
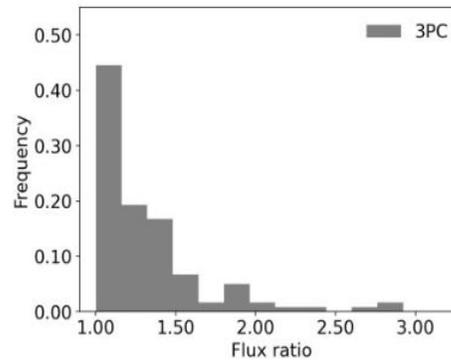
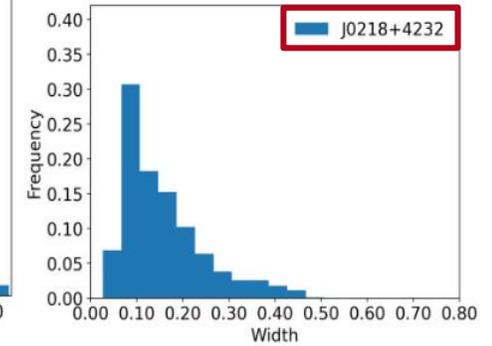
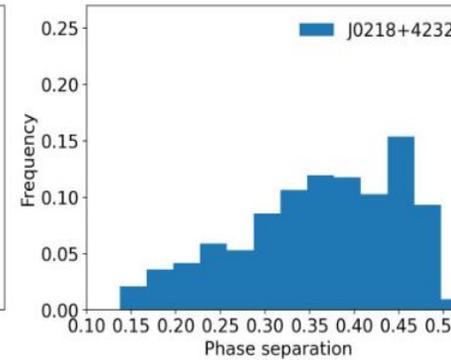
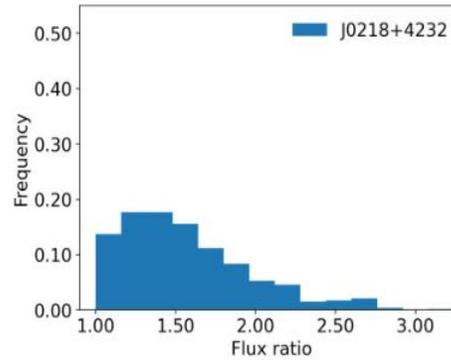
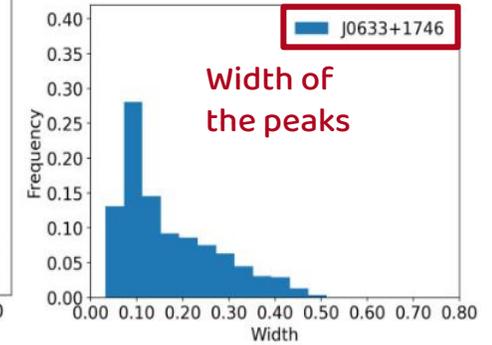
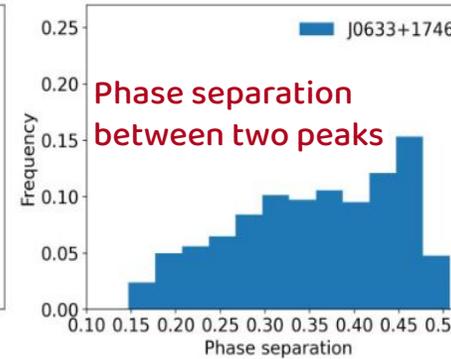
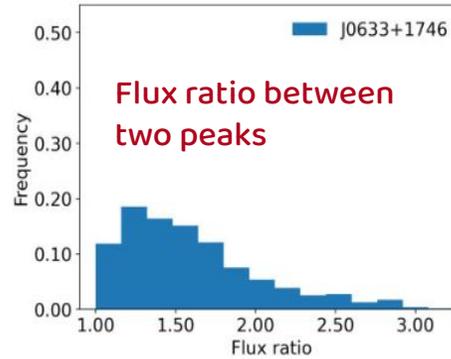
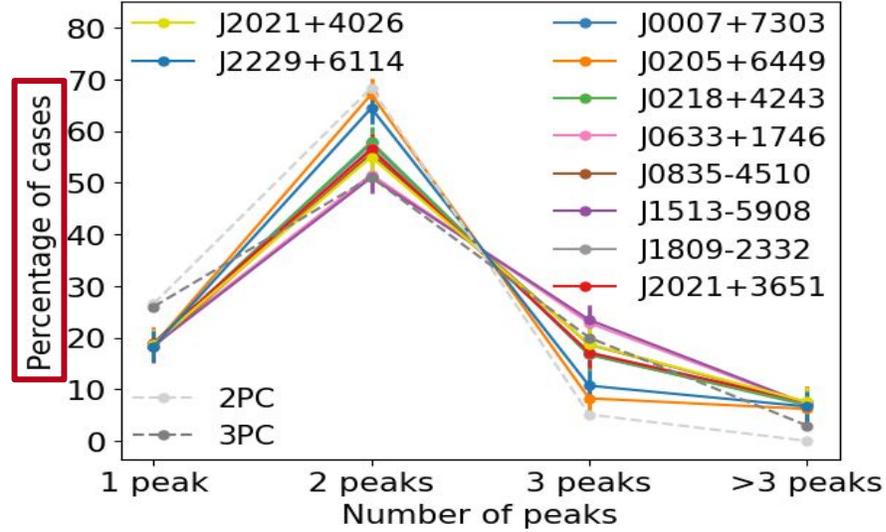
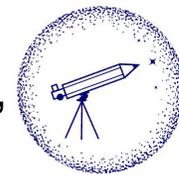
[Íñiguez-Pascual,
Torres & Viganò
2024, MNRAS,
530, 1550]

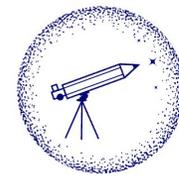
Outputs of the geometrical model



Light curves statistics

[Íñiguez-Pascual, Torres & Viganò 2024, MNRAS, 530, 1550]





Light curves fitting procedures

- Some considerations have to be made before light curve fitting: equal number of bins, background subtraction and normalization
- Light curve fitting is a non-trivial task, and there is not a perfect metric to be used
- We consider in parallel two methods:
 - A weighted reduced χ^2 in time domain

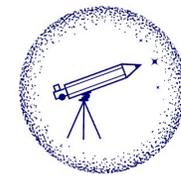
$$\overline{\chi_w^2} = \frac{1}{n-2} \sum_i \frac{(I_i^{obs} - I_i^{syn})^2 \bar{I}_i^{obs}}{(\delta I_i^{obs})^2}$$

- Euclidean distance in frequency domain

$$\|\hat{I}_k\| = K \left\| \sum_{j=0}^{n-1} I_j e^{-i2\pi k j/n} \right\| \quad ED = \sqrt{\sum_k \left(\|\hat{I}_k^{obs}\| - \|\hat{I}_k^{syn}\| \right)^2}$$

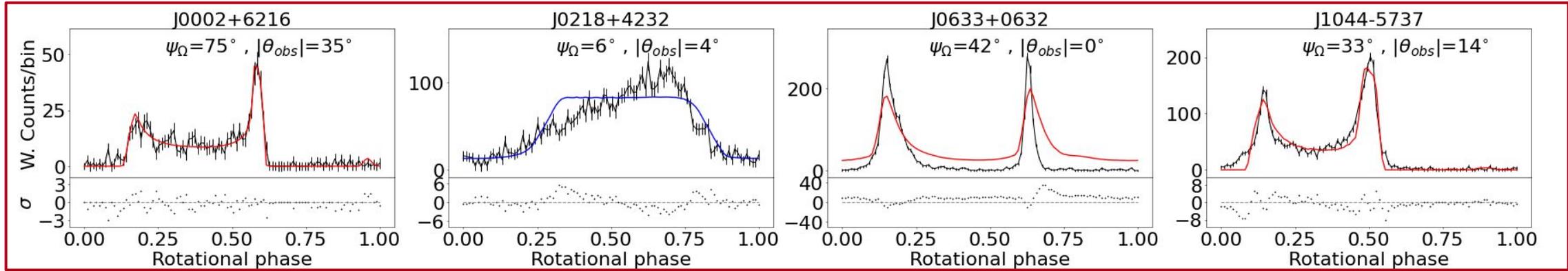
Light curves fitting

[Íñiguez-Pascual et al. 2025 (submitted)]

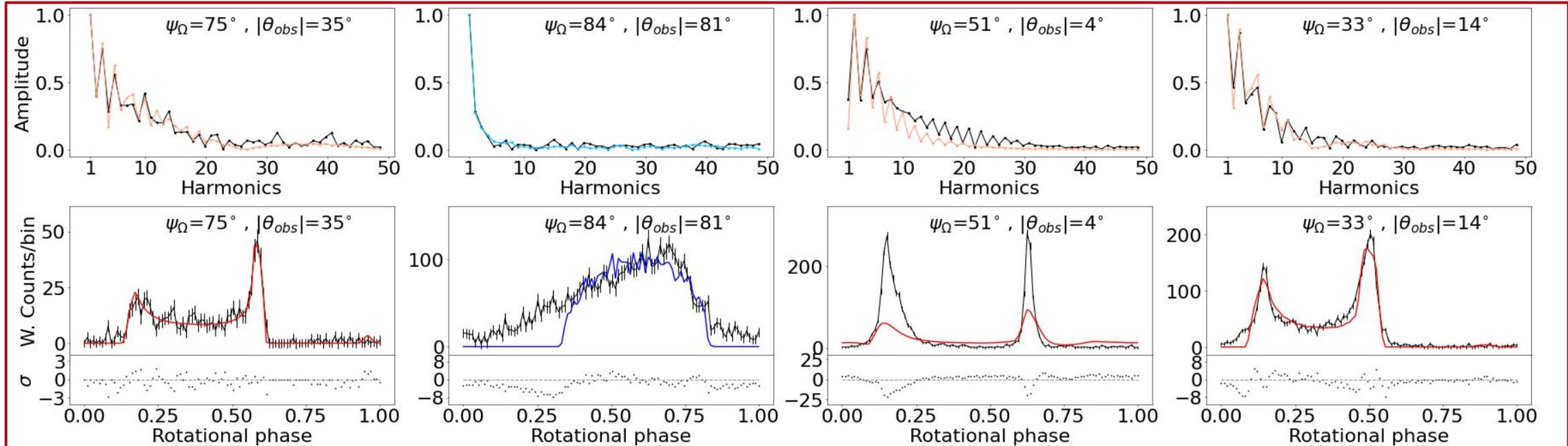


- Fitting synthetic light curves to observational ones, concurrently to the spectral fitting

X2 in time domain

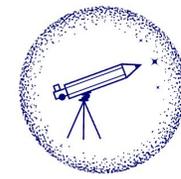


ED in frequency domain



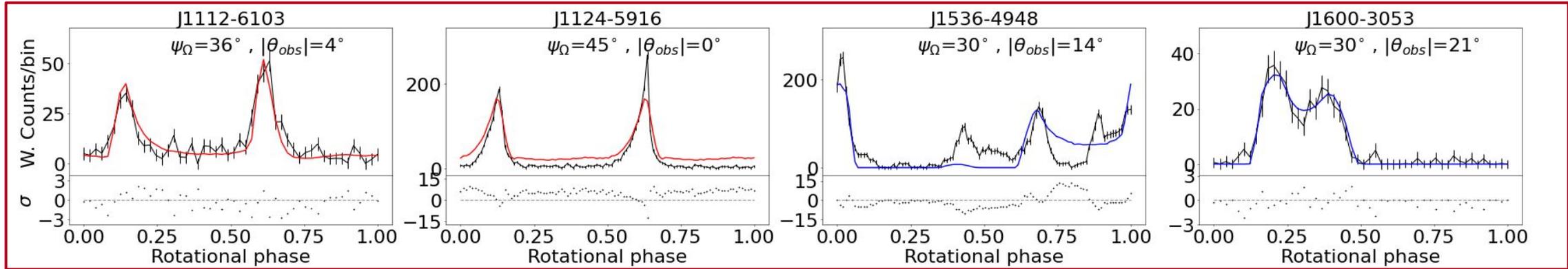
Light curves fitting

[Íñiguez-Pascual et al. 2025 (submitted)]

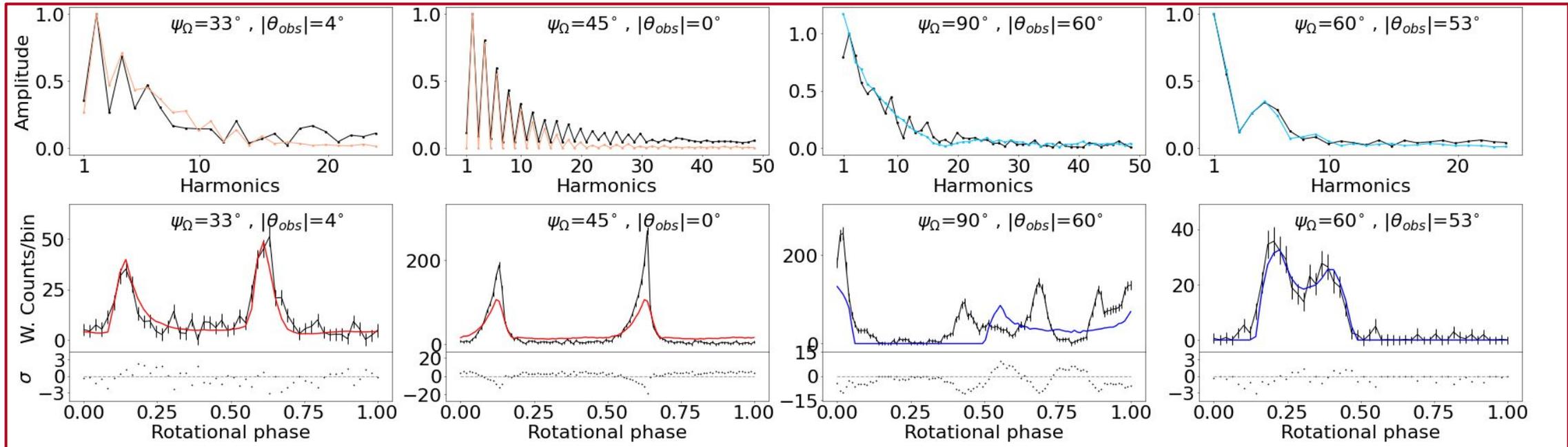


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X2 in time domain

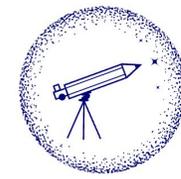


ED in frequency domain



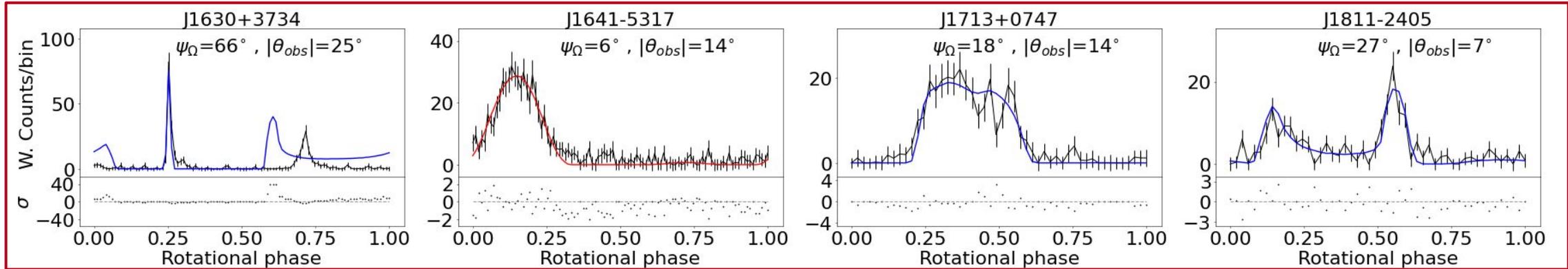
Light curves fitting

[Íñiguez-Pascual et al. 2025 (submitted)]

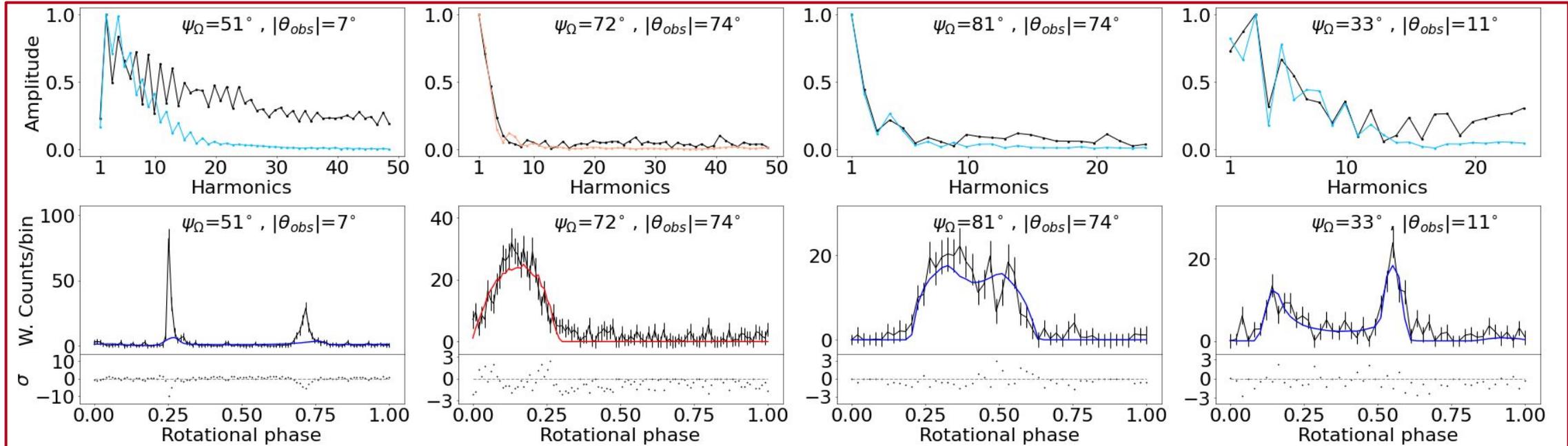


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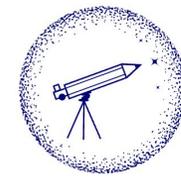


ED in frequency domain

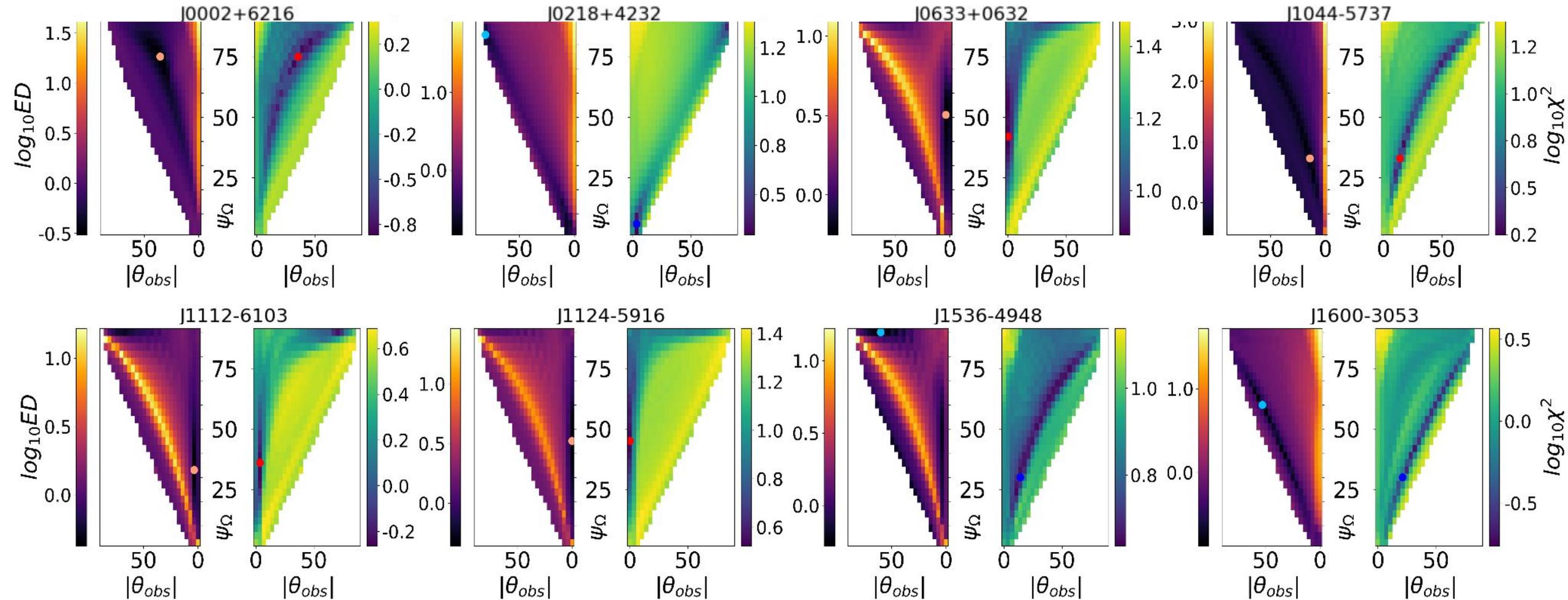


Light curves fitting

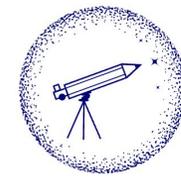
[Íñiguez-Pascual et al. 2025 (submitted)]



- Fitting synthetic light curves to observational ones, concurrently to the spectral fitting



Conclusions



- Concurrently fitting spectra and light curves of gamma-ray pulsars
- With three free spectral parameters, our spectral model is still able to fit the whole high-energy pulsar population
- We properly fit a number of gamma-ray light curves, involving just two free geometrical parameters
- Two different methods are used for the light curve fitting: a weighted reduced χ^2 in time domain and an euclidean distance in frequency domain, both with qualitatively similar results
- We will improve the model by including more realistic physics while keeping our effective approach

[Íñiguez-Pascual D., Viganò D., Torres D. F., 2022, MNRAS, 516, 2475 \(2208.05549\)](#)

[Íñiguez-Pascual D., Torres D. F., Viganò D., 2024, MNRAS, 530, 1550 \(2404.01926\)](#)

[Íñiguez-Pascual D., Torres D. F., Viganò D., 2025, submitted \(2504.01892\)](#)

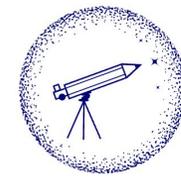
Thank you



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Spectral model: particle dynamics formulae



- The equation of motion of charged particles balances electric acceleration and synchro-curvature losses

$$\frac{d\vec{p}}{dt} = ZeE_{\parallel}\hat{b} - (P_{sc}/v)\hat{p}$$

$$\vec{p} = \Gamma m\vec{v}$$

[Viganò et al. (2015)]

- We solve it numerically, considering separately the components parallel and perpendicular to the trajectory

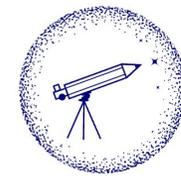
$$\frac{d(p \cos \alpha)}{dt} = ZeE_{\parallel} - \frac{P_{sc}}{v} \cos \alpha$$

$$\frac{d(p \sin \alpha)}{dt} = - \frac{P_{sc}}{v} \sin \alpha$$

- Local magnetic field strength and curvature radius are parametrized in an effective way:

$$B = B_{\star} \left(\frac{R_{\star}}{x} \right)^b \quad r_c = R_{lc} \left(\frac{x}{R_{lc}} \right)^{\eta}$$

Spectral model: synchro-curvature formulae



- Single-particle synchro-curvature power spectra:

$$\frac{dP_{sc}}{dE} = \frac{\sqrt{3}(Ze)^2\Gamma y}{4\pi\hbar r_{eff}} [(1+z)F(y) - (1-z)K_{2/3}(y)]$$

- Synchro-curvature power:

$$P_{sc} = \frac{2(Ze)^2\Gamma^4 c}{3r_c^2} g_r$$

[Cheng & Zhang (1996),
Viganò et al. (2015)]

$$r_{eff} = \frac{r_c}{\cos^2 \alpha} \left(1 + \xi + \frac{r_{gyr}}{r_c} \right)^{-1}$$

$$z = (Q_2 r_{eff})^{-2}$$

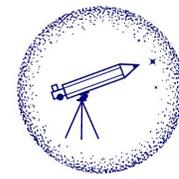
$$y = \frac{E}{E_c}$$

$$Q_2^2 = \frac{\cos^4 \alpha}{r_c^2} \left[1 + 3\xi + \xi^2 + \frac{r_{gyr}}{r_c} \right]$$

$$E_c = \frac{3}{2} \hbar c Q_2 \Gamma^3$$

$$F(y) = \int_y^\infty K_{5/3}(y') dy'$$

$$\xi = \frac{r_c}{r_{gyr}} \frac{\sin^2 \alpha}{\cos^2 \alpha}$$

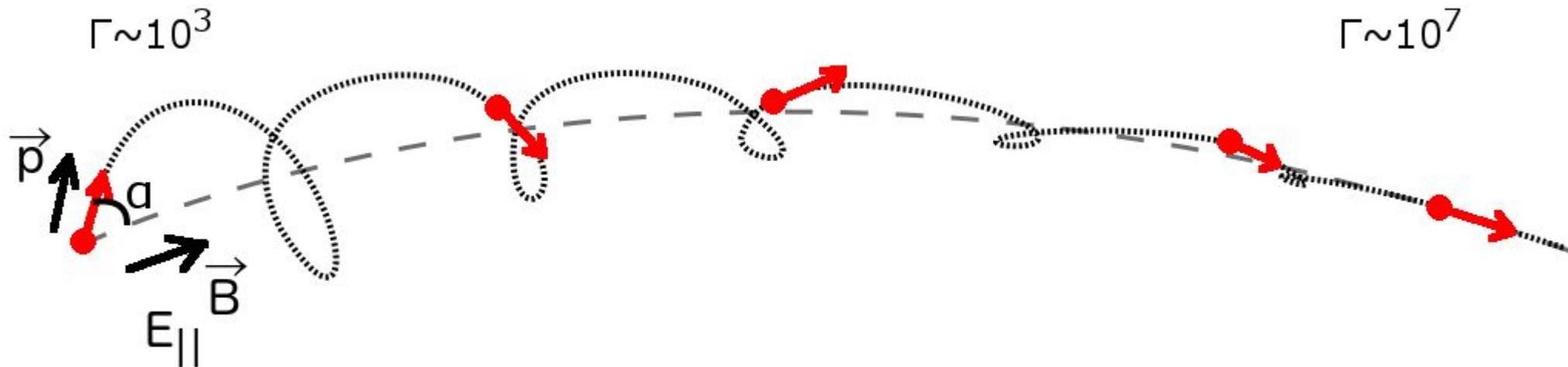


Back up slide: particle dynamics

- Power of synchrotron and curvature radiations depend on the Lorentz factor Γ of particles differently

$$P_{syn} = \frac{2}{3} \frac{(Ze)^4 B^2 (\Gamma^2 - 1) \sin^2 \alpha}{m^2 c^3}$$

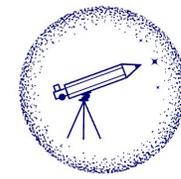
$$P_c = \frac{2}{3} \frac{(Ze)^2 c \Gamma^4}{r_c^2}$$



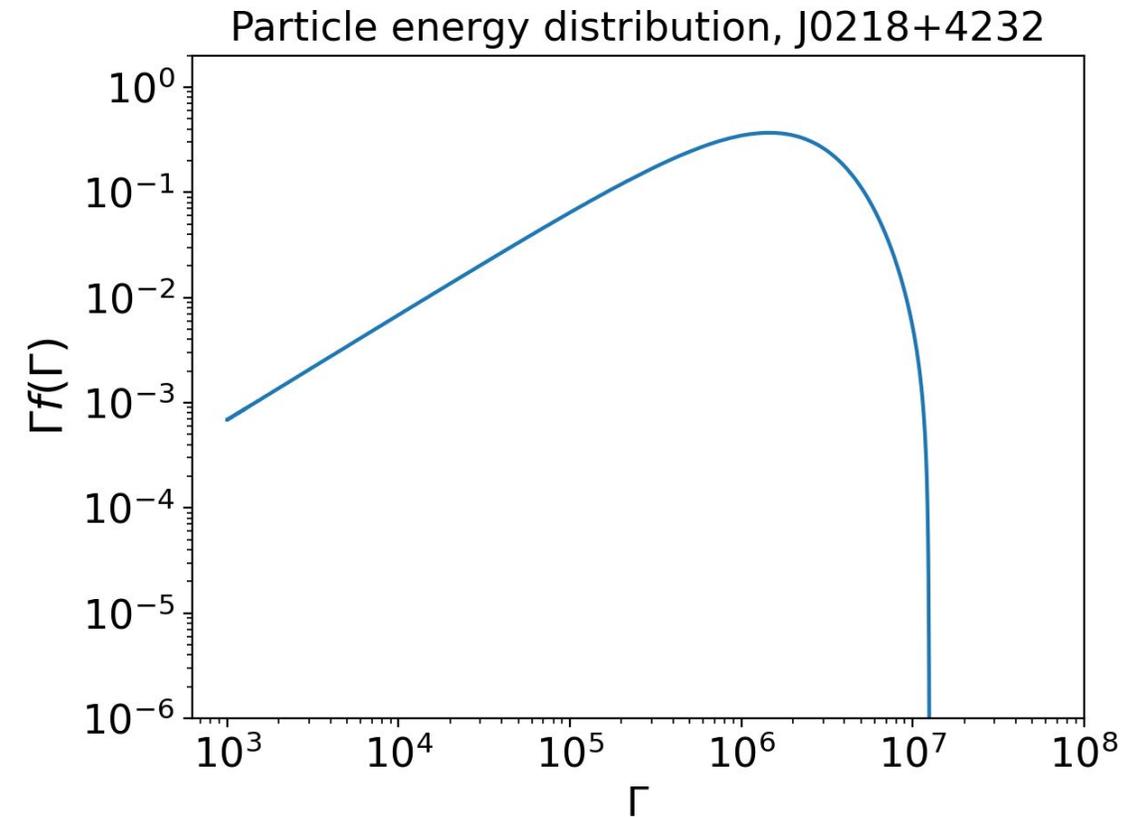
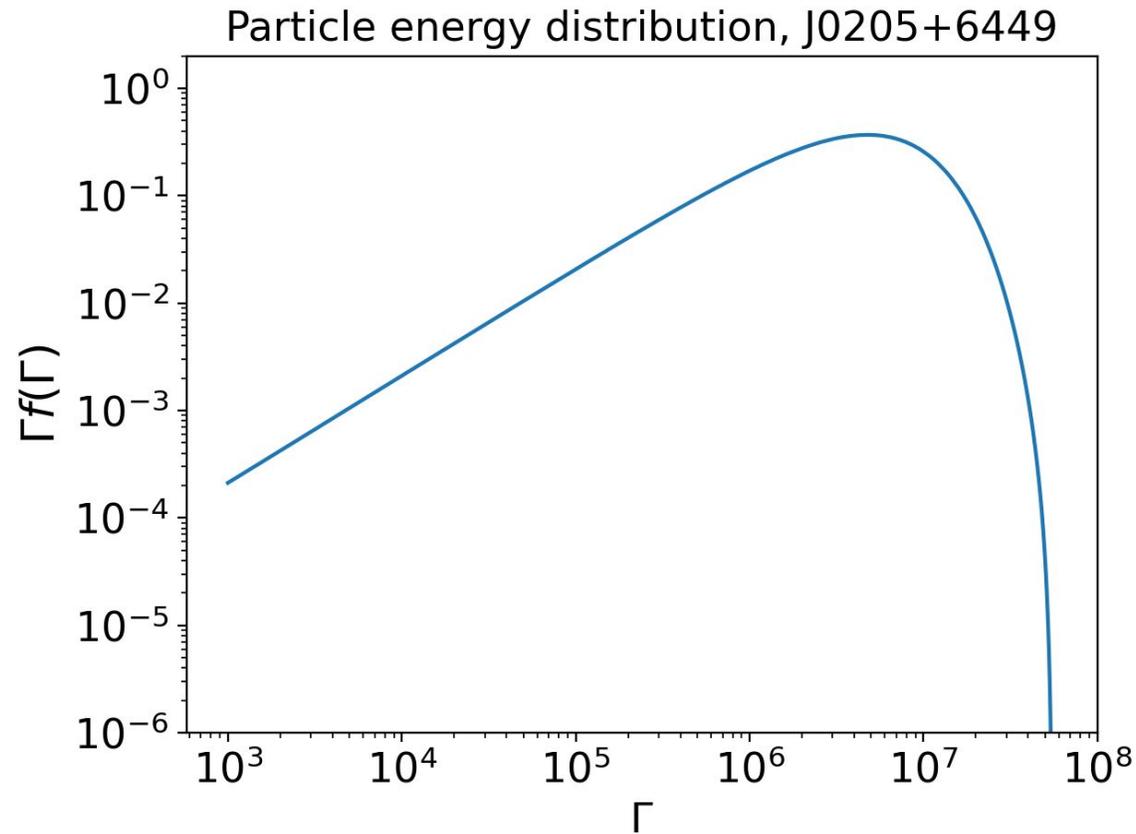
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[Viganò et al. 2015, MNRAS, 447, 1164]

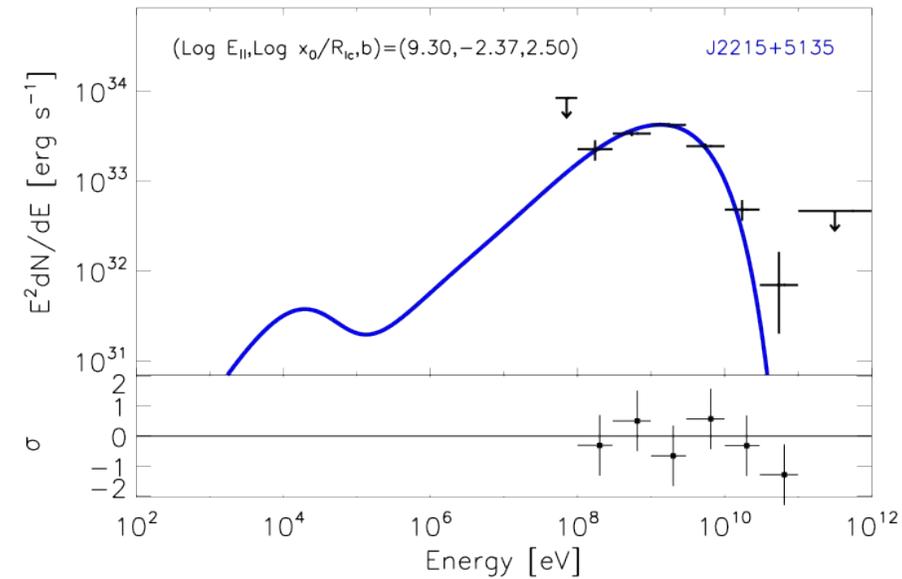
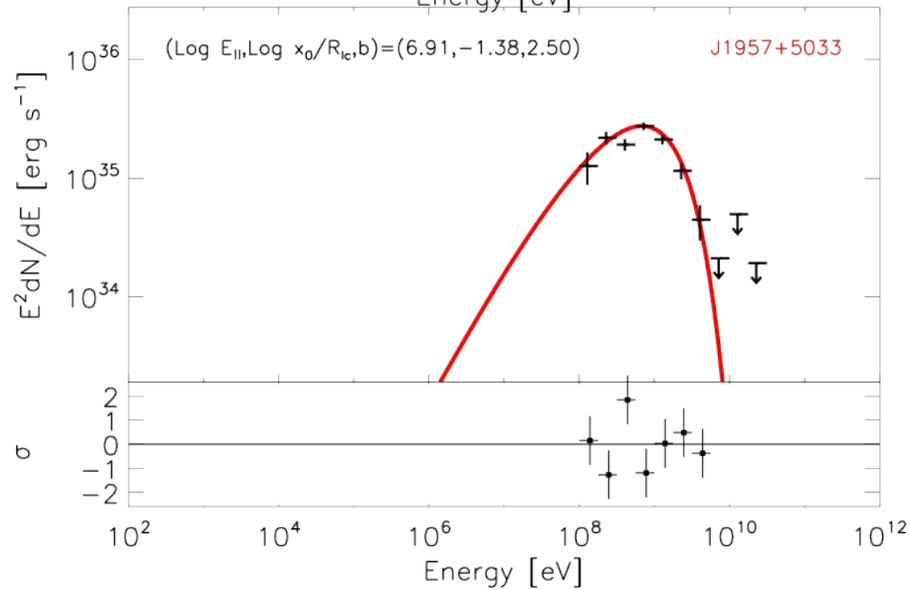
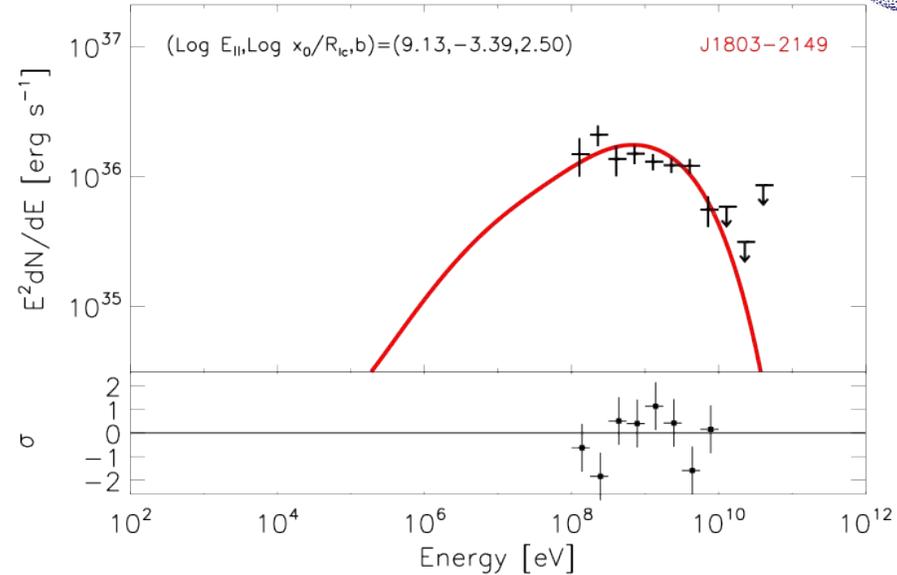
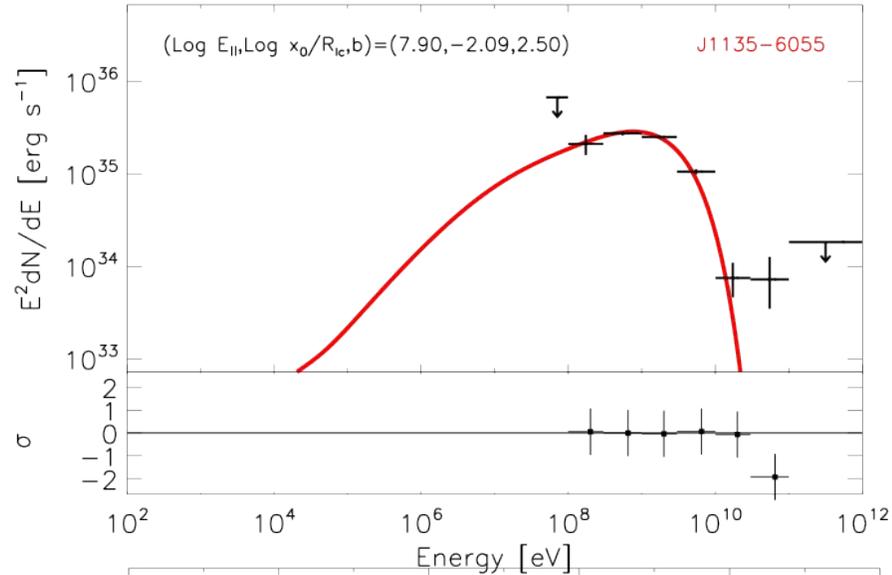
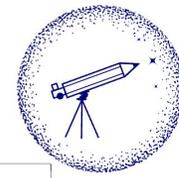
Back up slide: particle energy distribution



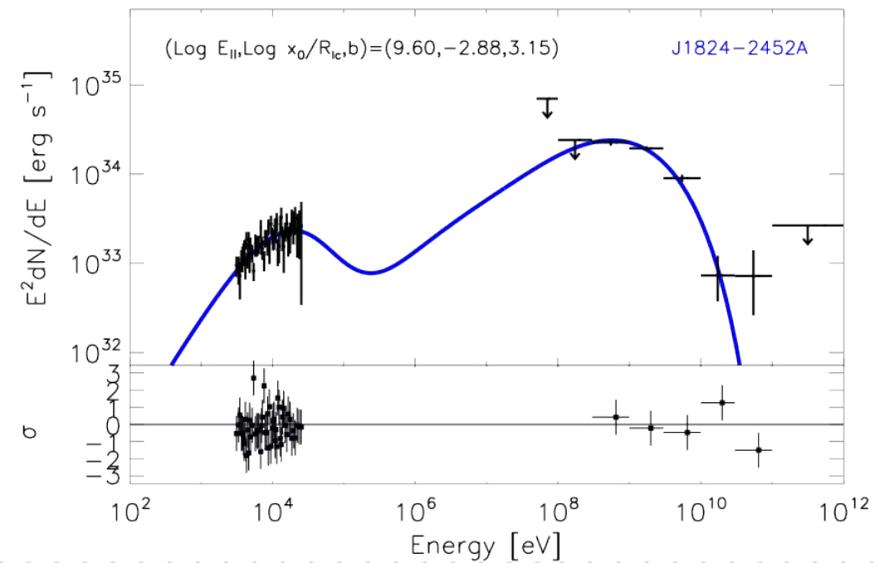
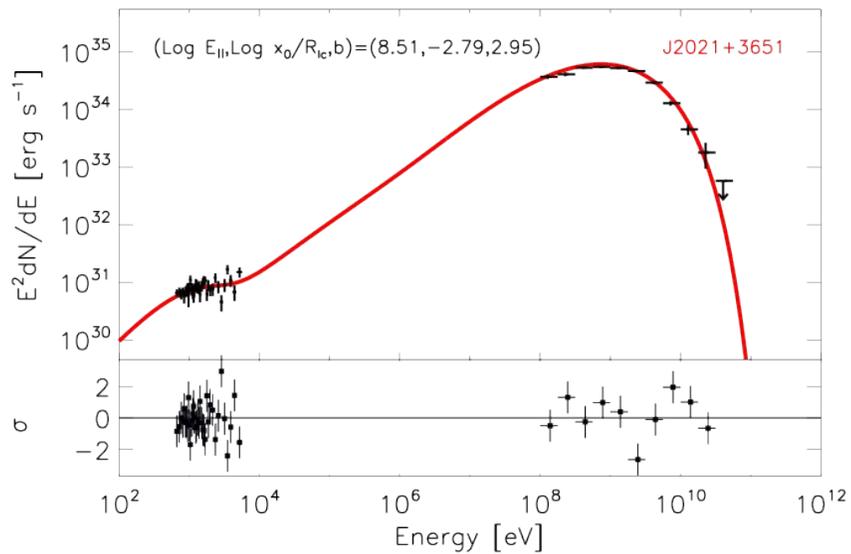
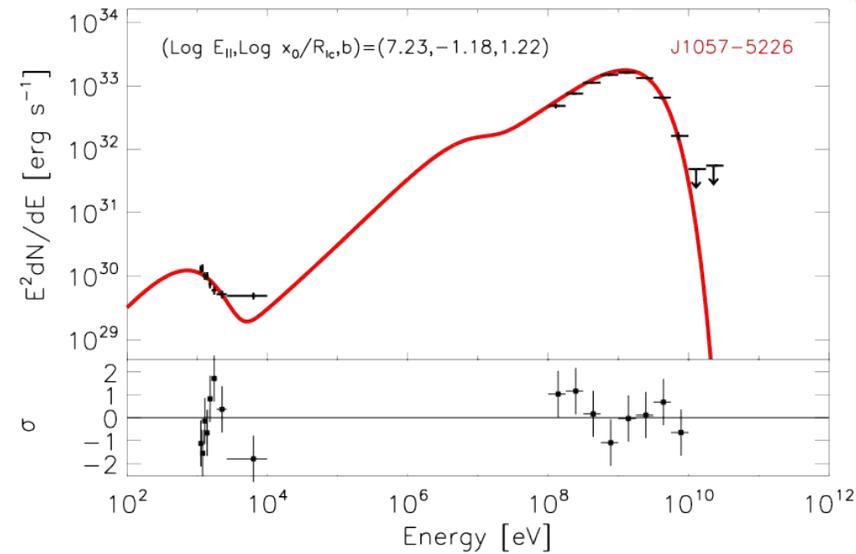
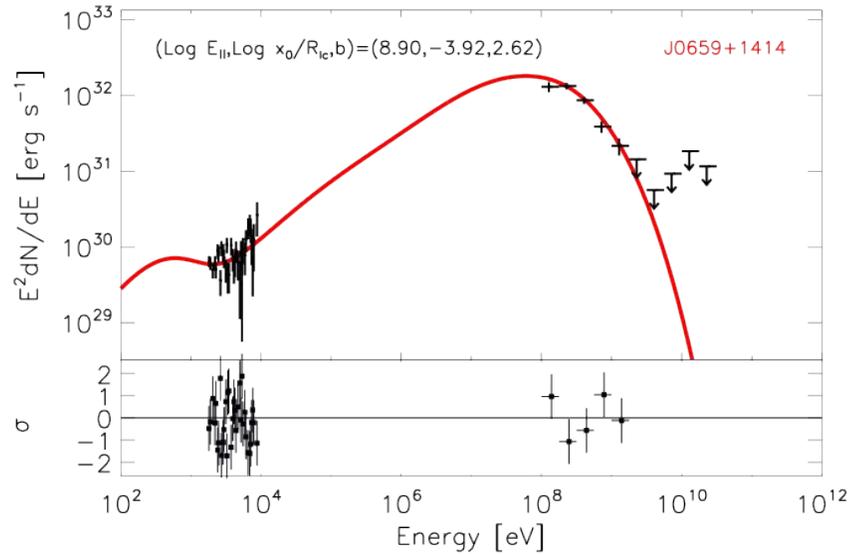
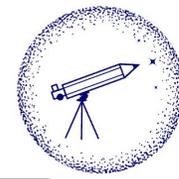
- Lorentz factors Γ typically range from 10^3 to 10^7 .



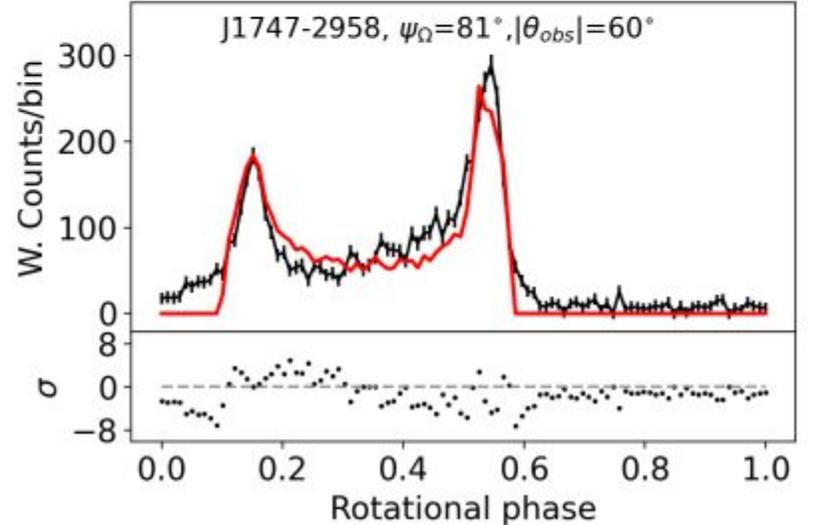
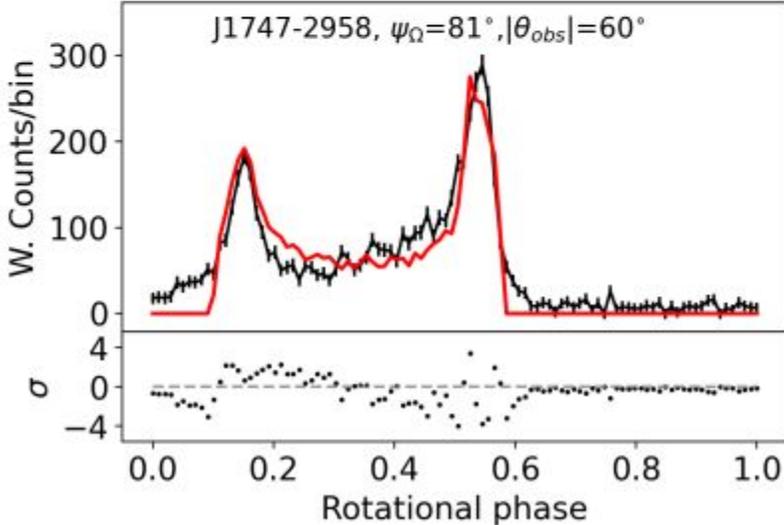
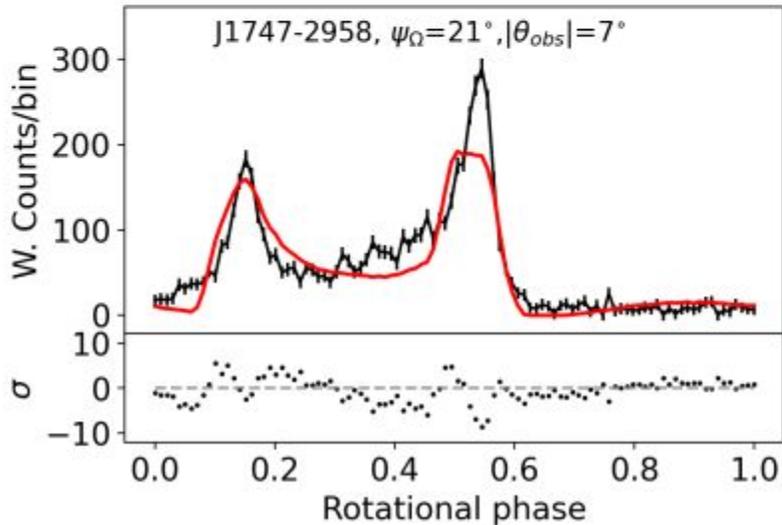
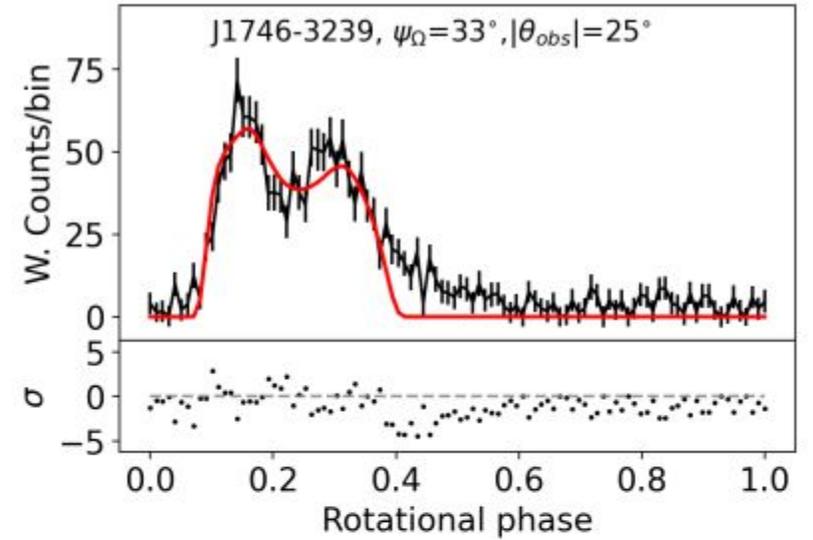
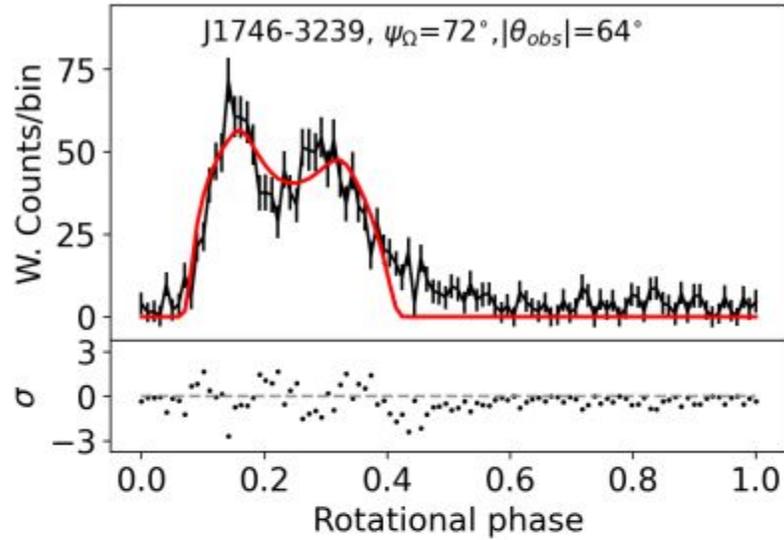
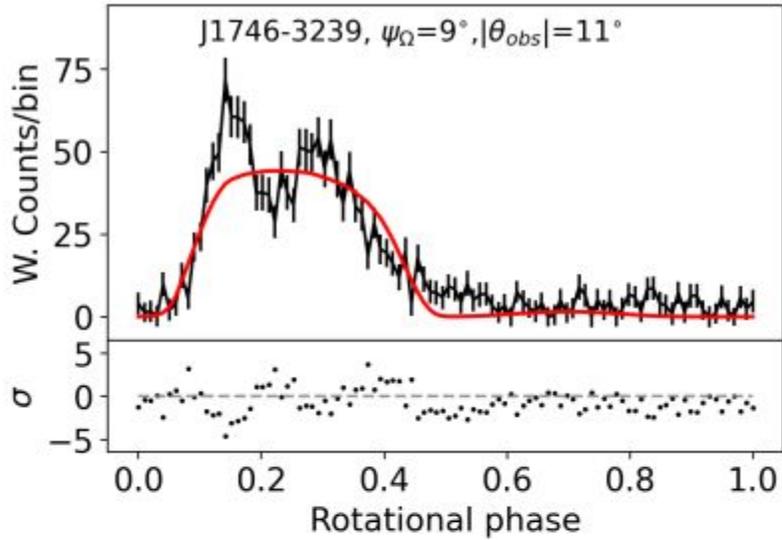
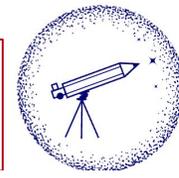
Back up slide: more spectral fits



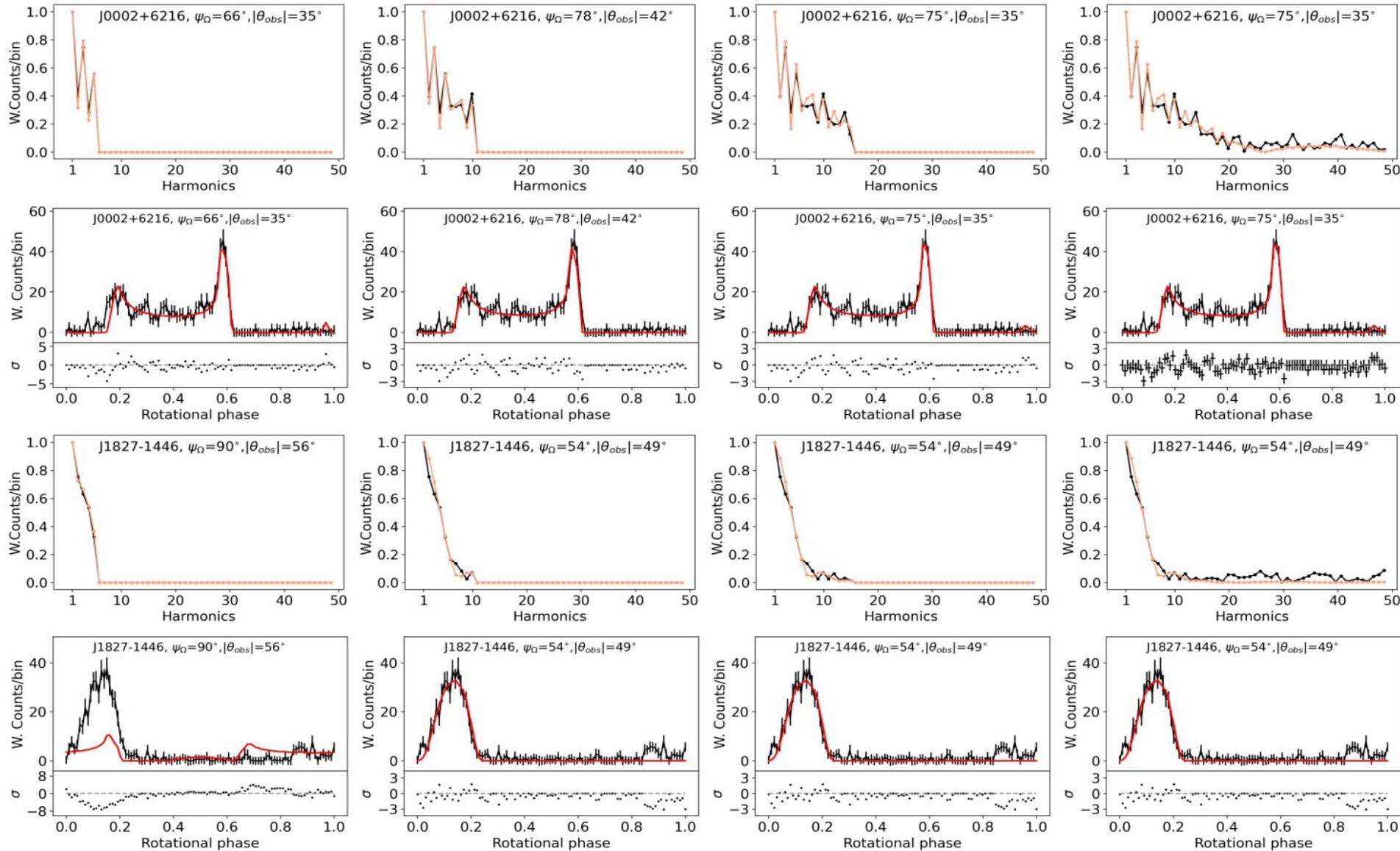
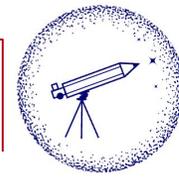
Back up slide: more spectral fits



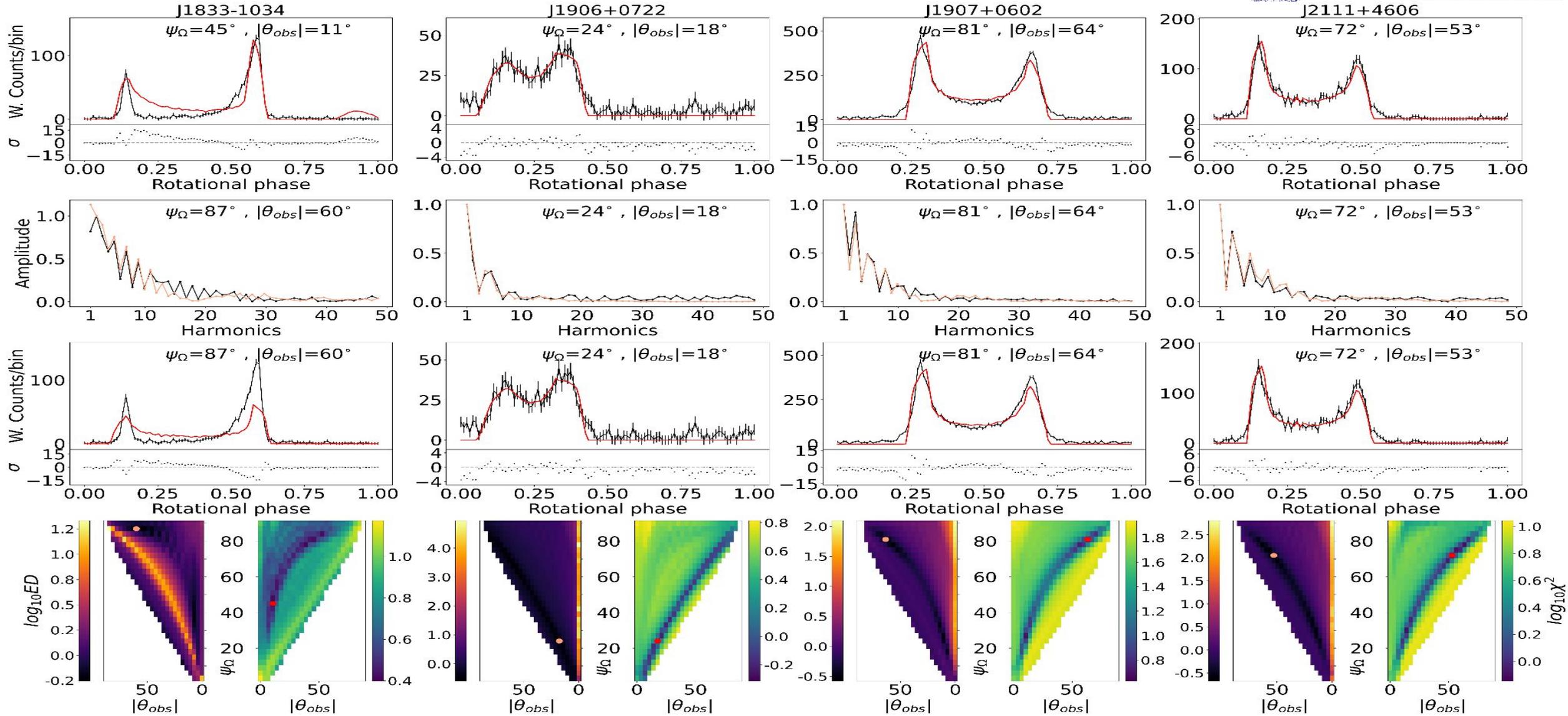
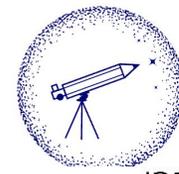
Back up slide: caveats



Back up slide: filtering out Fourier



Back up slide: more light curve fits



Back up slide: Degeneracy of geometrical parameters

