

Ultra-long Period Magnetars

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Andrey Timokhin (Zelona Gora)

Alice Harding (LANL)

Matthew Baring (Rice)

Jeremy Hare (NASA GSFC)

+ many others

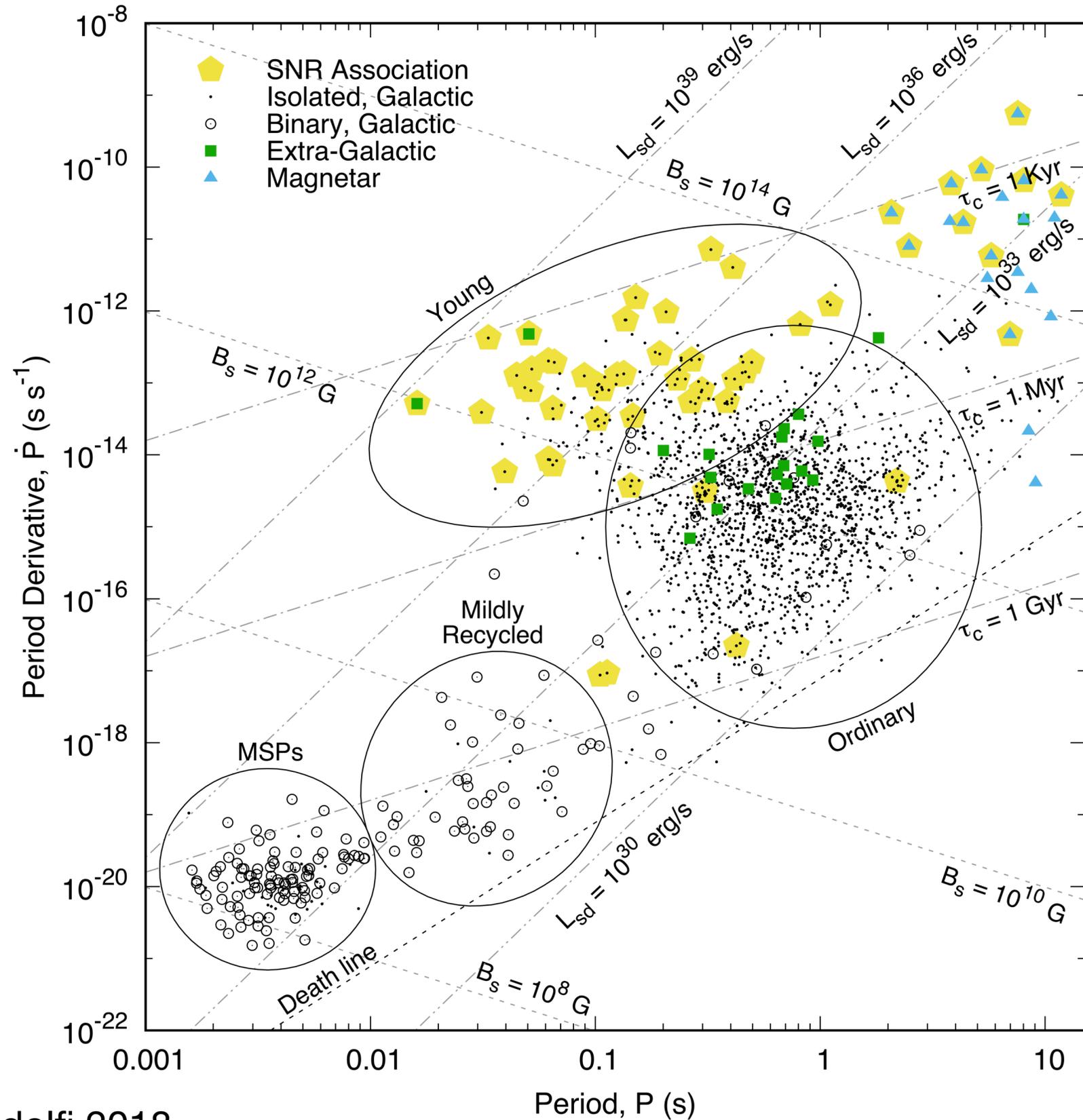
Relativistic Magnetospheres, Les Houches School Workshop

Les Houches, France, April 9, 2025

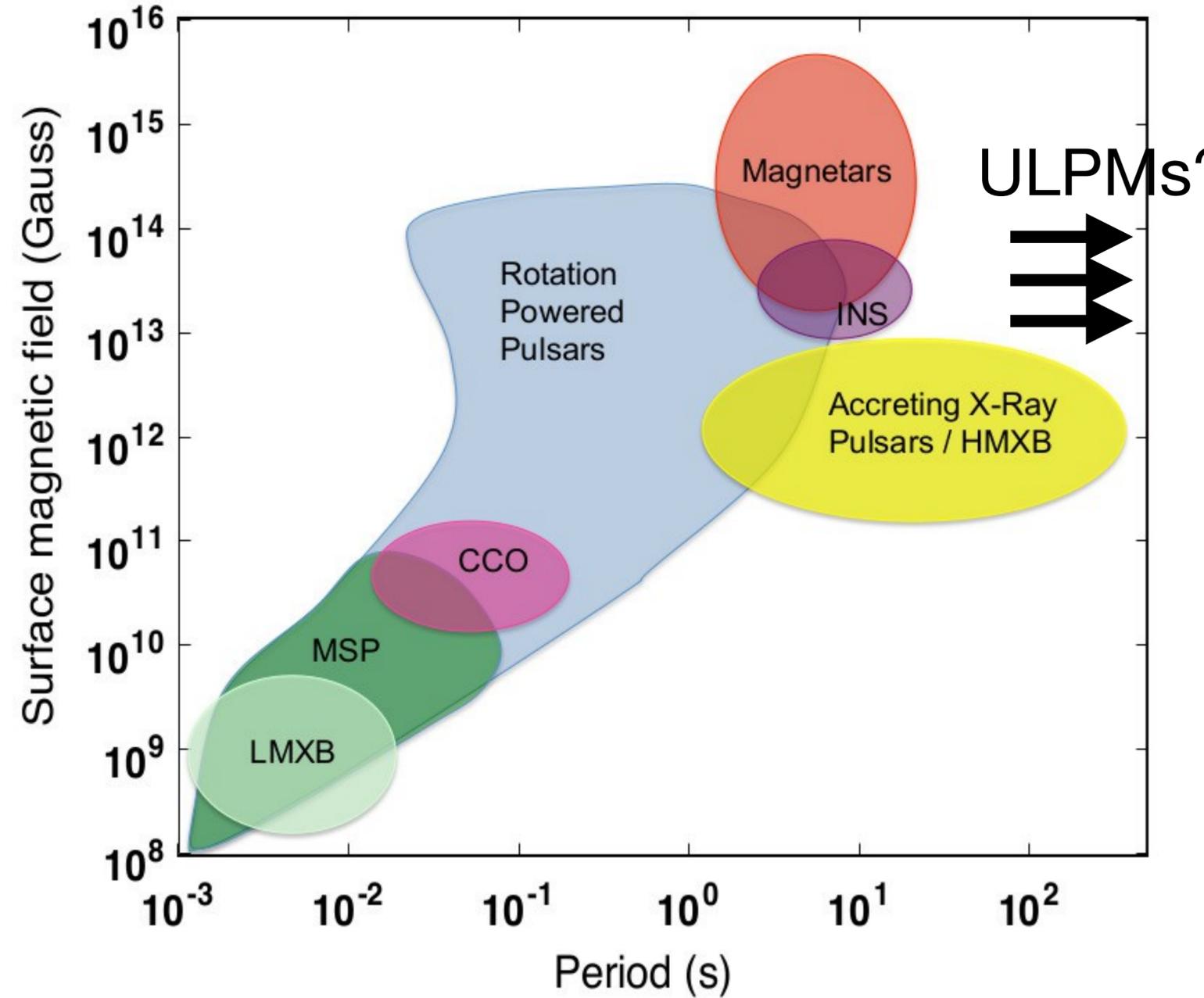
Outline

- Why and what
 - Why is this exciting
 - What: The zoo of long period radio transients and recent discoveries
- How
 - slow rotation magnetars, small twists, QED pair cascades, QED resonant Compton versus curvature photon gaps & death lines

The Neutron Star Zoo

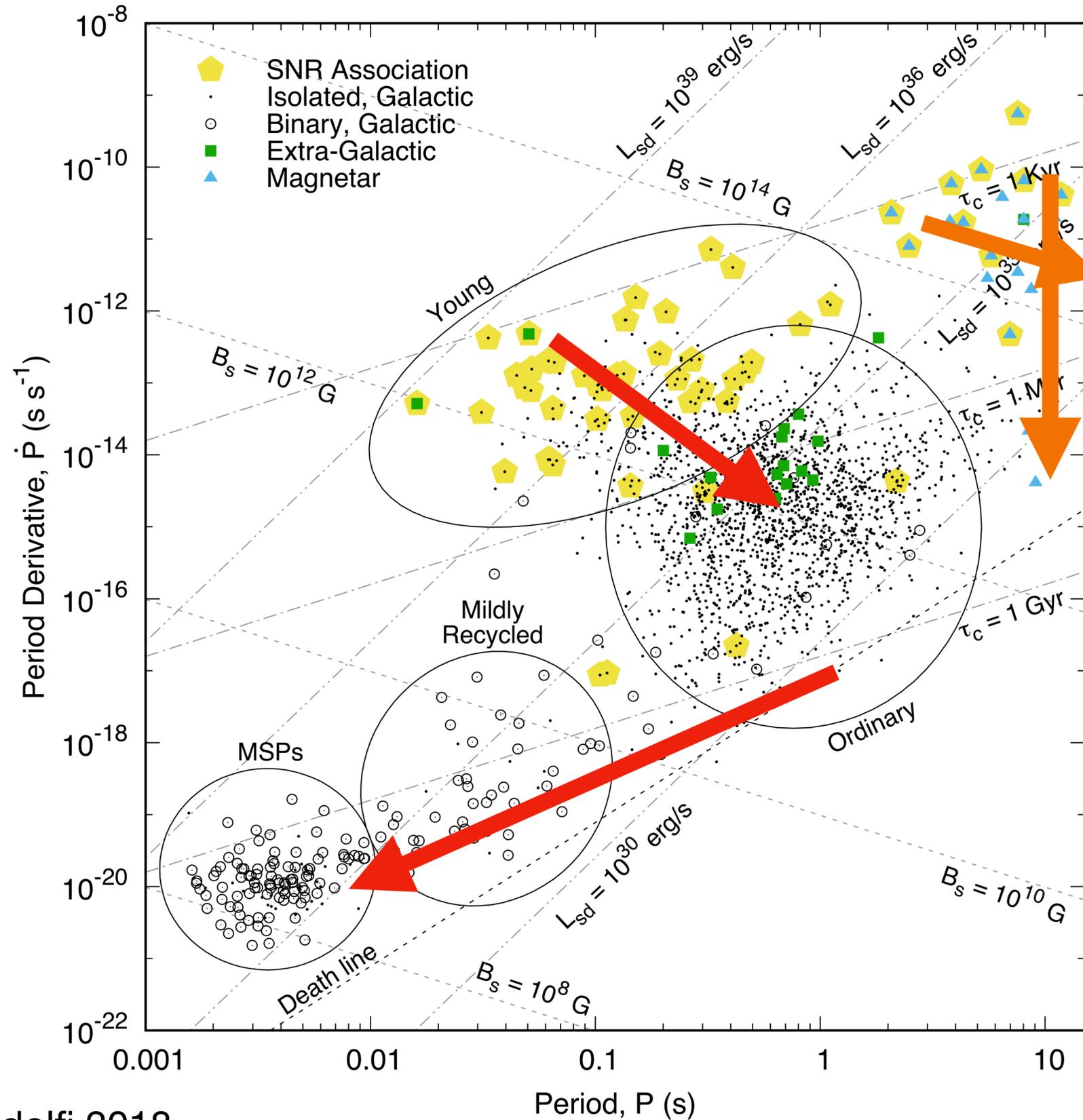


Ridolfi 2018

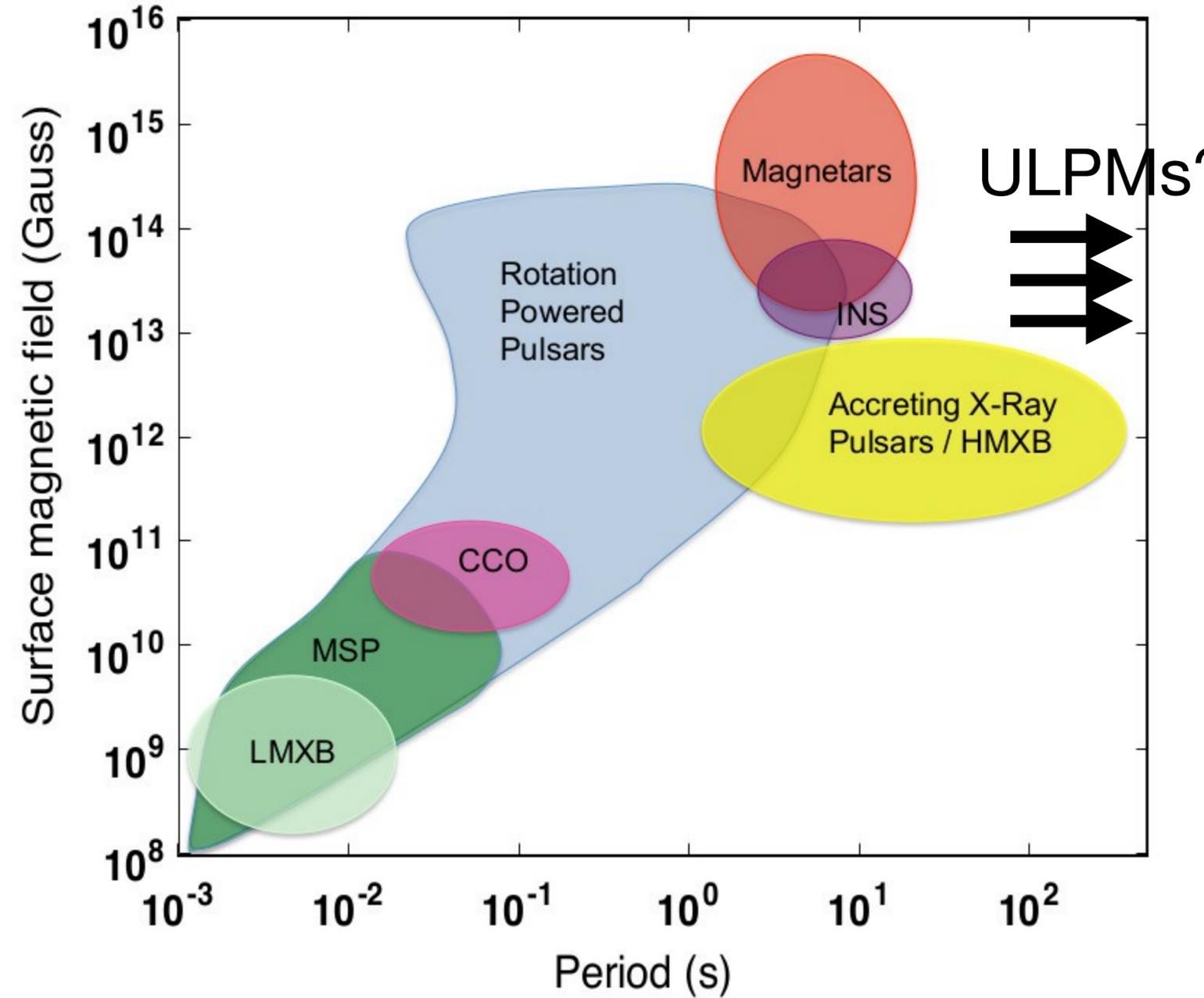


Harding 2013

The Neutron Star Zoo

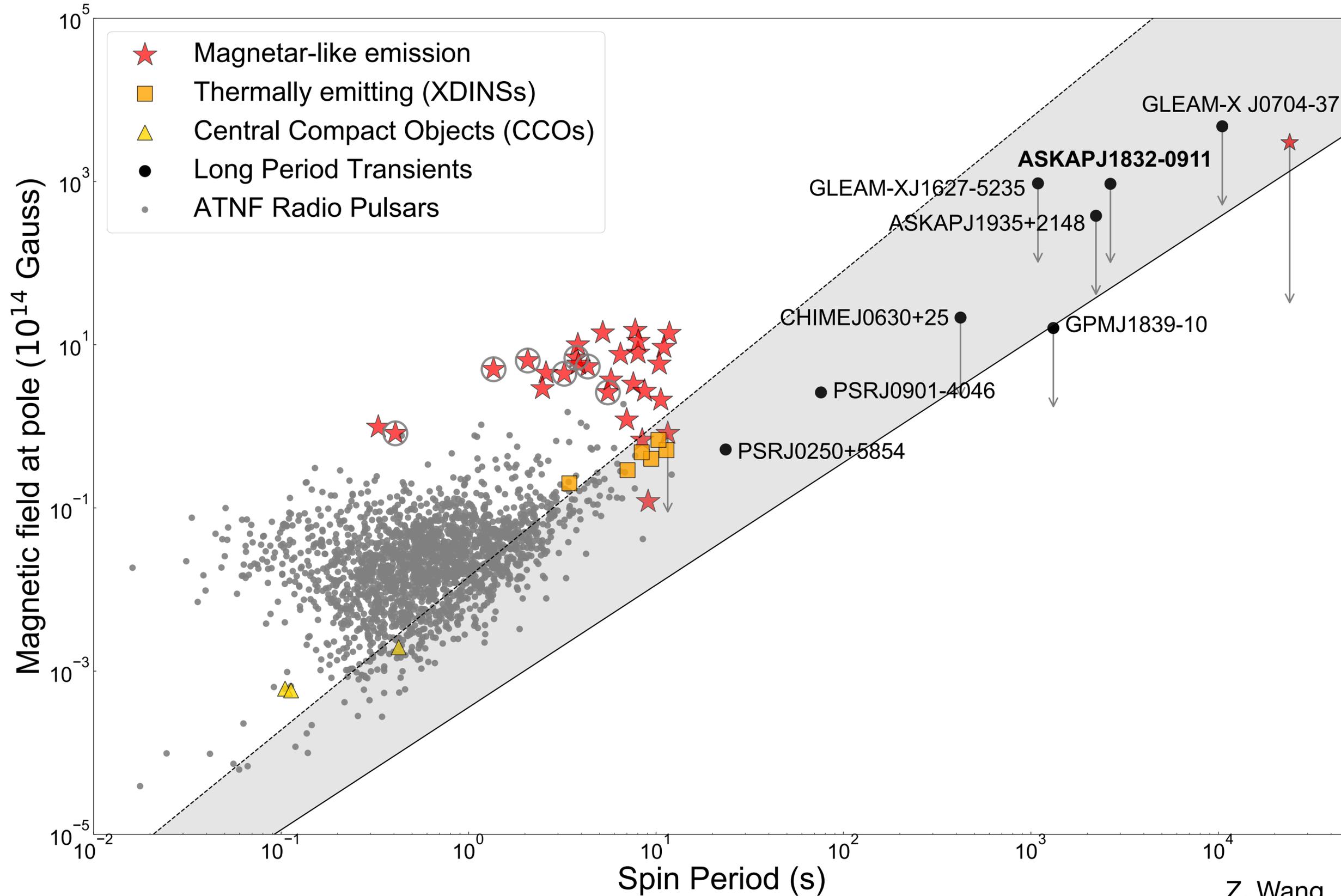


Ridolfi 2018

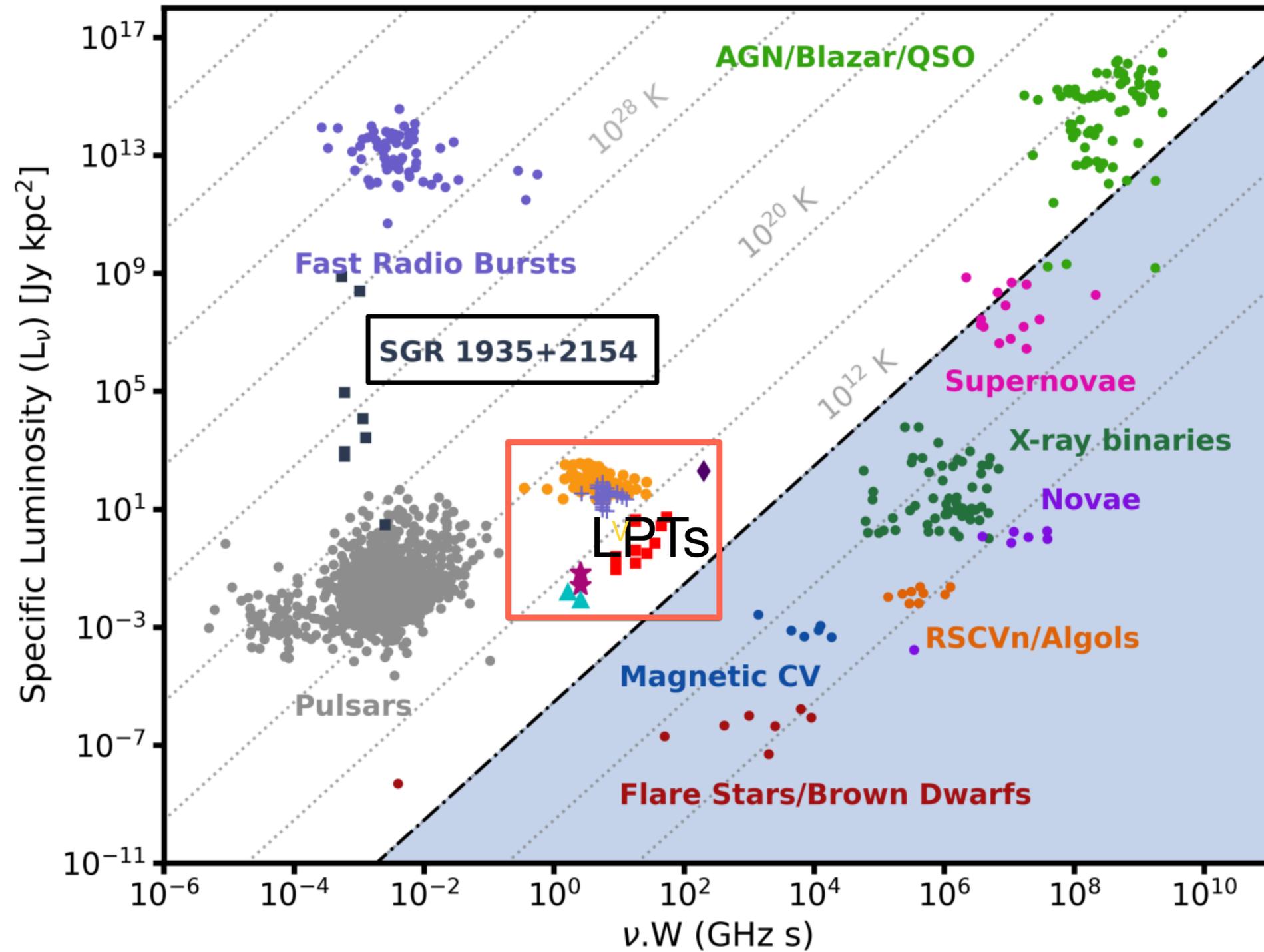


Harding 2013

The Neutron Star Zoo



The Expanding Coherent Radio Transient Phase Space



Motivation: Why are they intriguing?

if they are magnetized neutron stars

- Facilitates FRB generation and escape from inner magnetosphere, potential FRB progenitors (“low-twist” FRB model) (original motivation in **Wadiasingh&Timokhin** 2019; **Wadiasingh+**2020; Beniamini, **Wadiasingh**, Metzger+2020)
- New kind of common highly-magnetized neutron star (Beniamini, **Wadiasingh**, Hare+2023)
 - How do they form and spin down to long periods? Possibility: SN fall-back disks and magnetar winds (Beniamini, **Wadiasingh**, Metzger+2020; Beniamini, **Wadiasingh**, Hare+2023; Ronchi, Rea+ 2022)
- The objects are likely old! Implications for core field retention/evolution and superconductivity (Beniamini,**Wadiasingh**,Hare+2023; Lander 2024; Lander,Gourgouliatos,**Wadiasingh**,Antonopoulou 2024)
- Some pulses extremely bright in radio, exceeding 20-50 Jy, visible at extragalactic distances (N.B. LPRTs could have been discovered in the 1950s)
- Possible systems for test of QED processes, including pair cascades and photon splitting, mode switching and the vacuum resonance if ions are present in the magnetosphere (**Wadiasingh+**2020; Cooper&**Wadiasingh** 2024; Harding; **Wadiasingh**, Baring 2025)
- Interesting for electrospheres
- If involved in BNS/NSBH mergers, potentially greatly improves precursor detectability and EM torques imprints in GW waveform for 3G detectors (e.g., Cooper,Gupta, **Wadiasingh+** 2023; Skiathas, Kalapotharakos, **Wadiasingh+**2025)

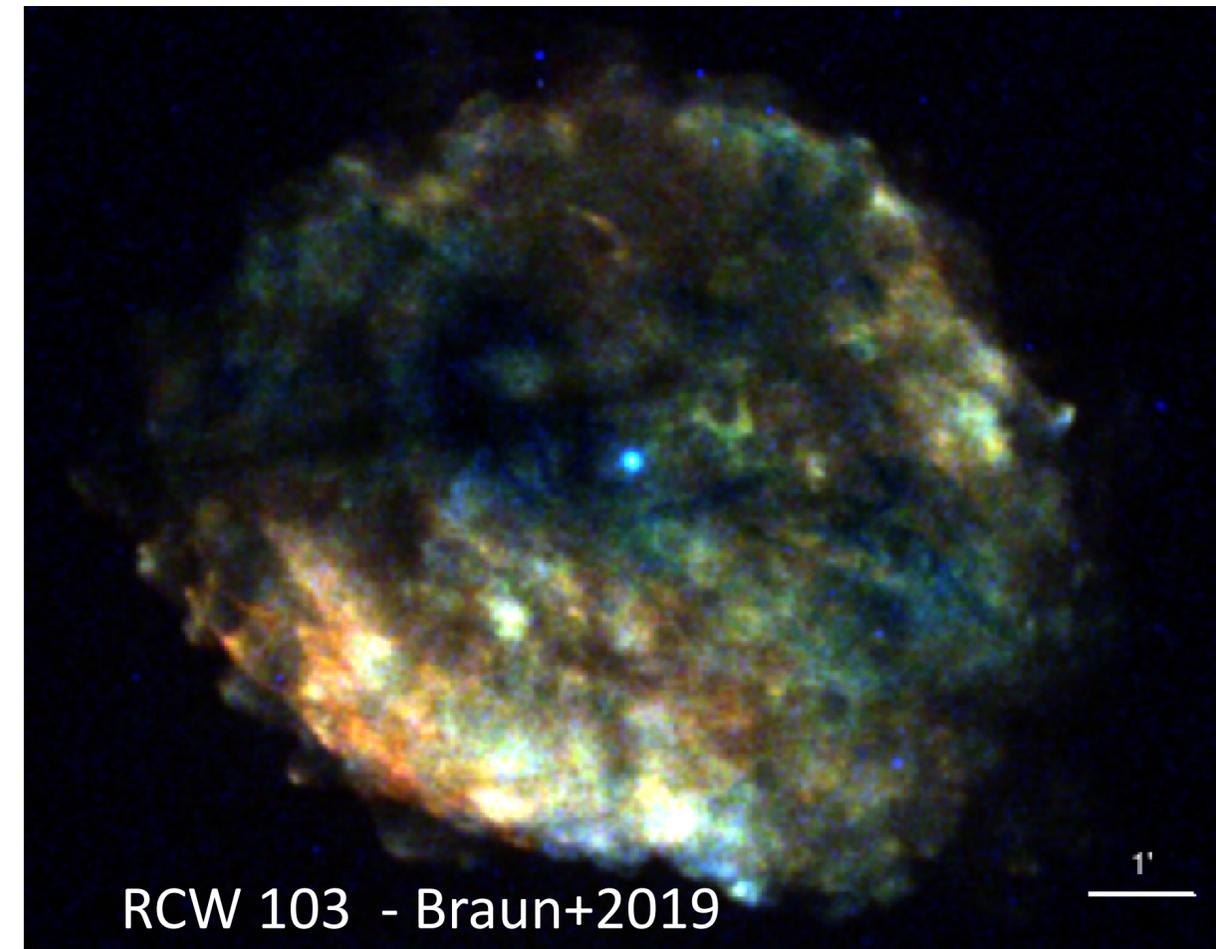
Note 1: 1E 161348–5055 in RCW 103

1E 161348–5055 - The magnetar CCO in RCW 103

Pulsating ($P \sim 6.7$ hr) central compact object in SNR RCW 103:

1. Millisecond duration short X-ray bursts - similar to magnetars
 2. Long-term outbursts and non-thermal hard X-ray emission
 3. Proper motion ~ 170 km/s from *Chandra* imaging – Wide binary would have been disrupted
 4. Companion hotter than M7 ruled out by HST observations – close binary should have been detected
- } Magnetar-like phenomenology

Credit: De Luca et al. 06, 08, Esposito et al. 11, D’Ai et al. 16, Rea et al. 16, Tendulkar et al. 17, Borghese et al. 18, Braun et al. 2019



Note 1 - 1E 161348-5055 in RCW 109

1E 161348-5055

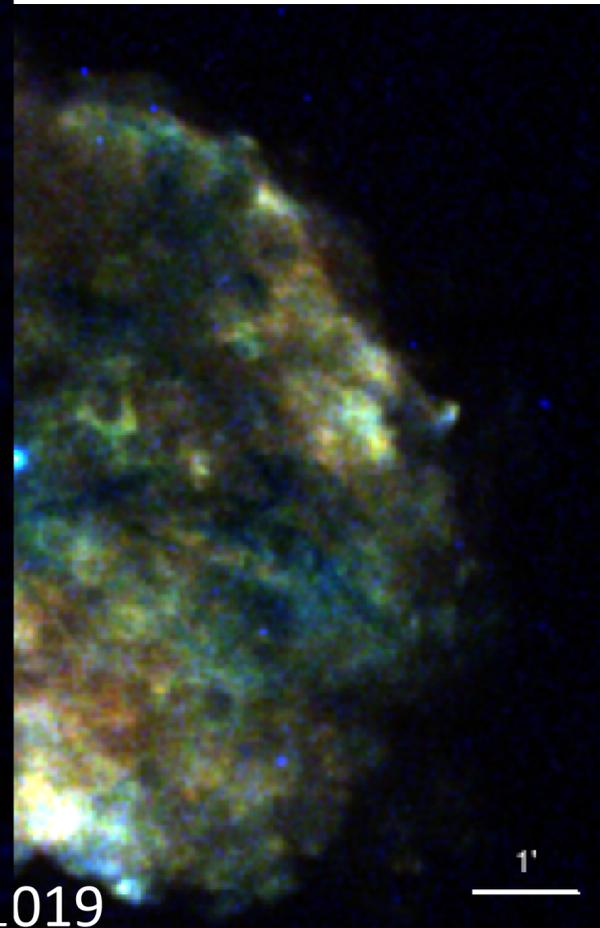
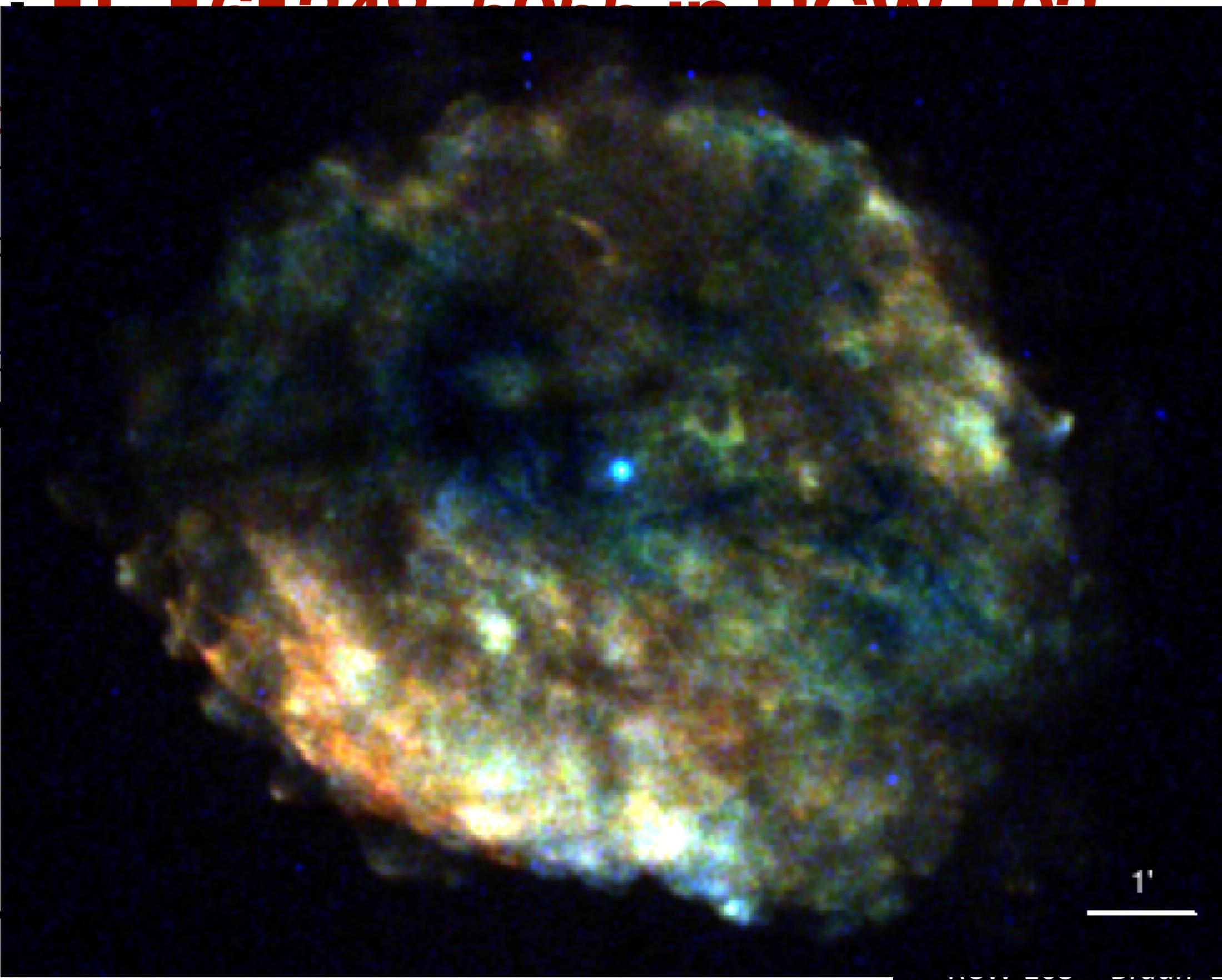
Pulsating (P ~ 6.7 ms)

1. Millisecond pulsar
2. Long-term orbital period
3. Proper motion
4. Companion star

phenology

noted

been detected



Credit: D
et al. 16,

2019

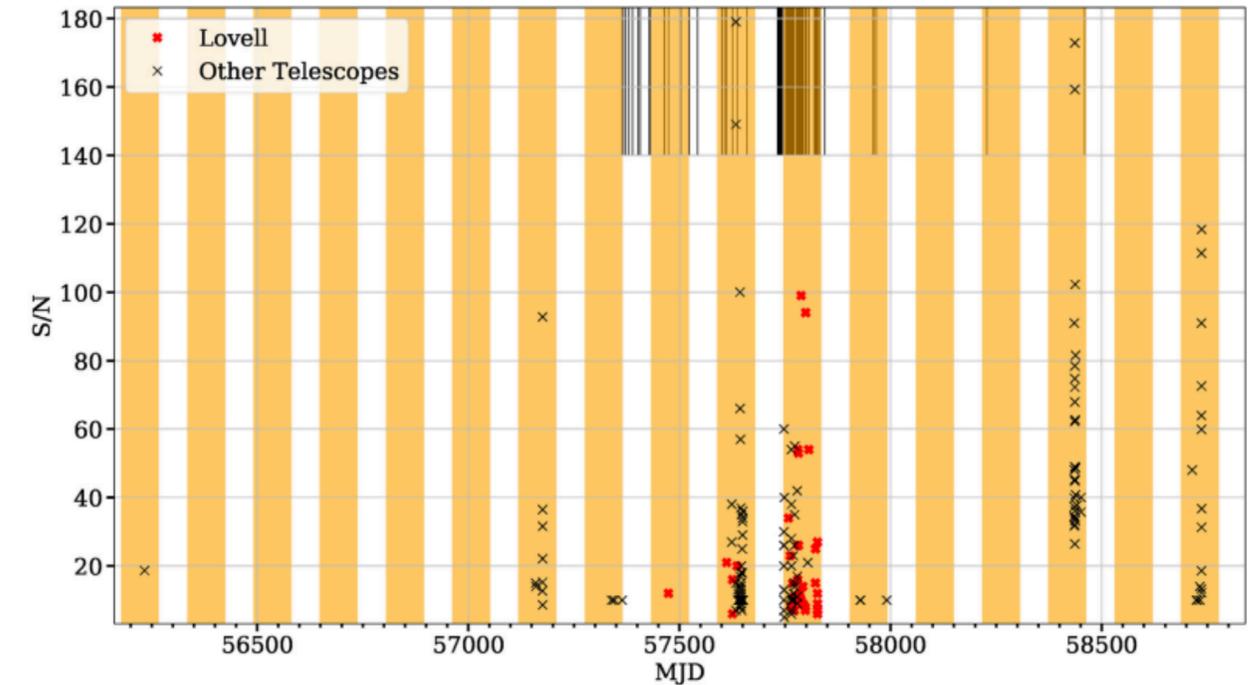
Note 2: Connection to FRBs

Magnetars are the prime candidate FRB sources, possible shared conditions for bright coherent emission

Secure long-term periodicity (16 days and 157 days) reported for two repeating FRBs - third one of ~120 days tentatively detected

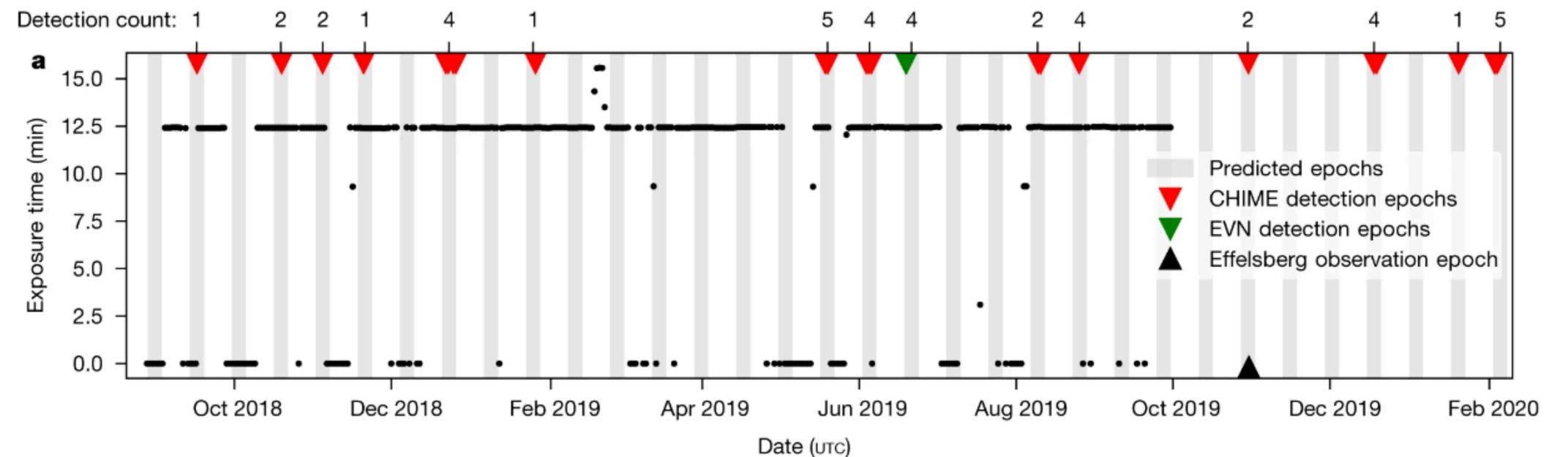
Precession, binarity, obscuration, or **long periods**

Long periods facilitate escape of FRBs from the inner magnetosphere



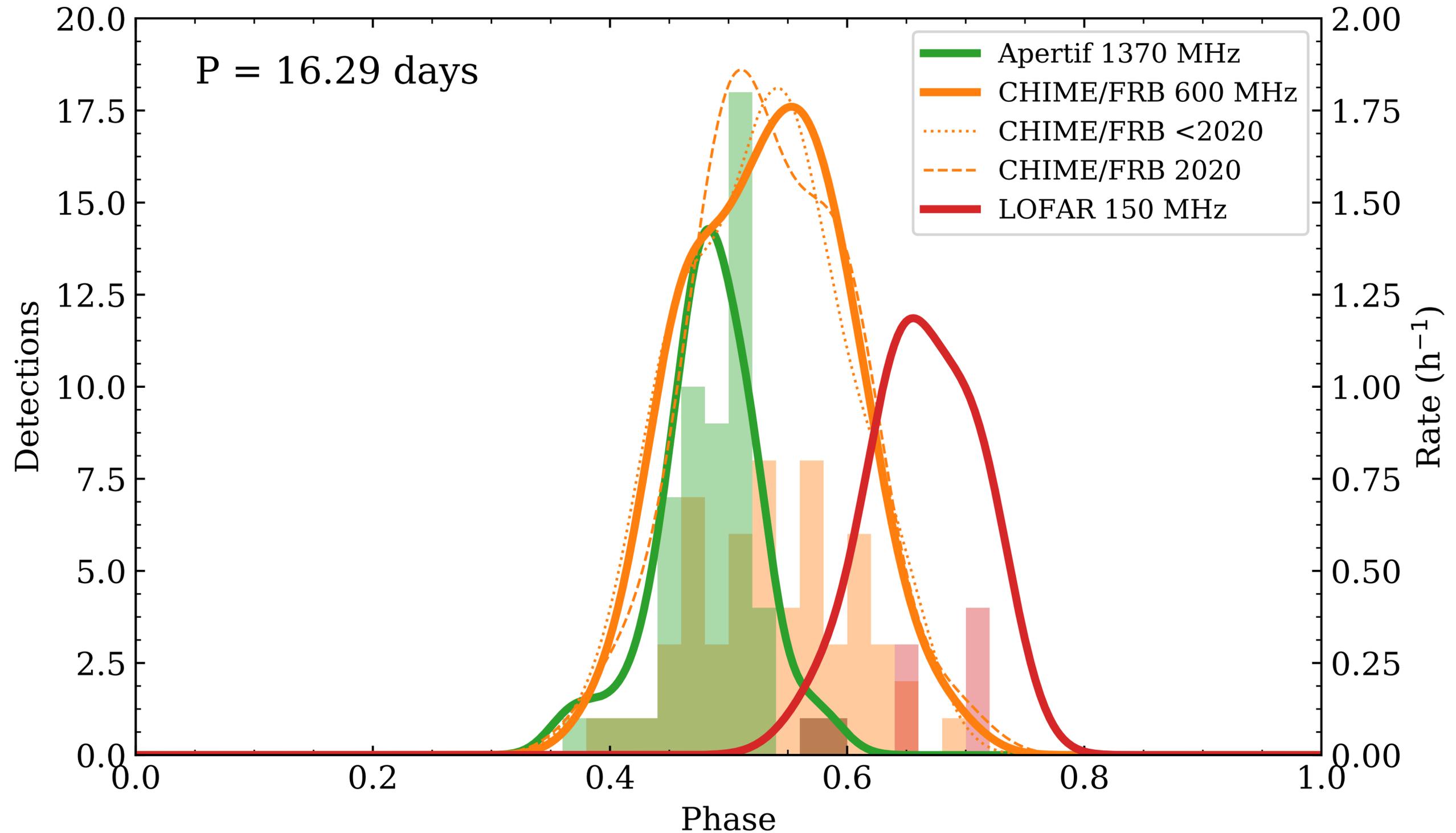
Rajwade et al. (2020)

e.g. Wadiasingh & Timokhin 2019
 Wadiasingh et al. 2020,
 Beniamini, Wadiasingh et al., 2020
 Beniamini, Wadiasingh et al. 2023
 Cooper & Wadiasingh 2024



CHIME/FRB Collaboration (2020)

Shrouding in a binary seems ruled out in FRB 180916



Pastor-Marazuela+ 2021, Nature, 2012.08348

Some observational highlights

Observations of long period transients

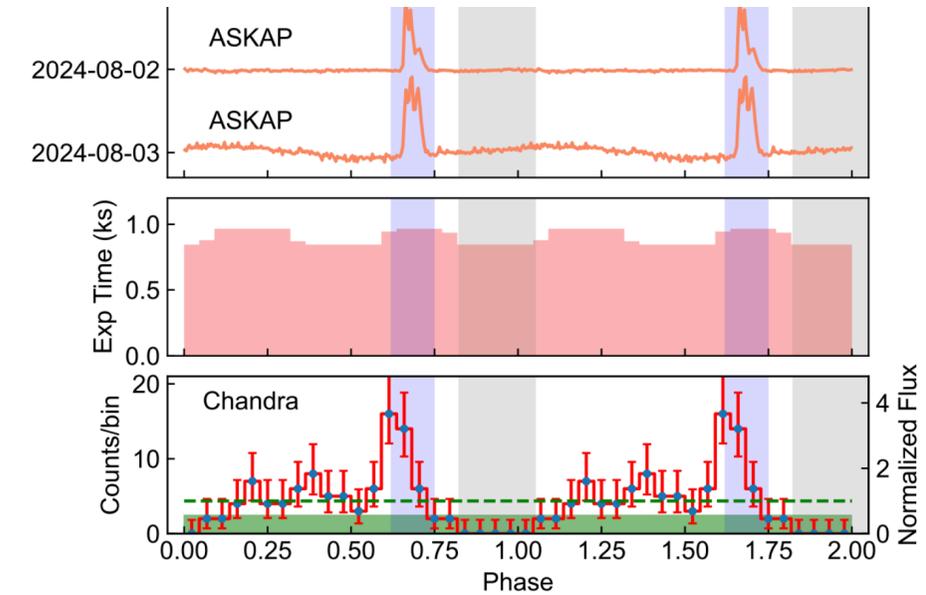
ASKAP J1935+2148: 54 minute period transient with 3 distinct emission states

CHIME J0630+25: 7 minute period transient, DM distance 170 ± 80 pc

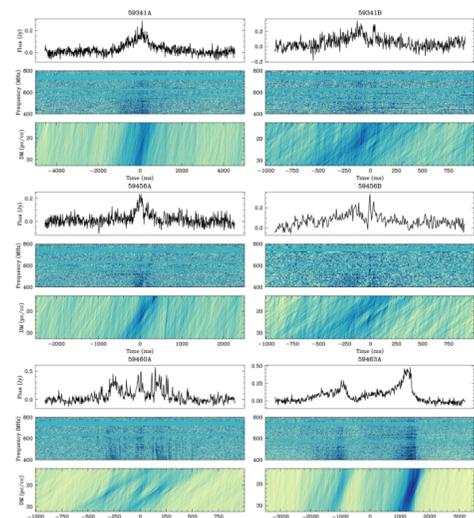
GLEAM-X J0704-37: 2.9hr period *binary* source with high Galactic latitude

ASKAP J1839-07: 6.5hr period source with interpulses

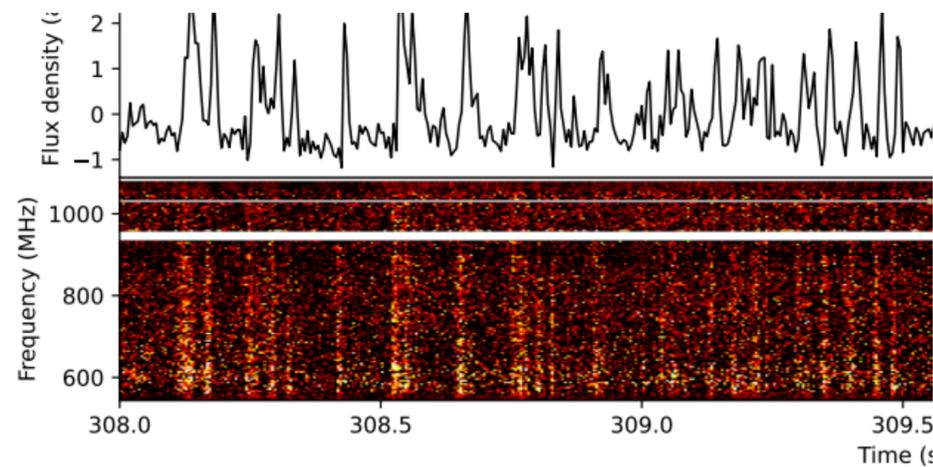
ASKAP J1832-0911: 44 min period with X-ray pulsations(!)



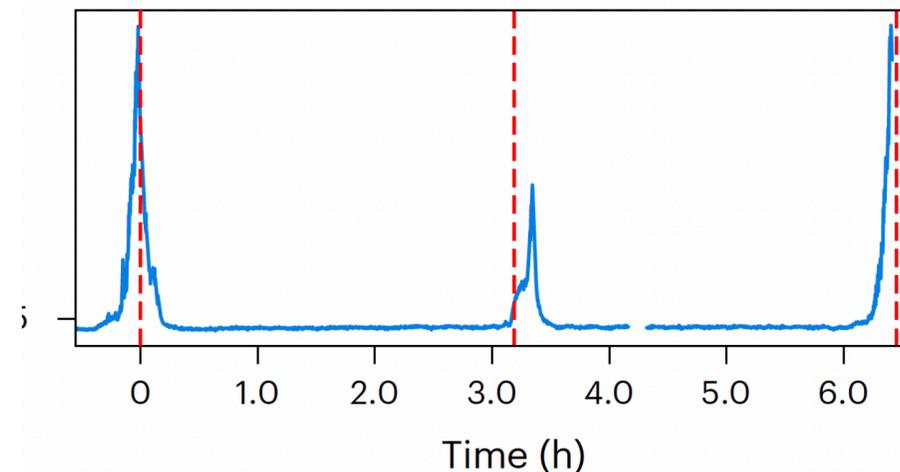
Wang et al., 2025 (ASKAP J1832)



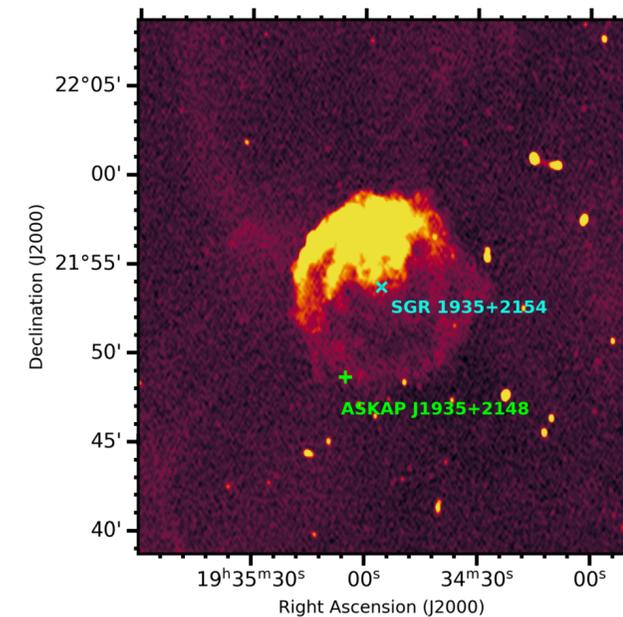
Dong et al., 2024 (CHIME)



Hurley-Walker et al., 2025 (GLEAM-X)



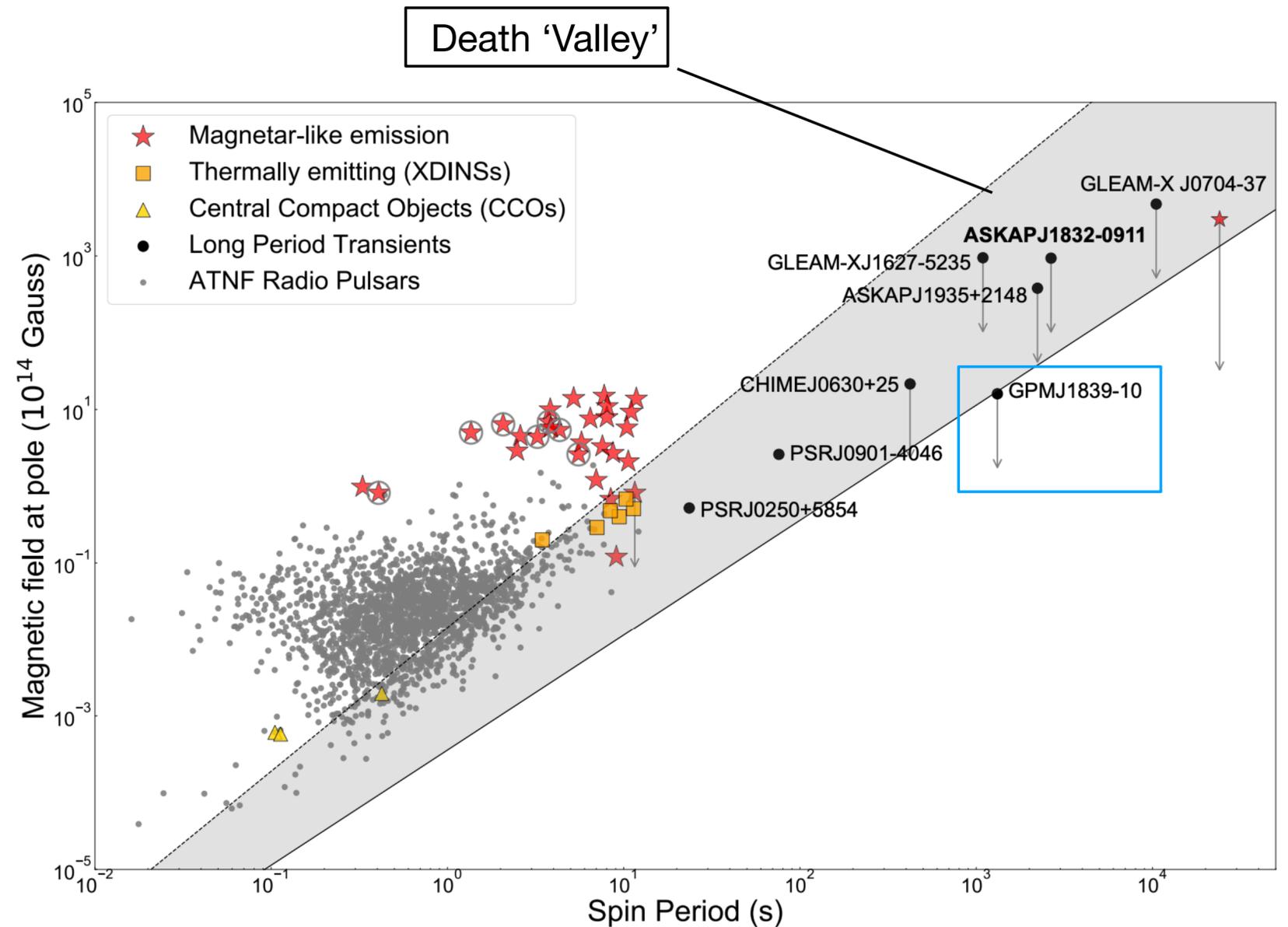
Lee et al., 2025 (ASKAP J1839)



Caleb et al., 2024 (ASKAP J1935)

Are they magnetars?

Source	GLEAM-X J1627	GPM J1839-10
P [min]	18.18	21
\dot{P}	$< 1.2 \times 10^{-9}$	$< 4.6 \times 10^{-13}$
Pulse [s]	30-60	30-300
Distance [kpc]	1.3 ± 0.5	5.7 ± 2.9
$F_{\nu, \text{radio}}$ [Jy]	5-40	0.1-10
L_{radio} [erg/s]	$\approx 10^{28-31}$	$\approx 10^{28}$
$L_{\text{spin-down}}$ [erg/s]	$\lesssim 1.2 \times 10^{28}$	$\lesssim 10^{25}$
$L_{\text{X}, 0.3-10\text{keV}}$ [erg/s]	$< 10^{32}$	$< 1.5 \times 10^{32}$
Duty cycle	≈ 2 months	$\gtrsim 33$ years

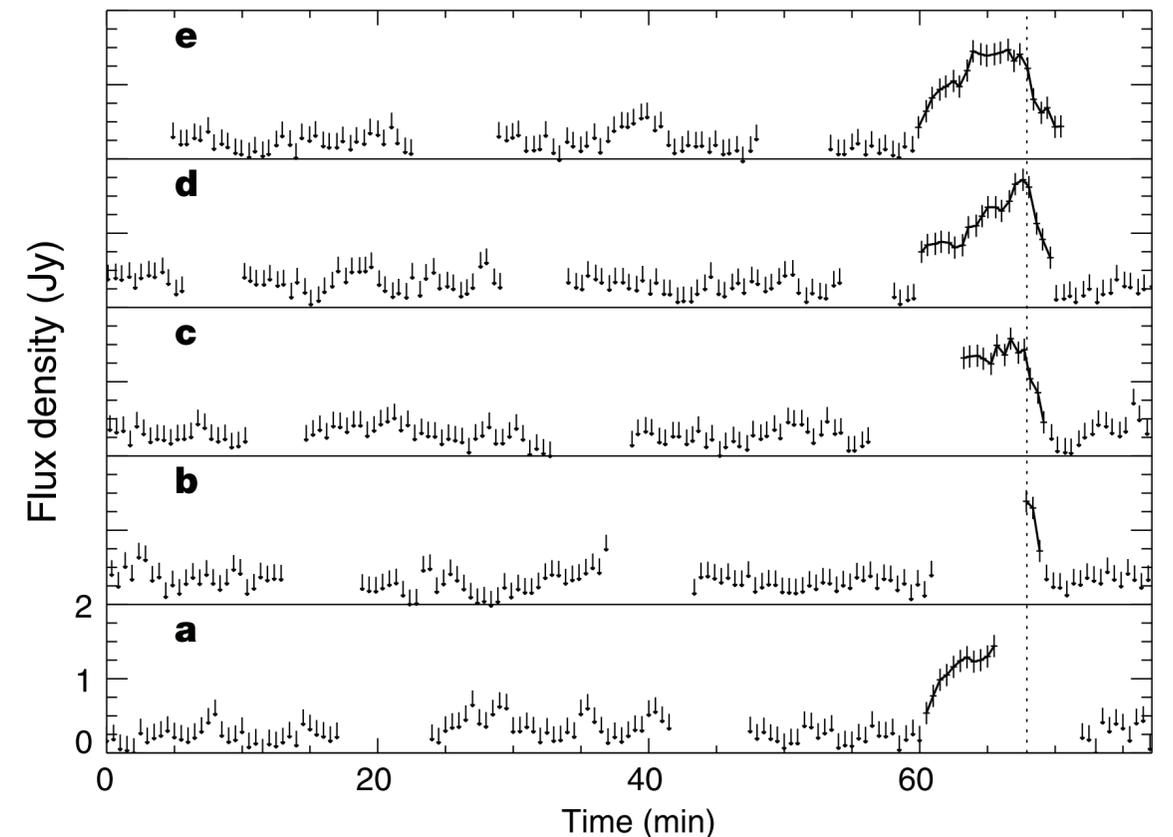


Galactic ULPM candidates - GCRT J1745–3009

GCRT J1745–3009

The Galactic “burper”. A $P \sim 77$ minute source discovered serendipitously by VLA

1. 10 minute wide “pulses”
2. $T_{\text{brightness}} \gg 10^{12}$ K for $D > 70$ pc
3. Optical observations rule out M type / brown dwarf nearby counterpart
4. If period is spin – cannot be rotation powered – suggestive of a magnetar origin



Credit: Hyman et al. 05, Kaplan et al. 08, Spreeuw et al. 09

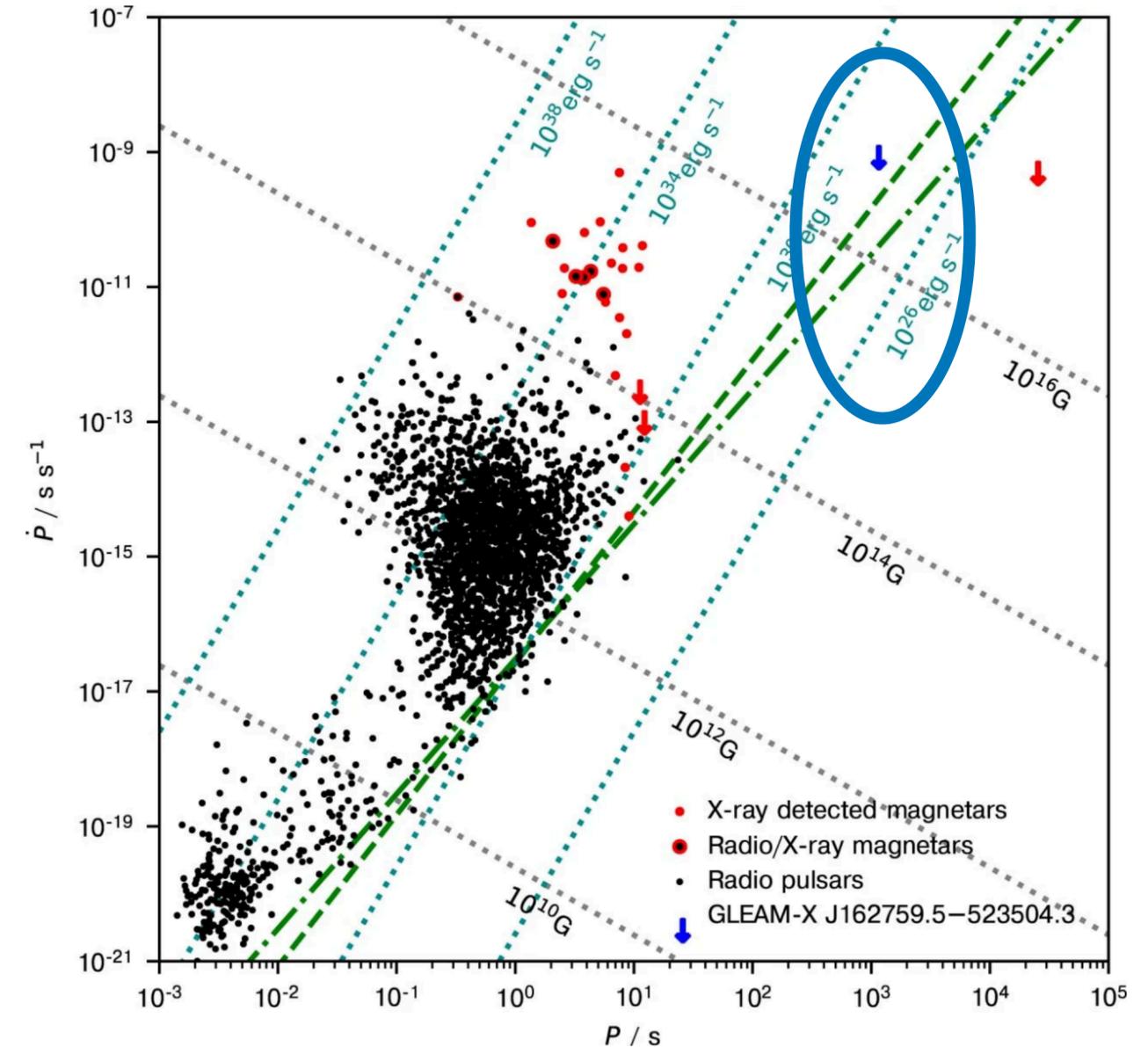
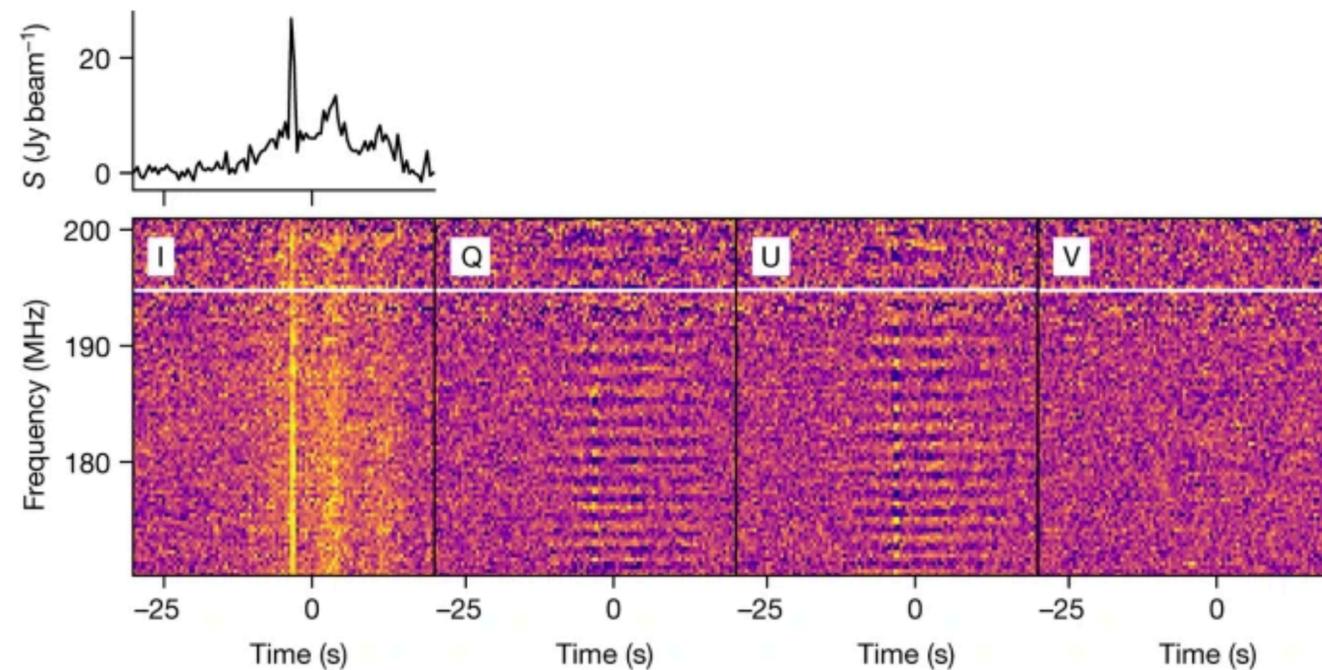
Galactic ULPM candidates - **GLEAM-X J162759.5–523504.3**

GLEAM-X J162759.5–523504.3

$P \sim 1091$ sec, $\dot{P} < 10^{-9}$ radio transient

1. Close to 100% linear polarization, up to 40 Jy pulses
2. Rapid (~ 0.5 s) variability suggesting compact object with brightness temperature $> 10^{16}$ K
3. Cannot be a rotation powered NS
4. 2% duty cycle
5. Beyond pulsar death-line for standard pulsar field strength
6. WD can largely be ruled out (Beniamini, Wadiasingh, Hare+2023)
7. No multi-wavelength counterpart - most binary companions ruled out

Rea et al. (2022) & Lyman et al. (2025)



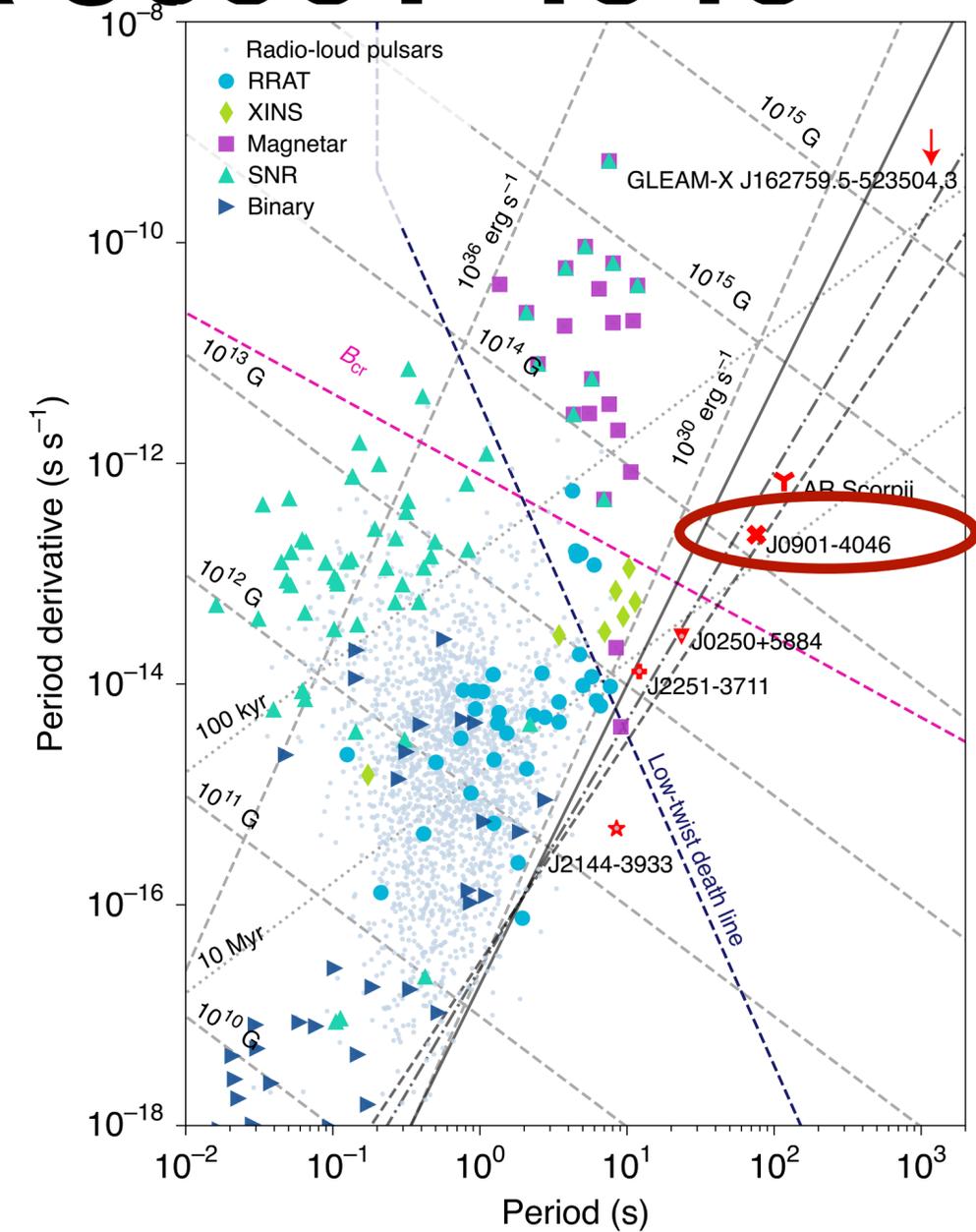
Credit: Hurley Walker et al. 2022

Galactic ULPM candidates - PSR J0901-4046

PSR J0901-4046 (formerly MTP0013)

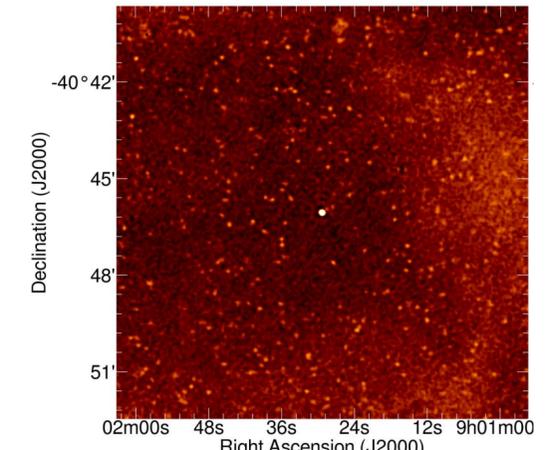
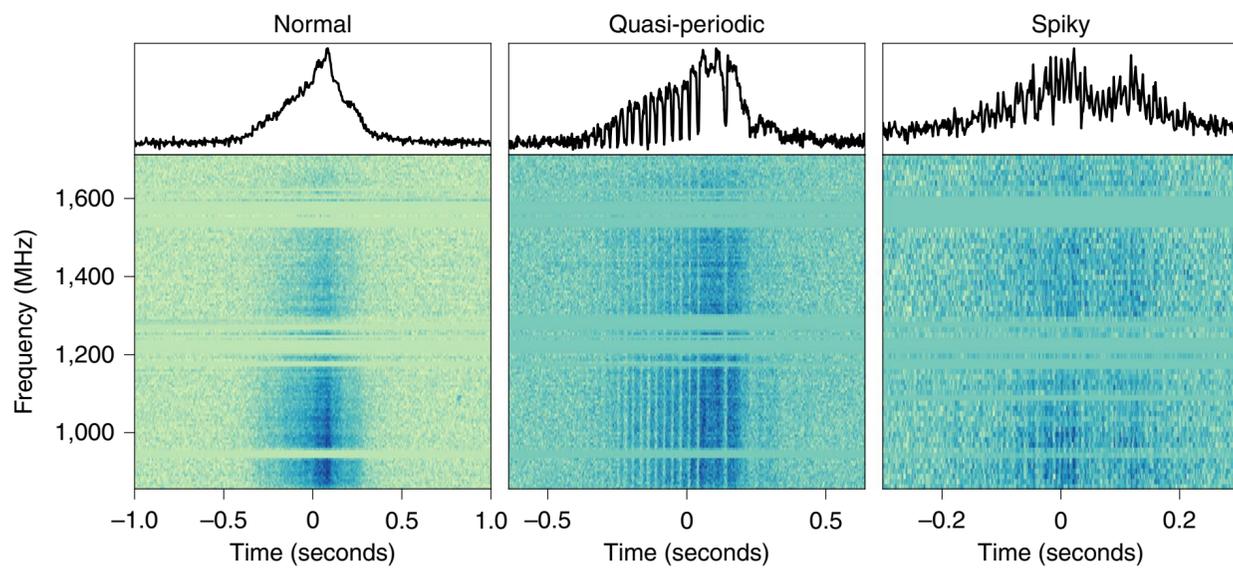
$P \sim 76$ s pulsar, well-measured $\dot{P} \sim 2 \times 10^{-13}$ $\rightarrow B_{\text{pole}} \sim 2 \times 10^{14}$ G, age > 5 Myr

1. Pulsar radio characteristics: high polarization fraction, PPA swings, variability in single pulses of flux and polarization
2. Very stable in timing — unusually stable for a magnetar
3. Only 328 pc away (YWM16) — implies many more exist
4. $L_x < 10^{30.5}$ erg/s
5. harmonic spaced QPOs at O(10) Hz— consistent with the existence of NS crust, unlikely to be magnetospheric Alfvénic modes



Discovery of a radio-emitting neutron star with an ultra-long spin period of 76 s

Manisha Caleb^{1,2,3,14}, Ian Heywood^{4,5,6,14}, Kaustubh Rajwade^{1,7}, Mateusz Malenta¹, Benjamin Willem Stappers^{1,4}, Ewan Barr⁸, Weiwei Chen⁸, Vincent Morello¹, Sotiris Sanidas⁹, Jakob van den Eijnden⁴, Michael Kramer^{1,8}, David Buckley^{9,10,11}, Jaco Brink^{9,10}, Sara Elisa Motta¹², Patrick Woudt¹⁰, Patrick Weltevrede¹, Fabian Jankowski¹, Mayuresh Surnis¹, Sarah Buchner⁶, Mechel Christiaan Bezuidenhout¹, Laura Nicole Driessen^{1,13} and Rob Fender⁴



Credit: Caleb et al. 2022

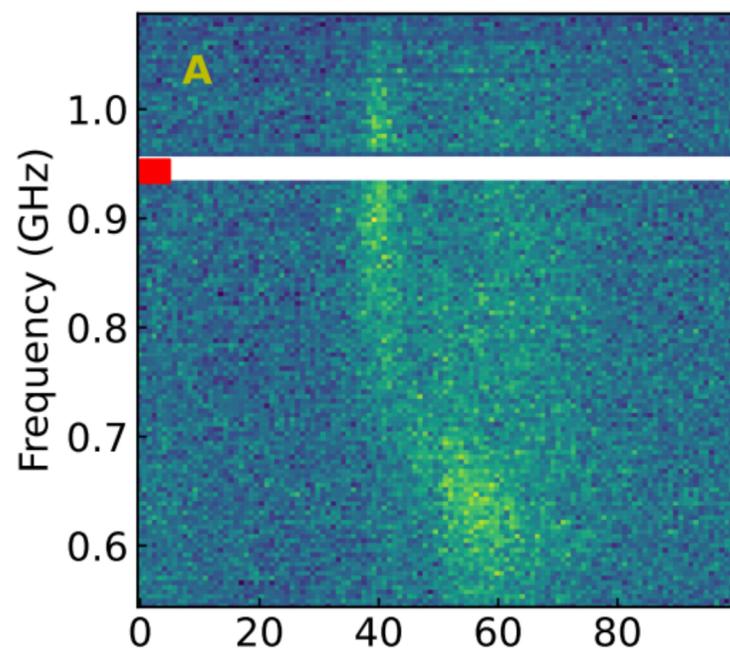
Observations of long period transients - GPM J1839-10

GPM J1839-10, with pulses dating back decades

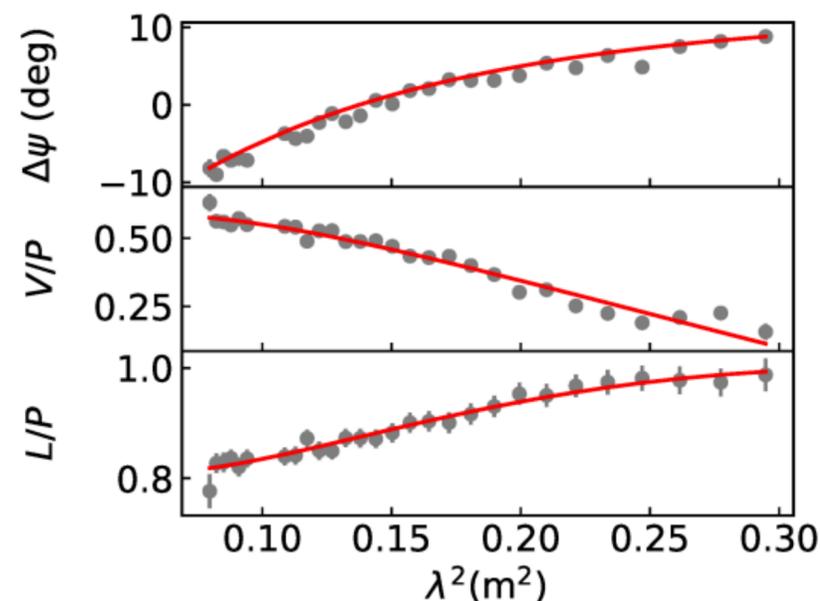
21 minute period, 30-300s pulses with rich substructure, FRB-like drifting and linear-to-circular polarisation conversion

Strong spin period derivative $< 3 \times 10^{-13}$ s/s

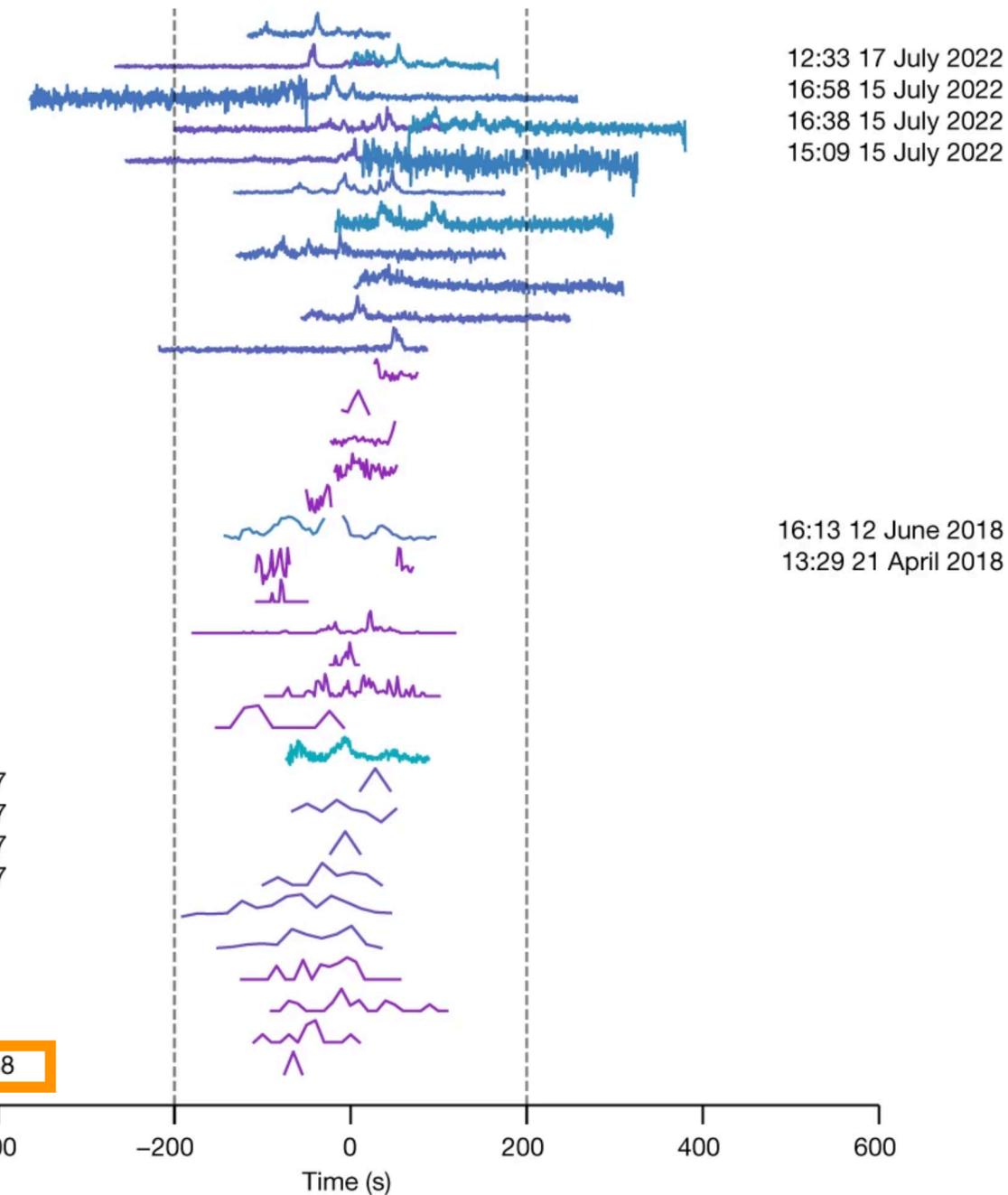
No multi-wavelength counterpart



Men et al., 2025

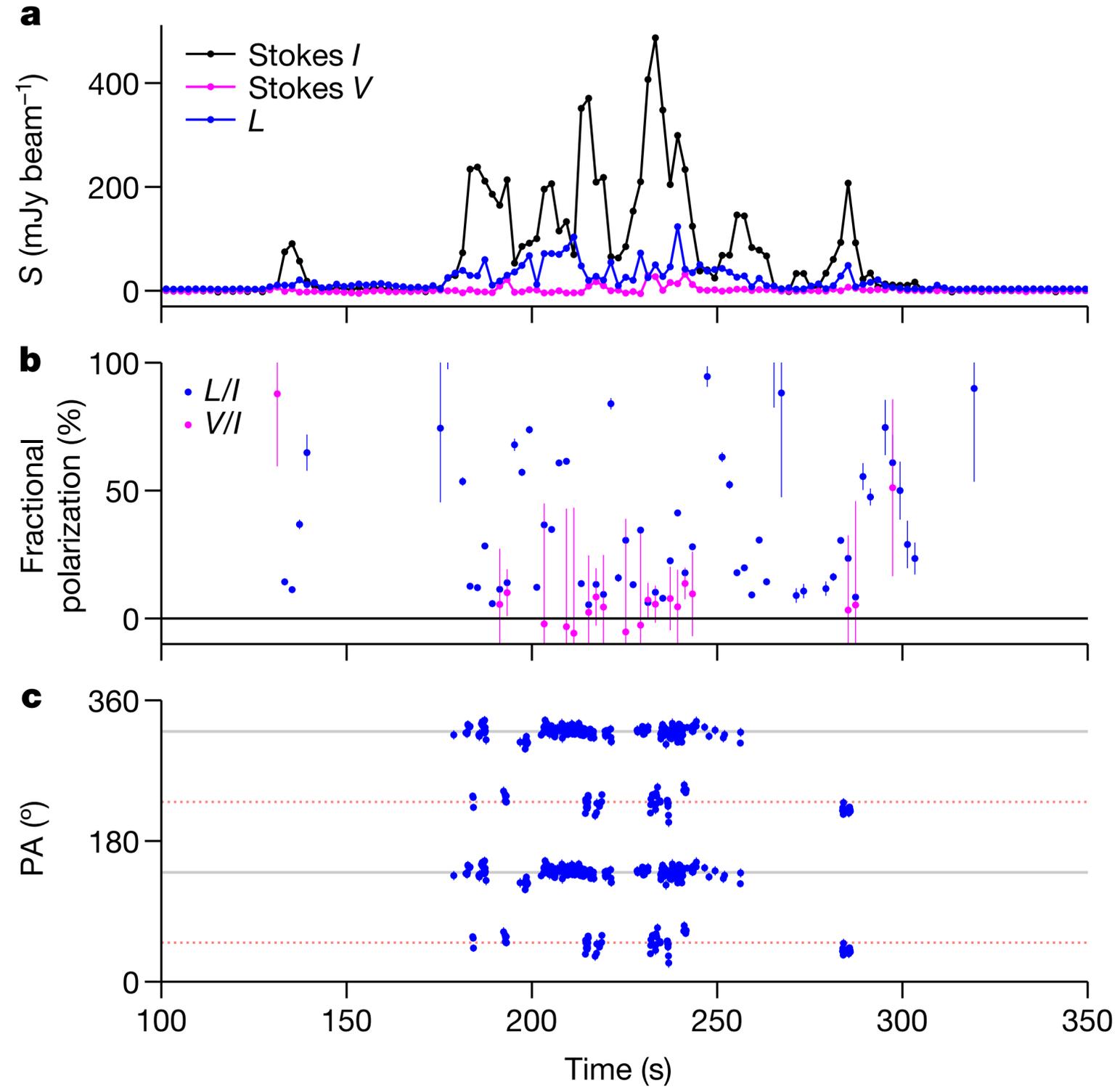


- 12:53 17 July 2022
- 12:28 17 July 2022
- 16:53 15 July 2022
- 16:33 15 July 2022
- 15:04 15 July 2022
- 14:44 15 July 2022
- 14:25 15 July 2022
- 14:56 14 July 2022
- 14:36 14 July 2022
- 16:17 30 June 2022
- 15:52 30 June 2022
- 11:53 19 May 2022
- 08:03 18 May 2022
- 08:59 6 May 2022
- 13:09 19 April 2022
- 13:37 24 February 2021
- 16:11 12 June 2018
- 13:27 21 April 2018
- 06:08 17 January 2015
- 03:33 17 January 2015
- 09:41 3 January 2015
- 07:06 3 January 2015
- 22:05 8 May 2014
- 13:26 27 August 2013
- 11:58 14 December 2007
- 13:27 29 November 2007
- 13:06 29 November 2007
- 12:54 17 November 2007
- 21:16 12 May 2002
- 19:05 12 May 2002
- 02:20 31 August 2001
- 13:21 28 April 2001
- 11:30 28 April 2001
- 22:12 23 September 1988

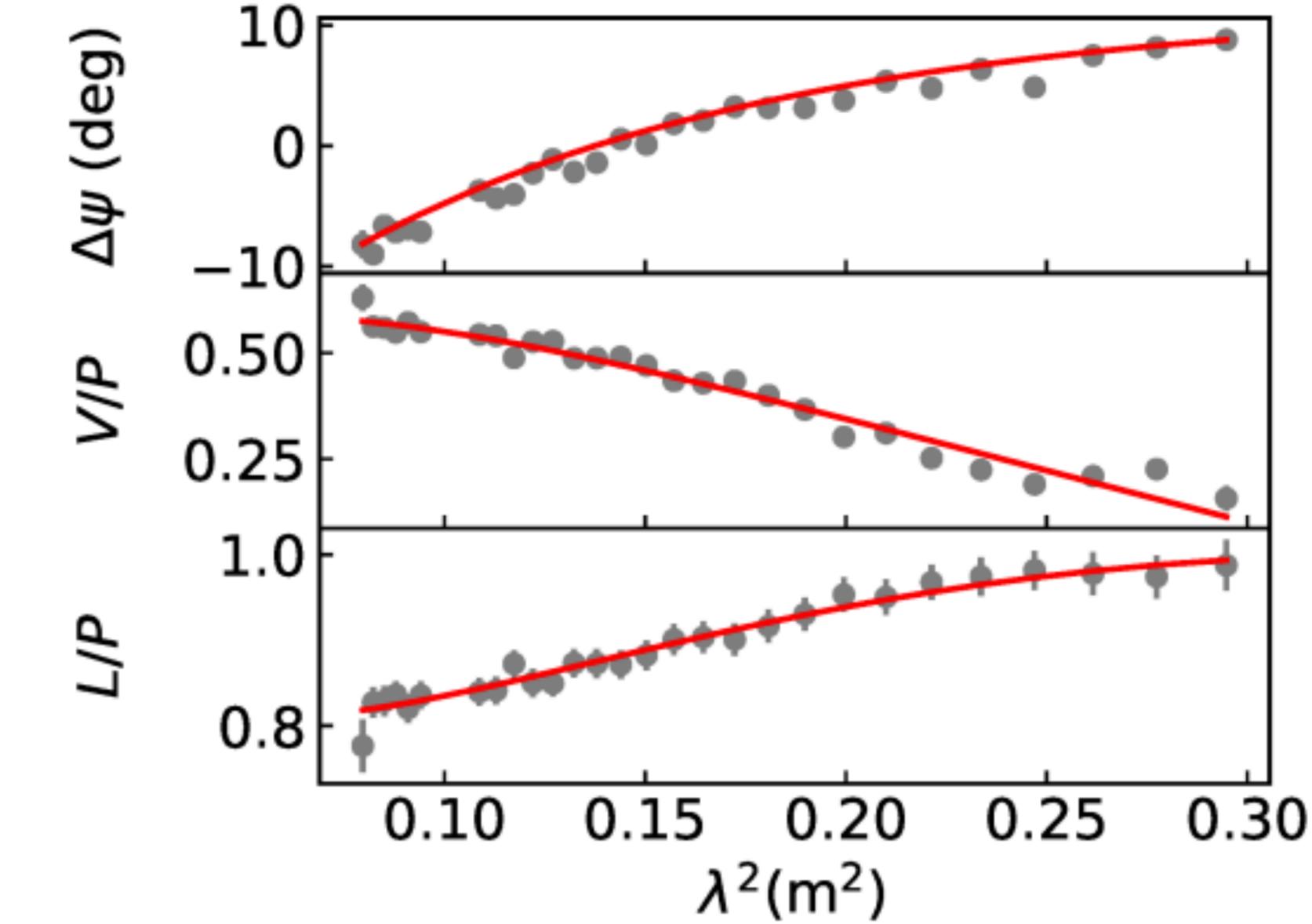
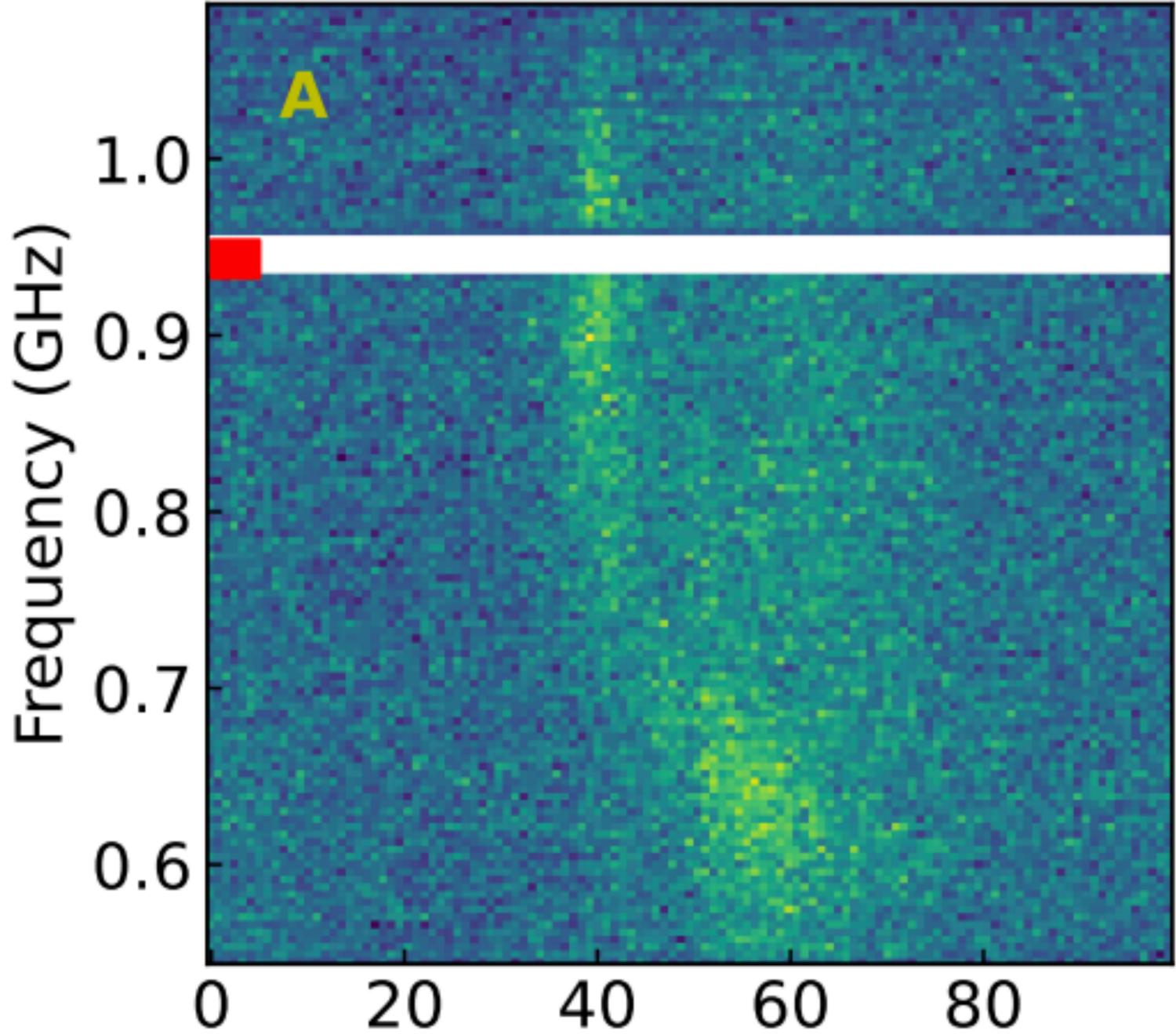


Hurley-Walker et al., 2023

Observations of long period transients - GPM J1839-10

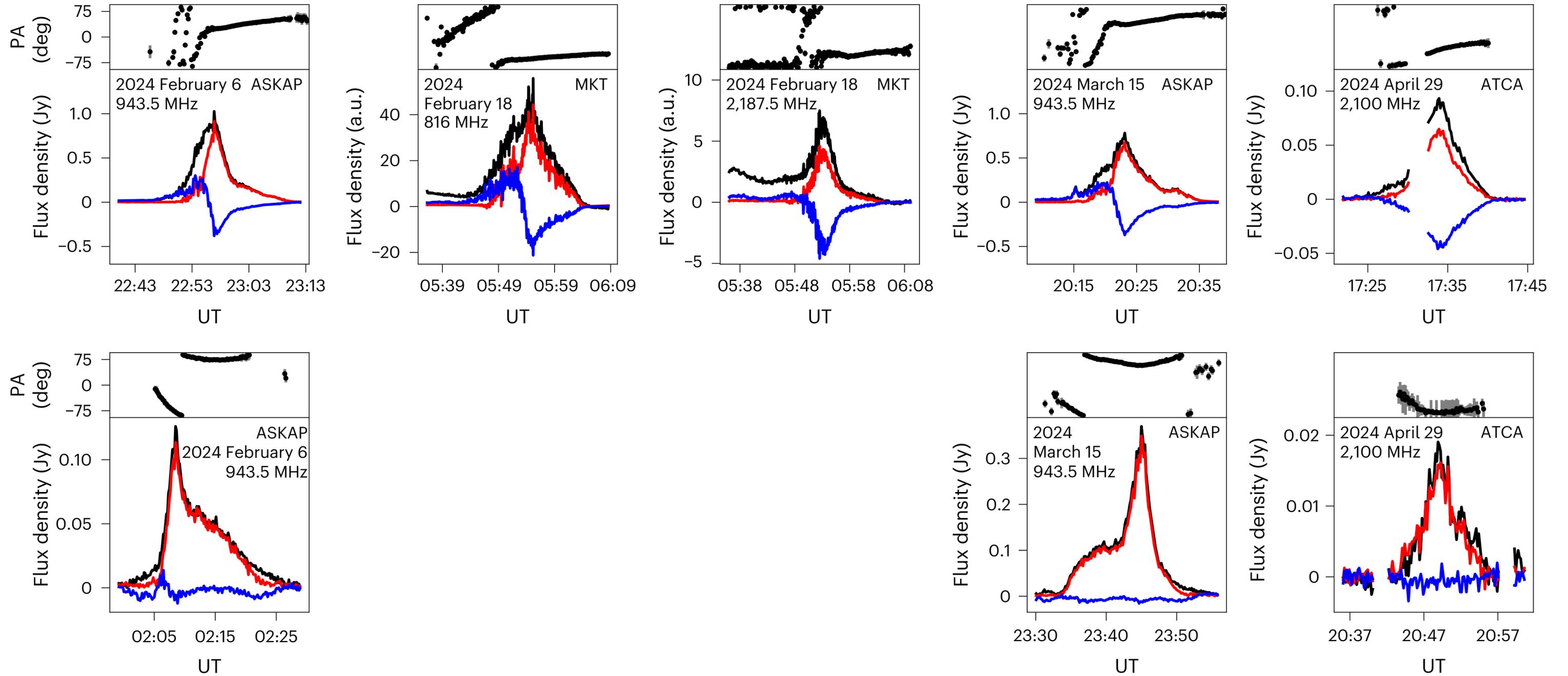


Observations of long period transients - GPM J1839-10



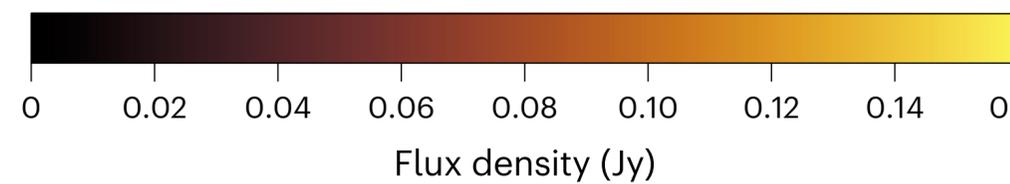
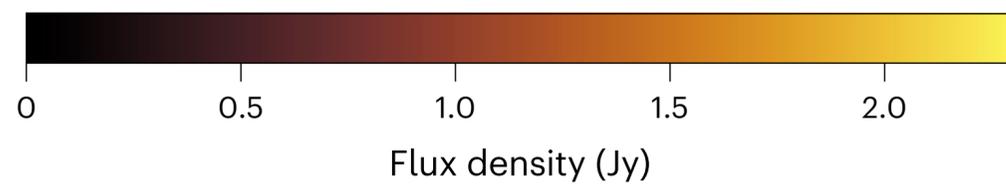
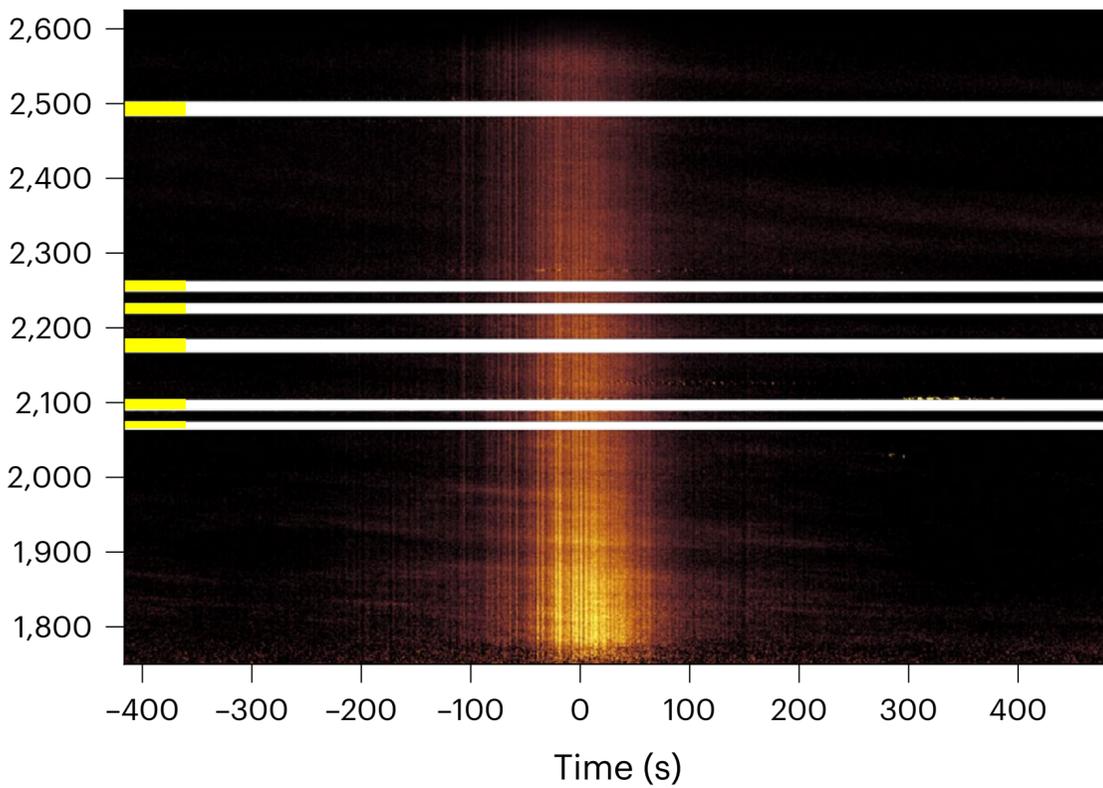
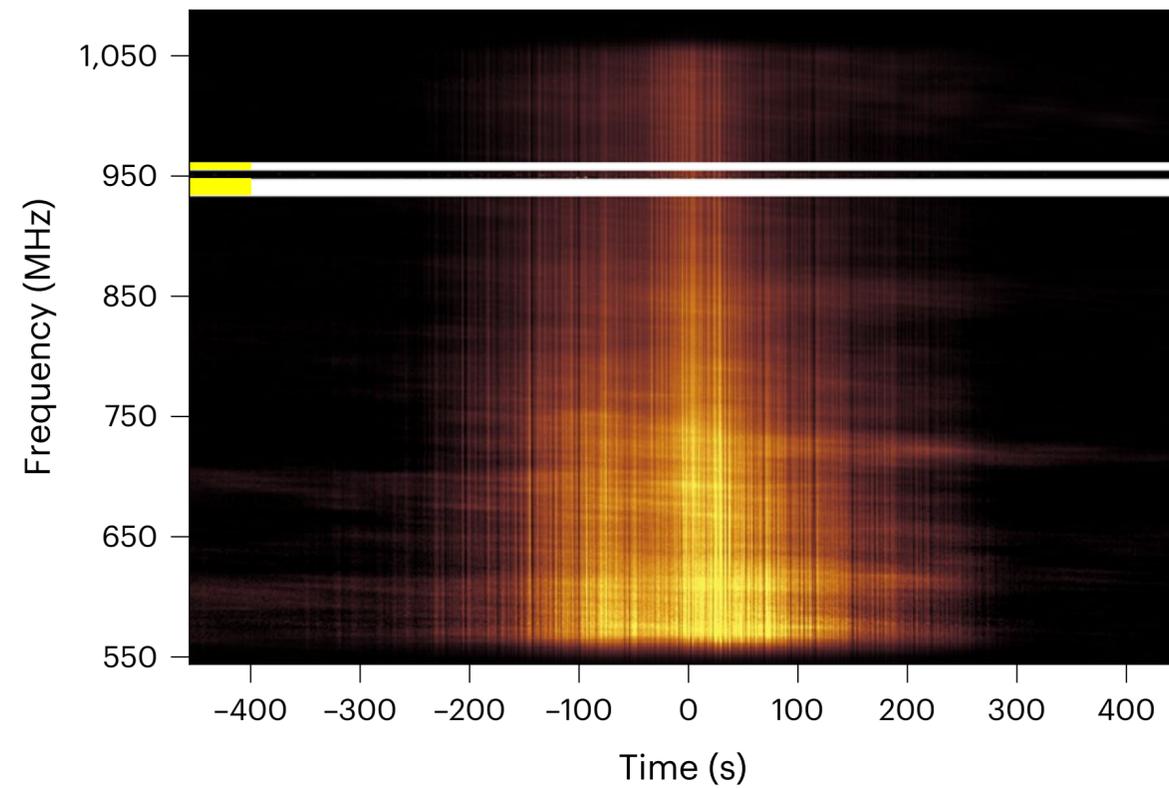
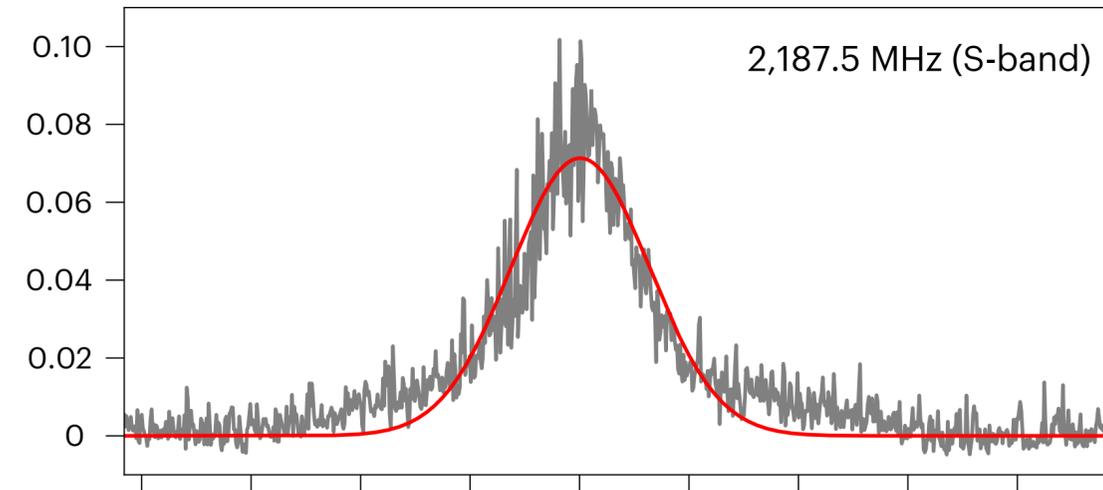
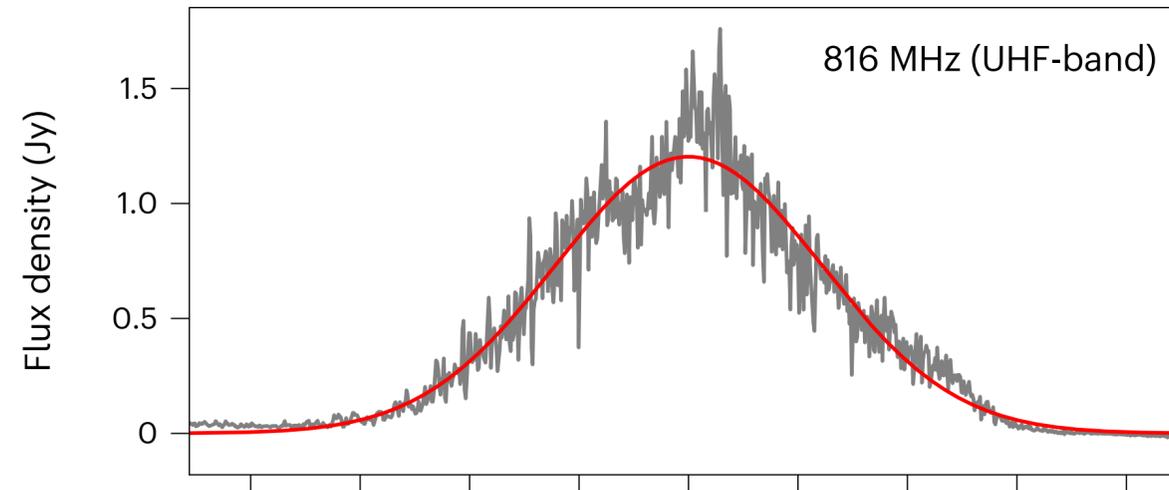
Men et al., 2025

A 6.45 hr LPRT with an interpulse - ASKAP J1839-07



No apparent optical / binary counterpart

A 6.45 hr LPRT with an interpulse - ASKAP J1839-07



No apparent optical / binary counterpart

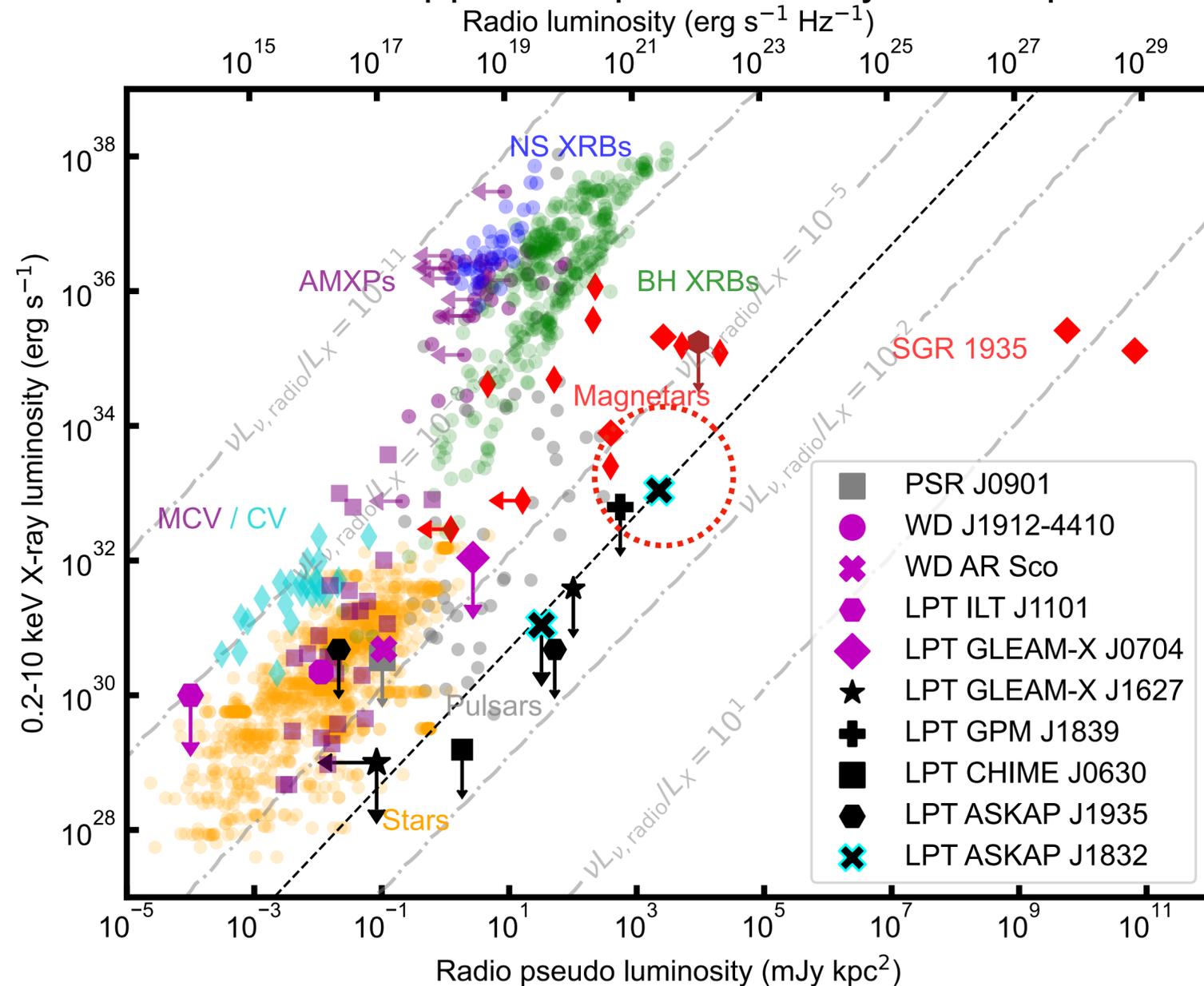
Lee, Caleb... **Wadiasingh** +2025

Observations of long period transients - ASKAP J1832-0911

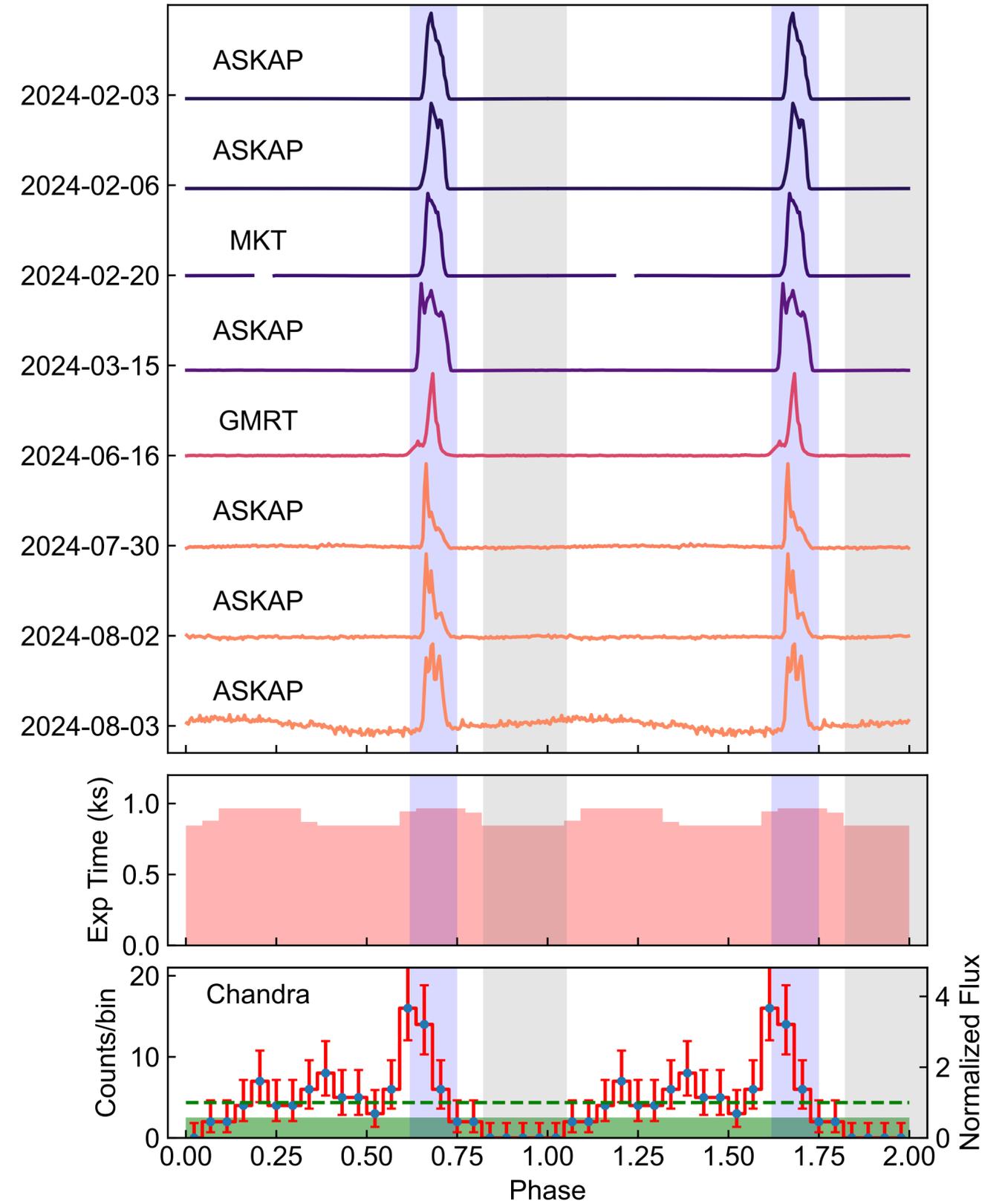
10^{33} erg/s radio and X-rays - highly efficient radio

Some radio pulses 20 Jy!

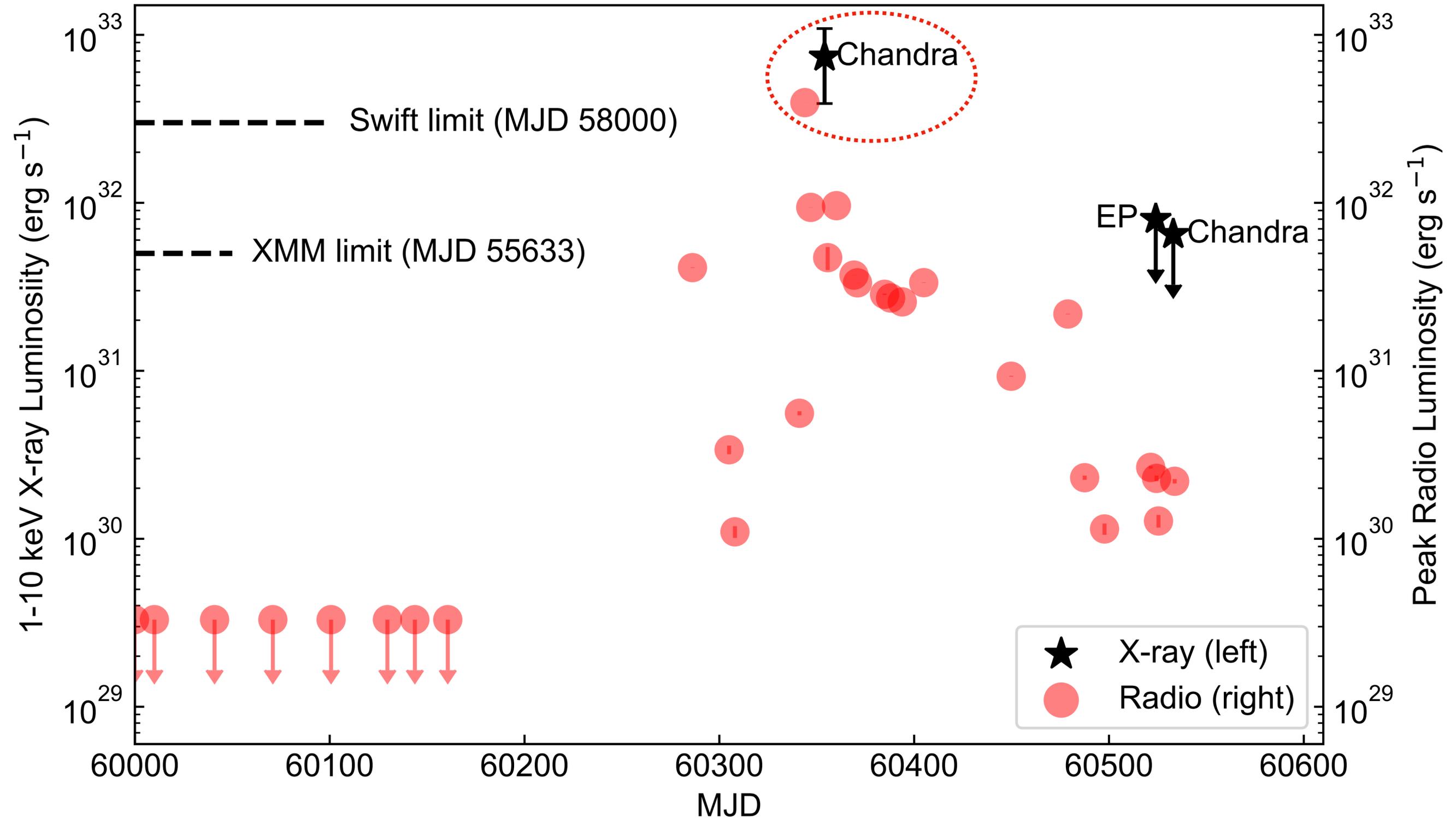
No apparent optical / binary counterpart



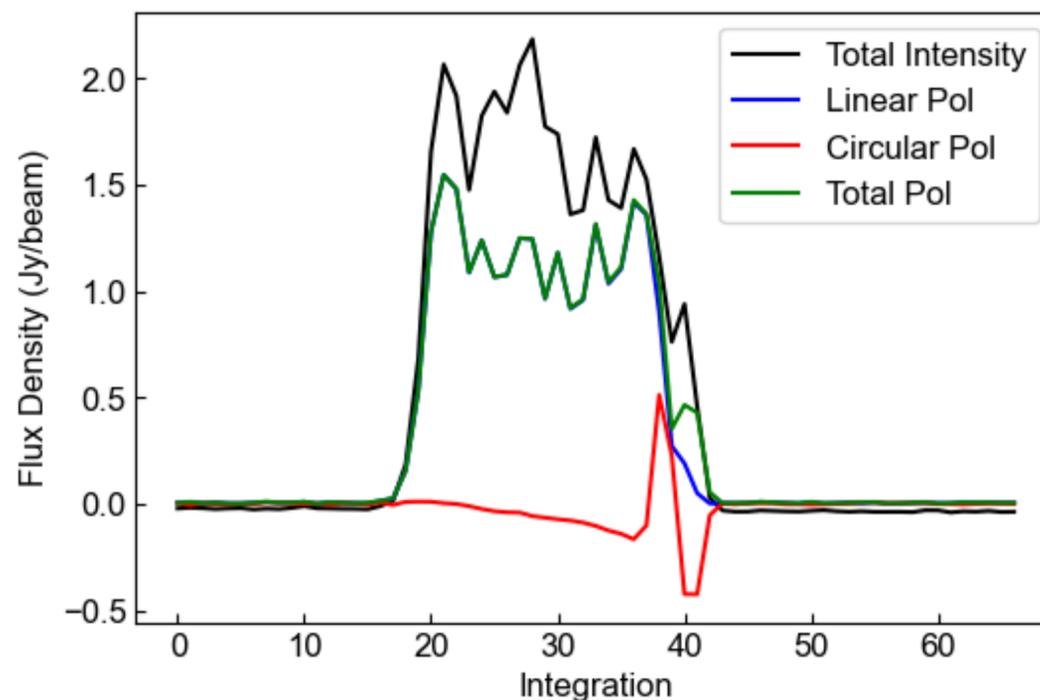
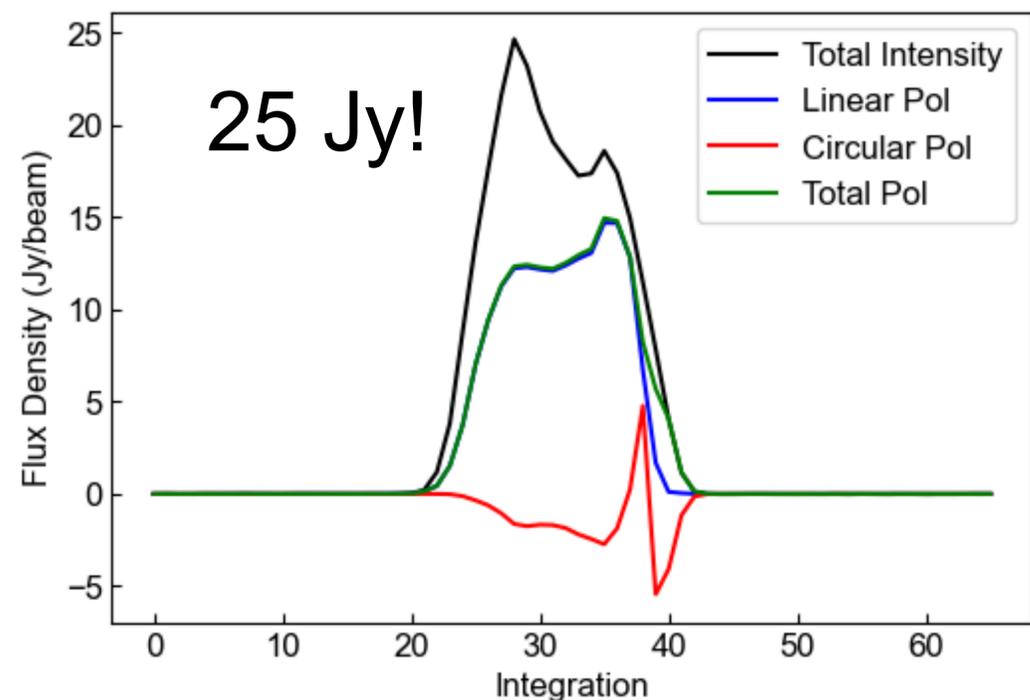
Ziteng Wang, Rea, ..., Wadiasingh et al. 2025 23



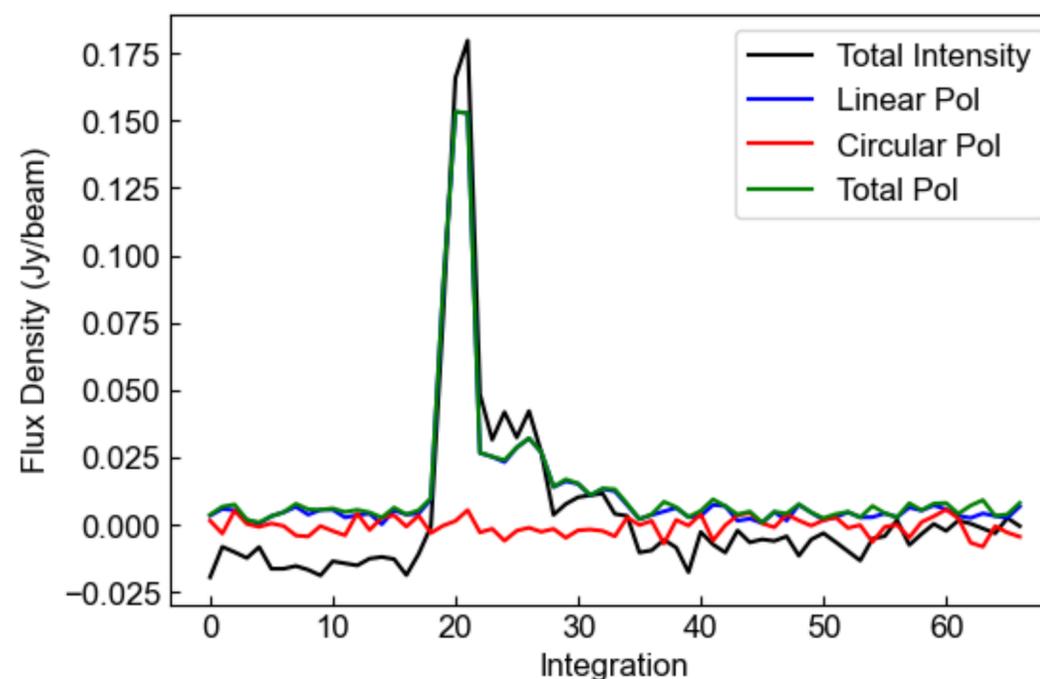
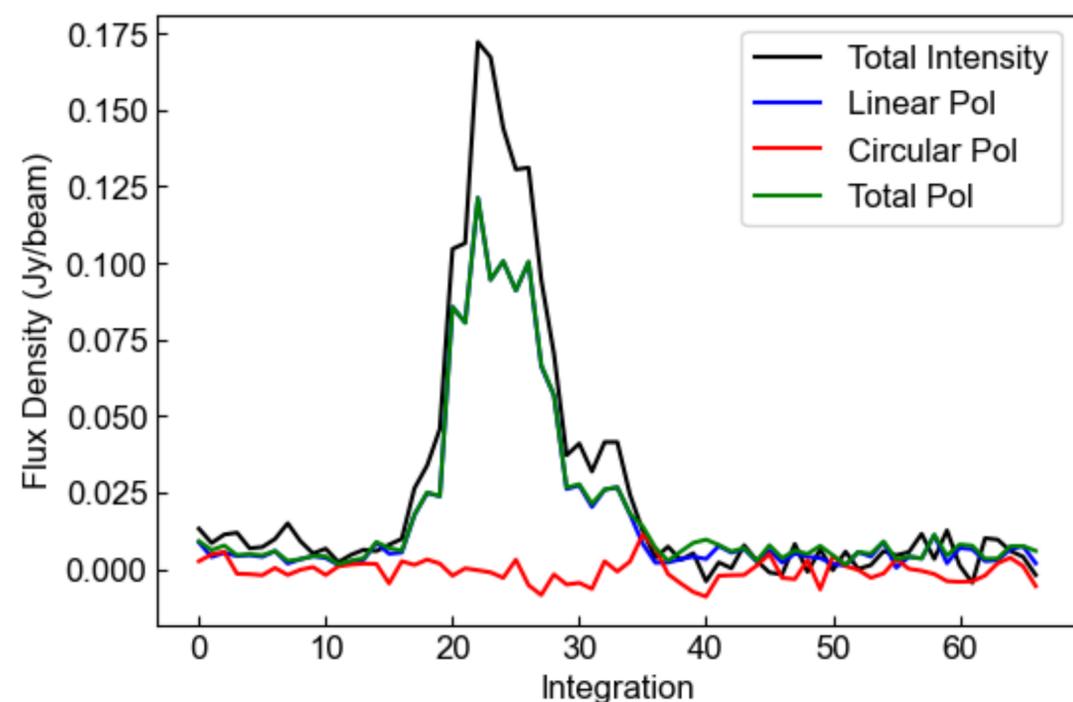
Observations of long period transients - ASKAP J1832-0911



Follow-ups of ASKAP J1832-0911



Ziten Wang, Rea,...Wadiasingh
et al. 2025



Variable pulse
morphology, polarisation
properties, radio spectral
index...

How might they exist?

Many possible mechanisms to spin down magnetars to long periods

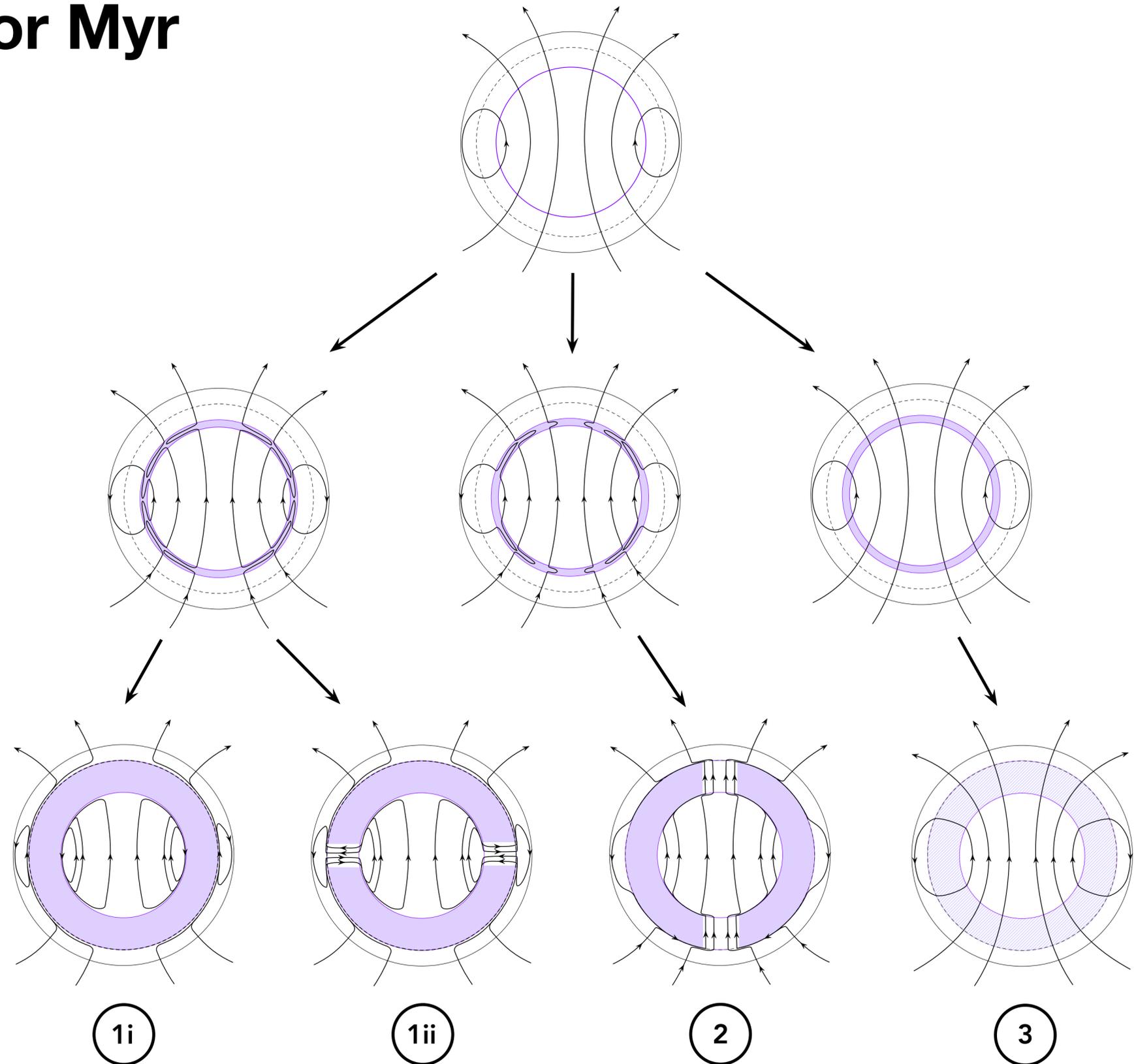
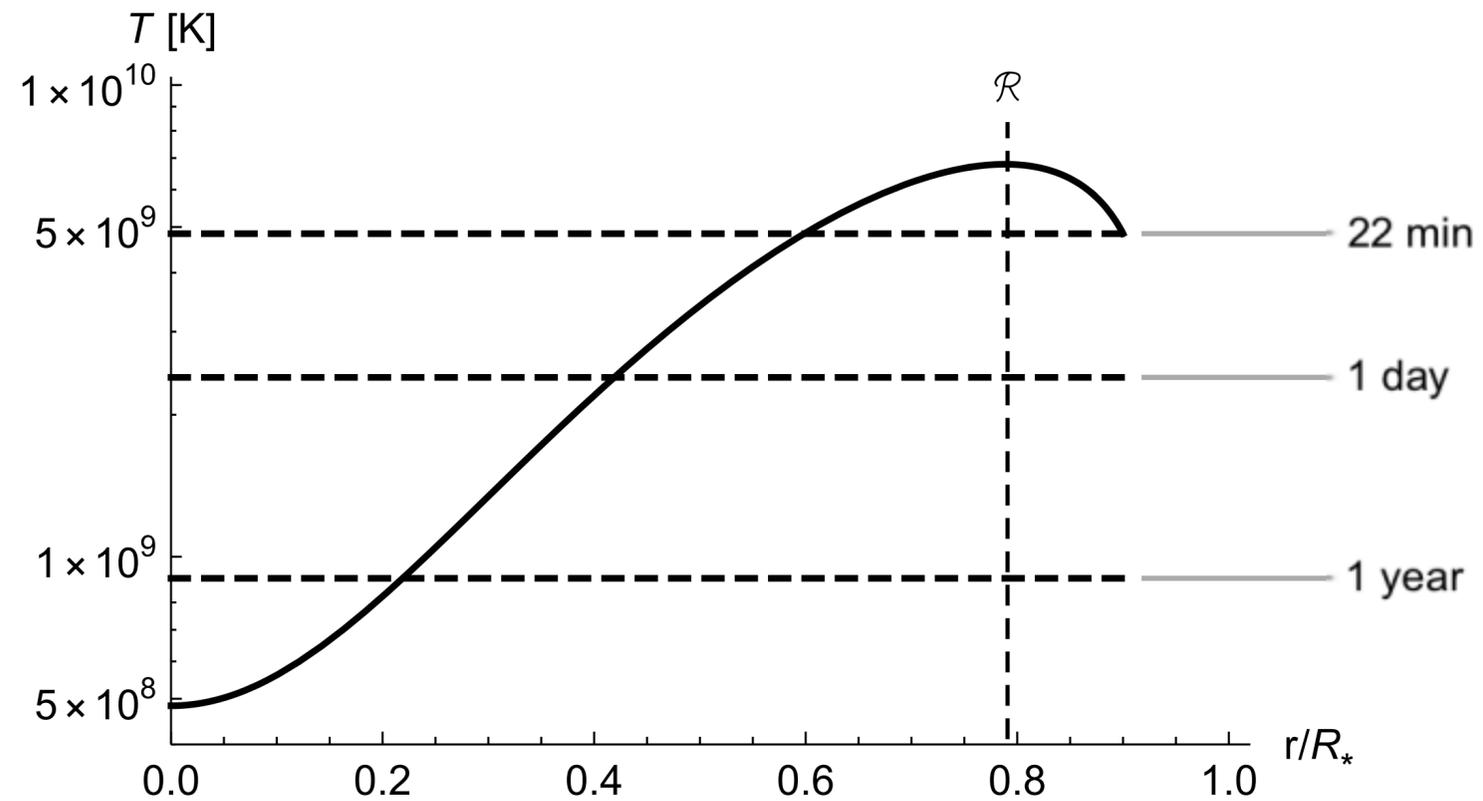
There is much phenomenological evidence for epochs of enhanced spindown in Galactic magnetars.

Physical mechanisms for attaining long periods:

- **Fallback disks** (Beniamini, **Wadiasingh**, Metzger+2020; Xu+2021; Ronchi, Rea+ 2022, Beniamini, **Wadiasingh**, Hare+2023;)
- **Enhanced spindown from monopolar particle winds and opening of magnetic flux** (Beniamini, **Wadiasingh**, Metzger+2020)
- **Giant flare kicks** (Beniamini, **Wadiasingh**, Metzger+2020)
- **Regular magnetic dipole spin-down persisting on a long-lived strong field**
- **Some or all of the above operating over the lifetime of the object**

Core Fields and Superconductivity

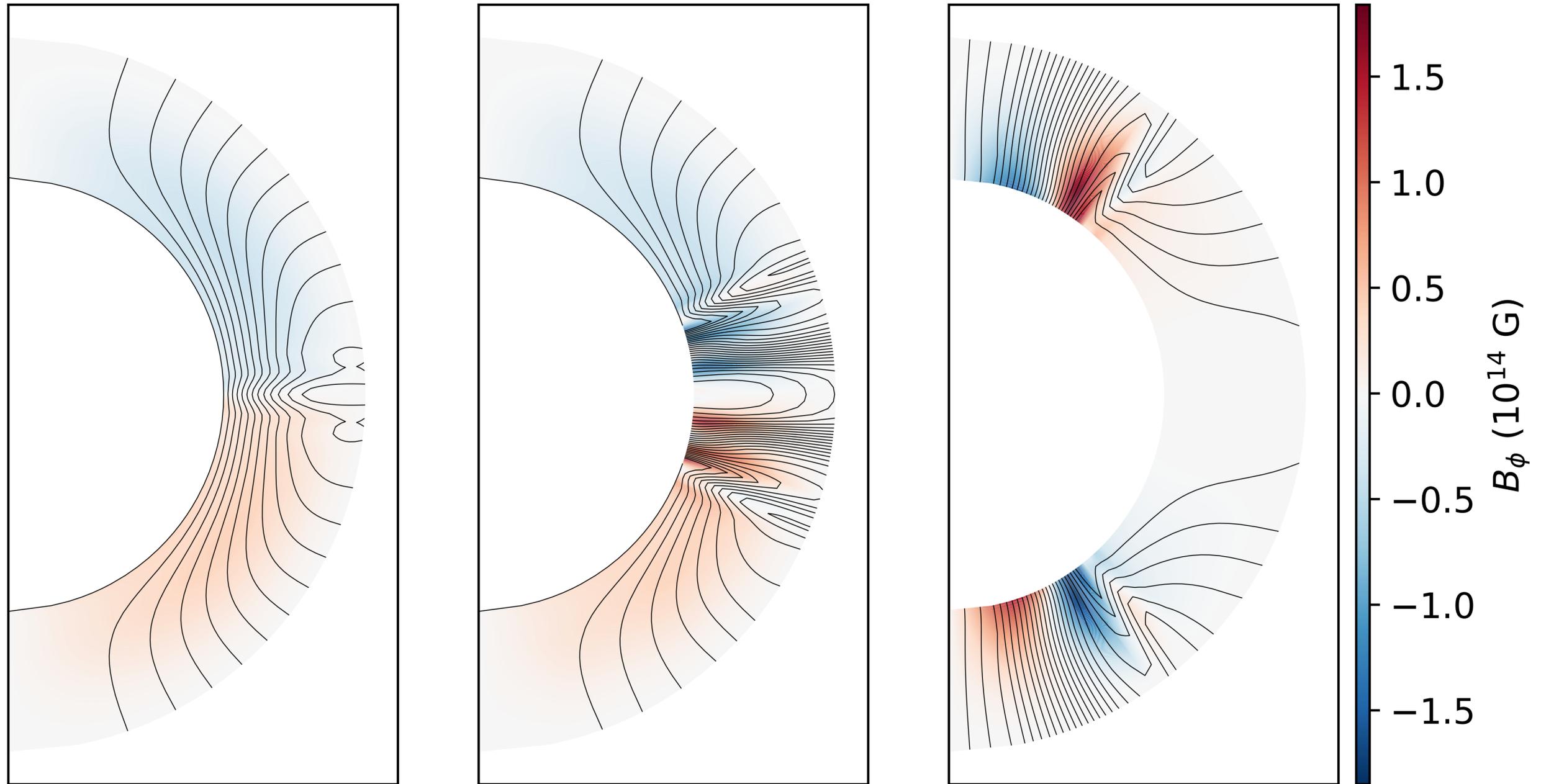
Survival of $>10^{15}$ G fields for Myr



Ho+2018;
Beniamini, **Wadiasingh**, Hare+2023;
Lander 2024;
Lander, Gourgoulatos, **Wadiasingh**, Antonopoulou 2024

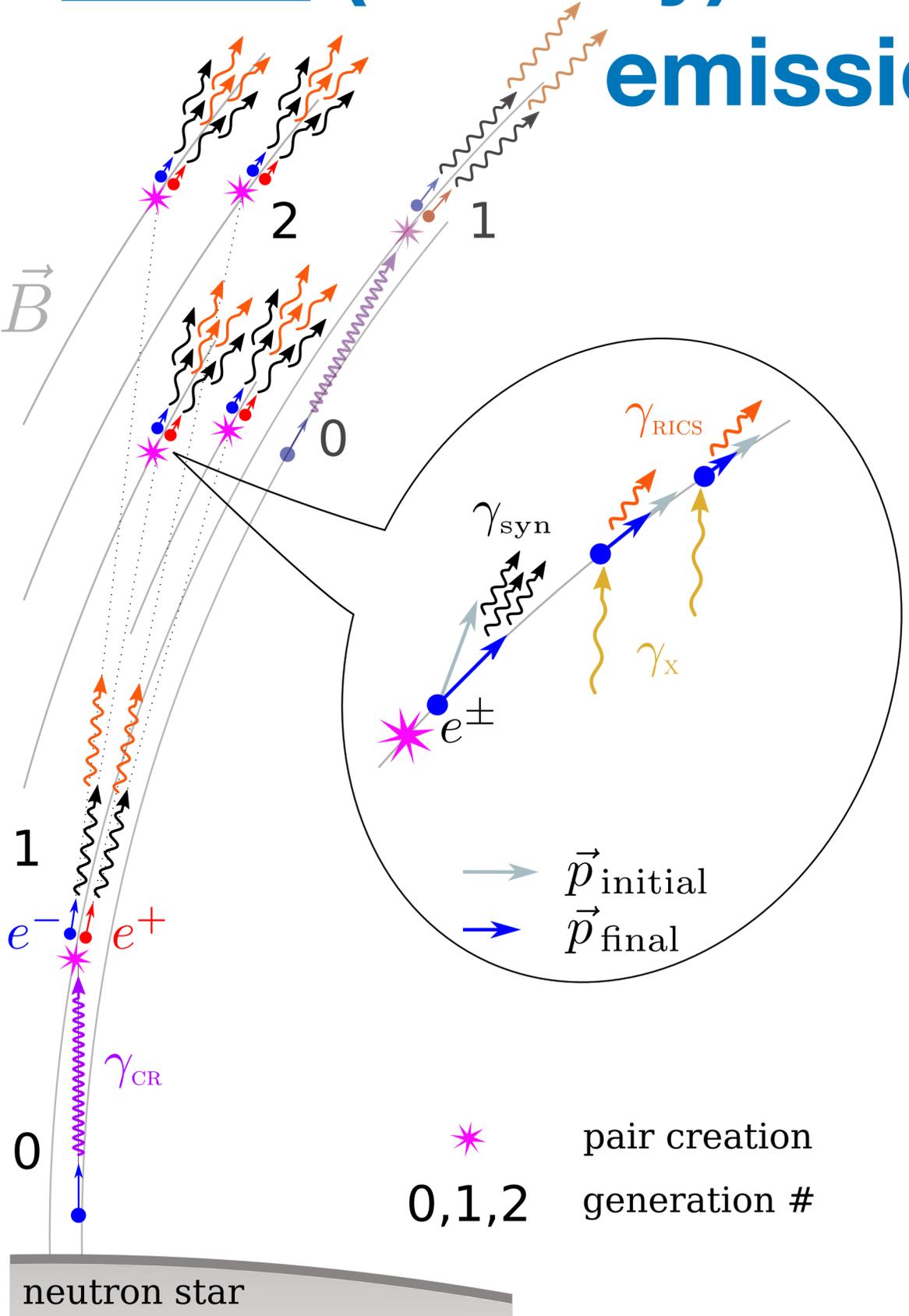
Core Fields and Superconductivity

Survival of $>10^{15}$ G fields and activity for Myr



**How might the radio emission
be produced?**

We know (broadly) what is required for pulsar-like radio emission: **pair cascades**



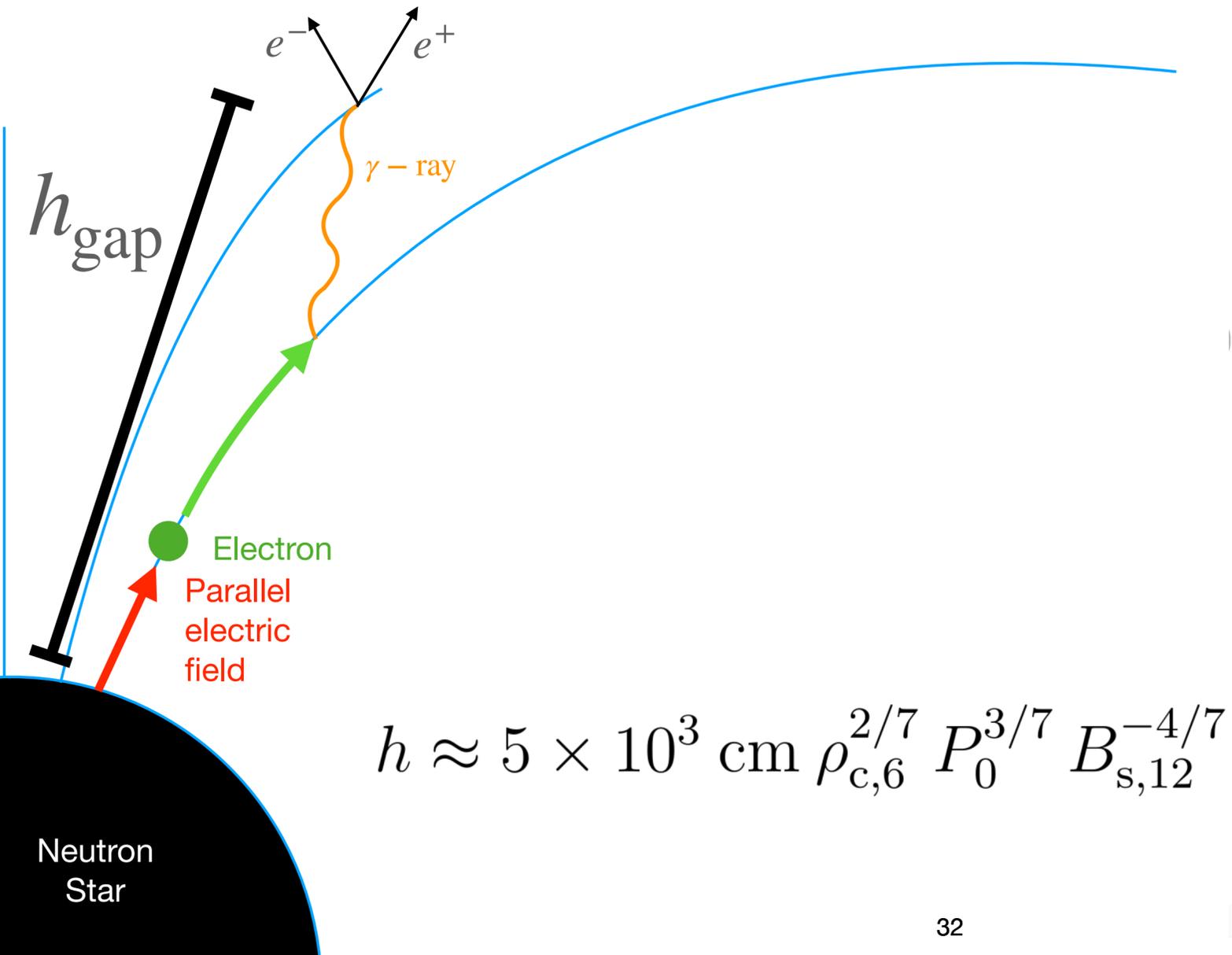
$$\gamma + B \rightarrow e^+ + e^-$$

Pairs are produced in the ground Landau state in magnetar high-B regime

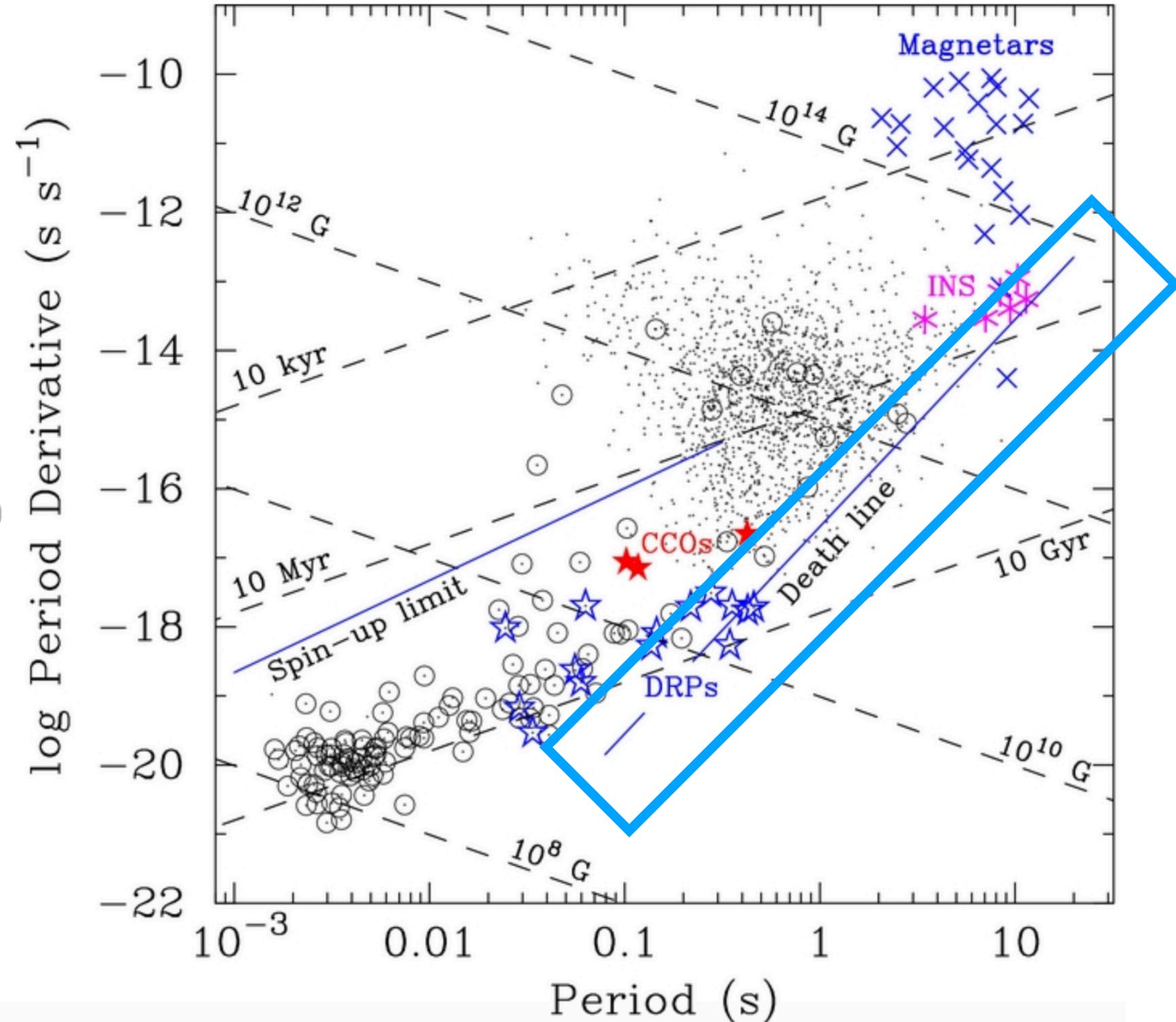
e.g., Sturrock (1971), Baring & Harding (1997), Baring & Harding (2001), Timokhin (2010), Timokhin & Arons (2013) Timokhin & Harding (2015, 2019)

The Pair Curvature Death Line

- Computing the unscreened gap height allows luminosity estimates



$$h \approx 5 \times 10^3 \text{ cm } \rho_{c,6}^{2/7} P_0^{3/7} B_{s,12}^{-4/7}$$

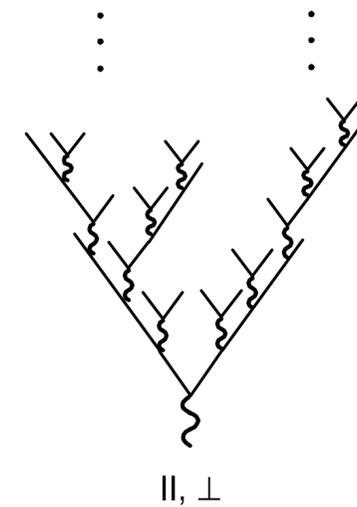


How are magnetar pair cascades different?

- Resonant inverse Compton scattering (RICS, effectively cyclotron absorption+ immediate emission) channel photons can compete versus the curvature radiation
- Radiative losses can be catastrophic in the RICS channels
- Photon splitting can suppress pair production efficiency
- Pairs produced are largely in the ground Landau state, so do not effectively radiate synchrotron photons to continue a cascade
- Non-resonant scatterings are suppressed

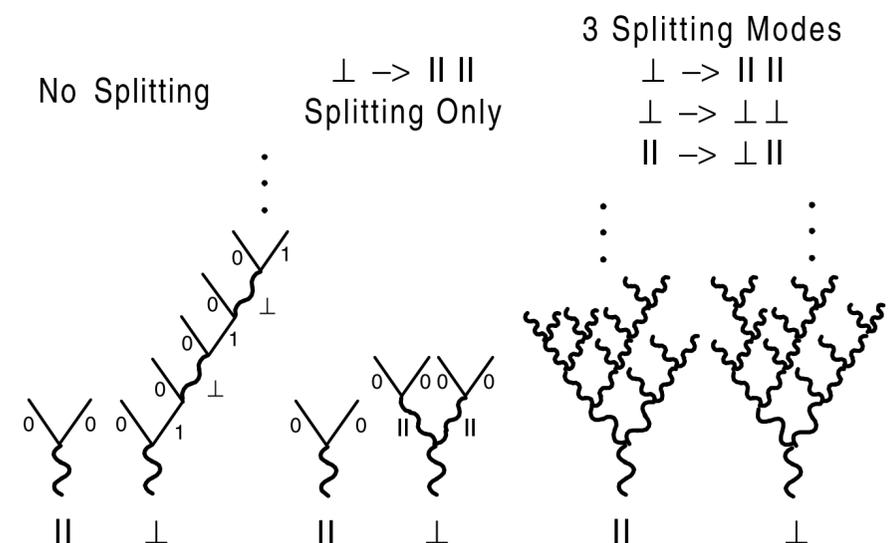
LOW FIELD CASCADES

$$B < 0.1 B_{cr}$$



HIGH FIELD CASCADES

$$B > 0.1 B_{cr}$$



Baring & Harding (2001)

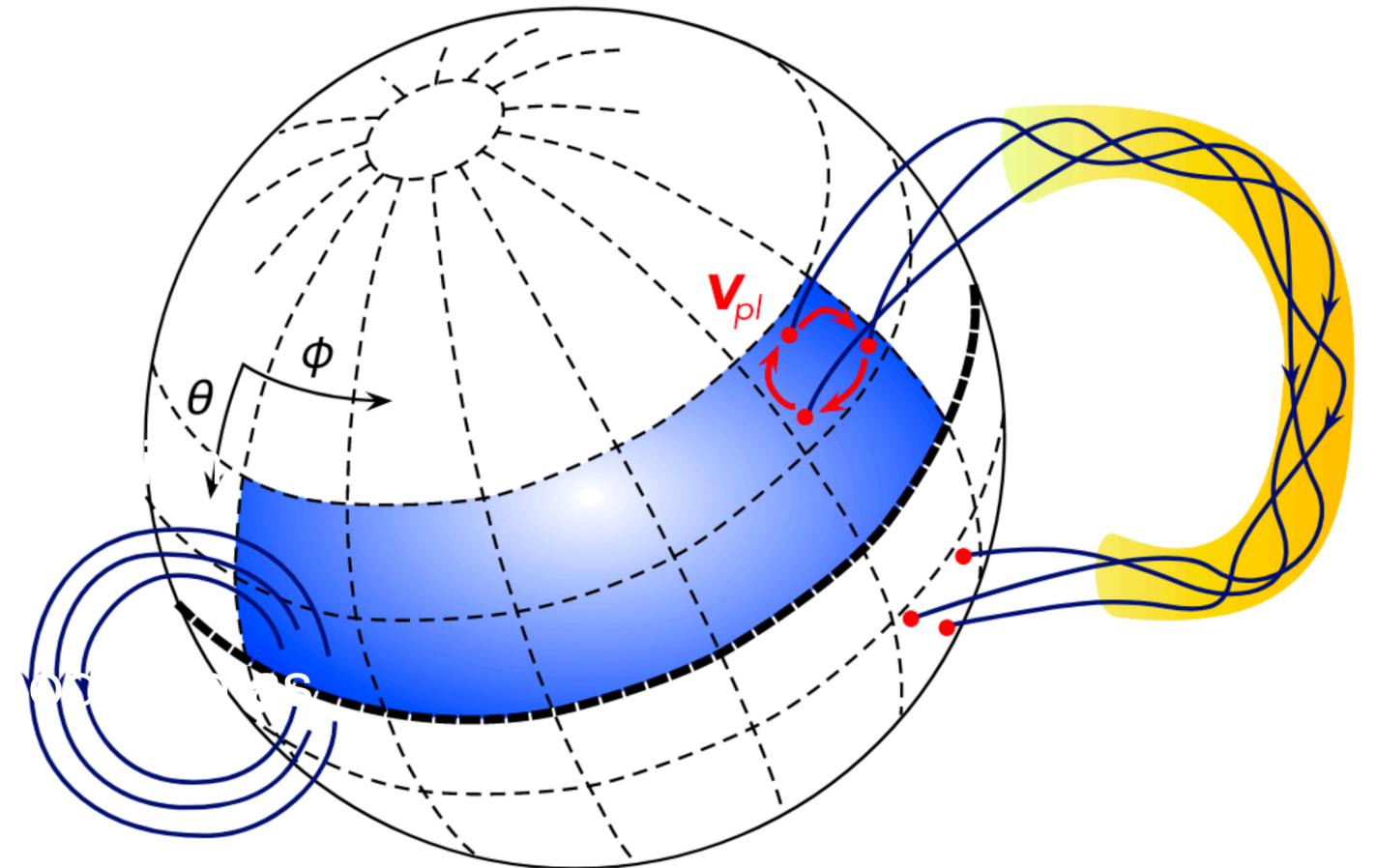
Ingredient 1: Plastic Flow powered twists

- At $B > 10^{15}$ G **plastic flow** (magnetic stress driven ‘continental drift’) dominates crustal evolution, and can twist the magnetic field

See e.g. Lander 2019; Gourgouliatos & Lander 2021; Lander 2023

$$\psi \equiv \frac{B_{\phi}}{B_{dp}} < 1$$

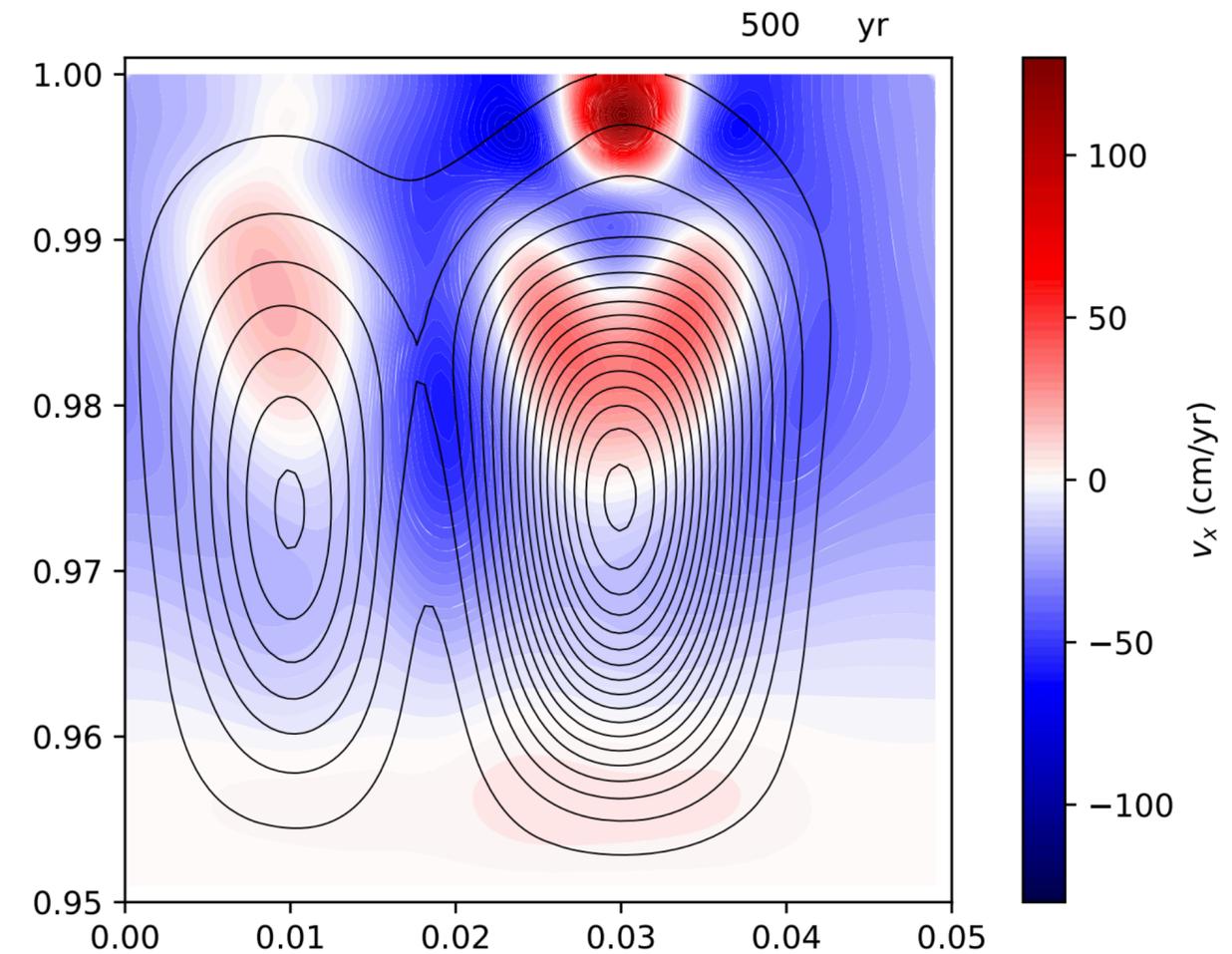
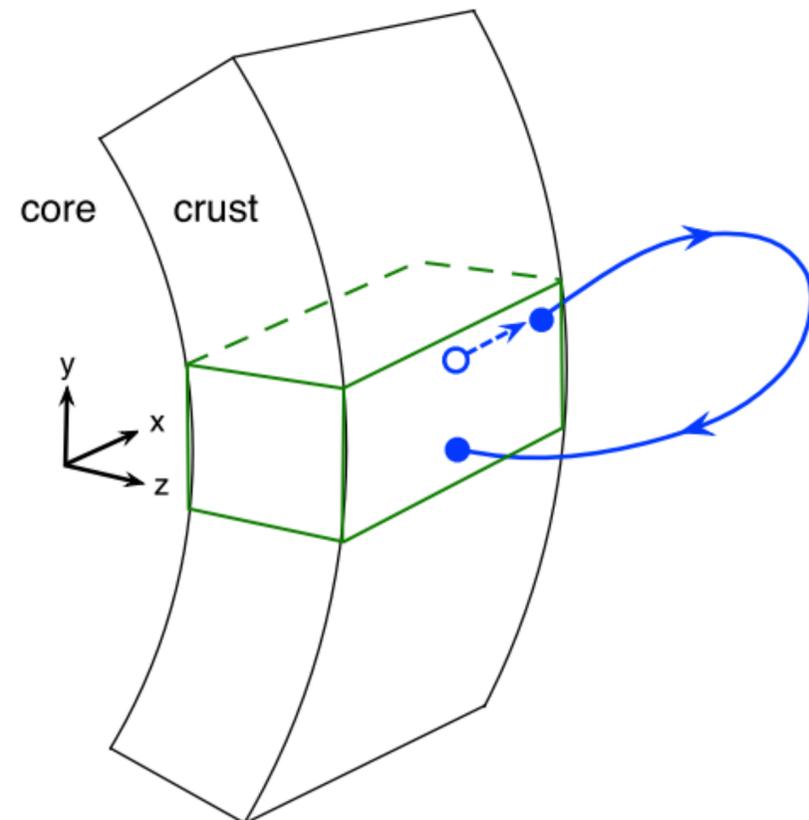
Twists require charge similarly to rotation
– Analogous to preexisting “low-twist” model for FRBs proposed by Wadiasingh & Timokhin 2019, Wadiasingh+2020



Lander et al., 2023

Ingredient 1: Plastic Flow powered twists

- At $B > 10^{15}$ G **plastic flow** (magnetic stress driven ‘continental drift’) dominates crustal evolution, and can twist the magnetic field
- Lander et al (2019) show that plastic flow of 1km^2 patches has velocities of **1-100 cm/yr** for months to decades

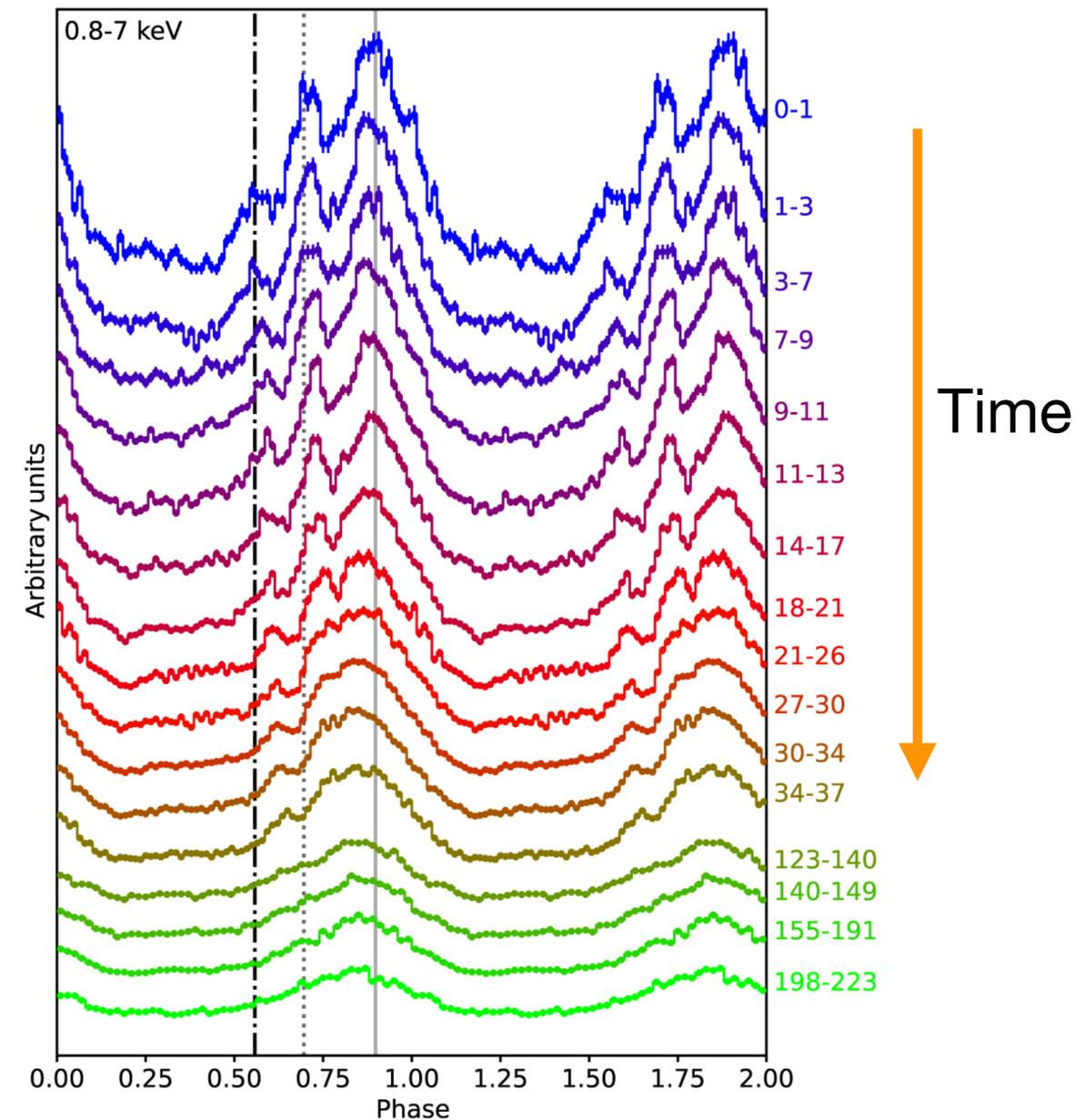


Lander, 2019

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- Lander et al (2019) show that plastic flow of 1km^2 patches has velocities of **1-100 cm/yr** for months to decades
- Younes et al. (2022) posit a connection between X-ray magnetar pulse peak migration plastic flow, requiring 10^6 cm/yr

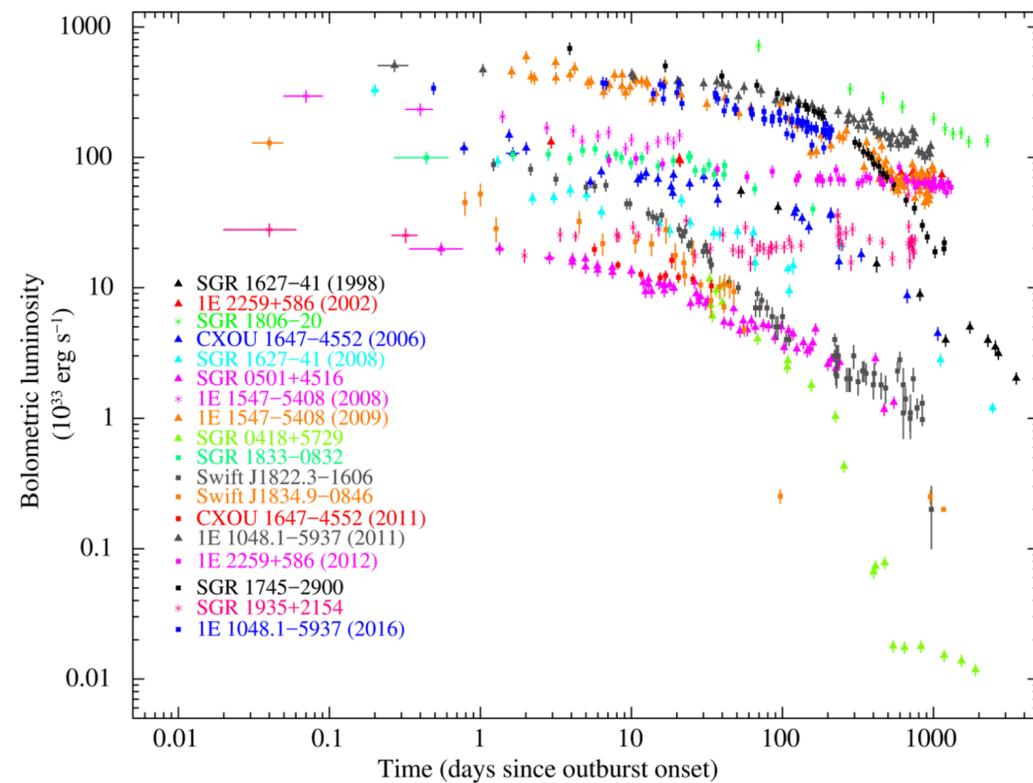
See also Younes et al. (2025)



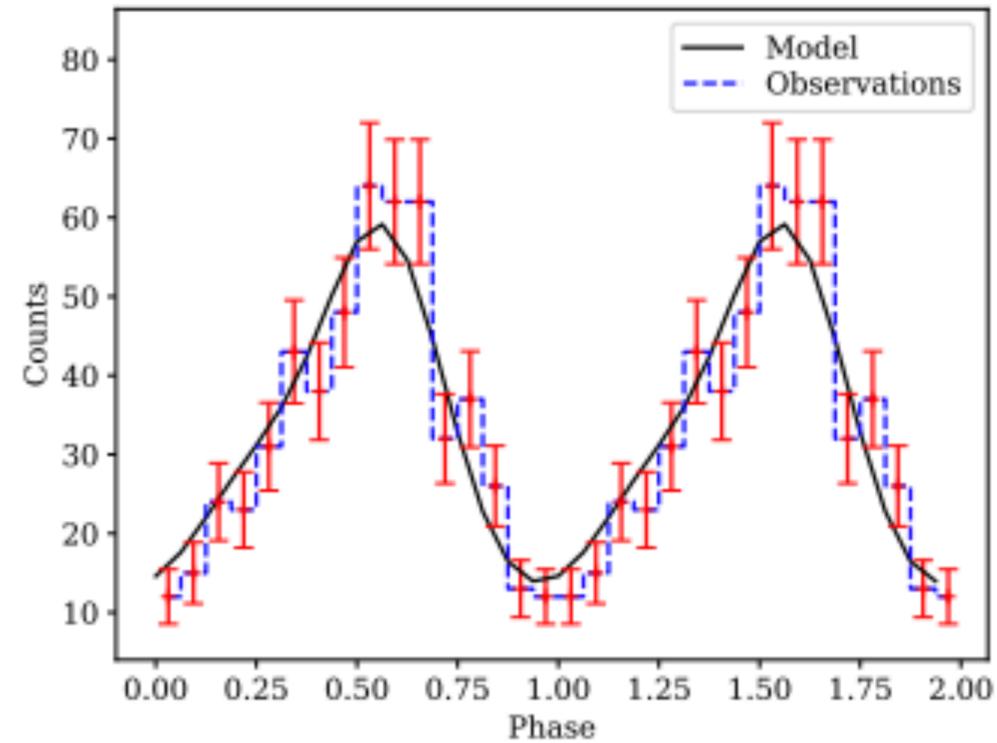
Younes+ 2022

Ingredient 2: Twist-induced radio emission

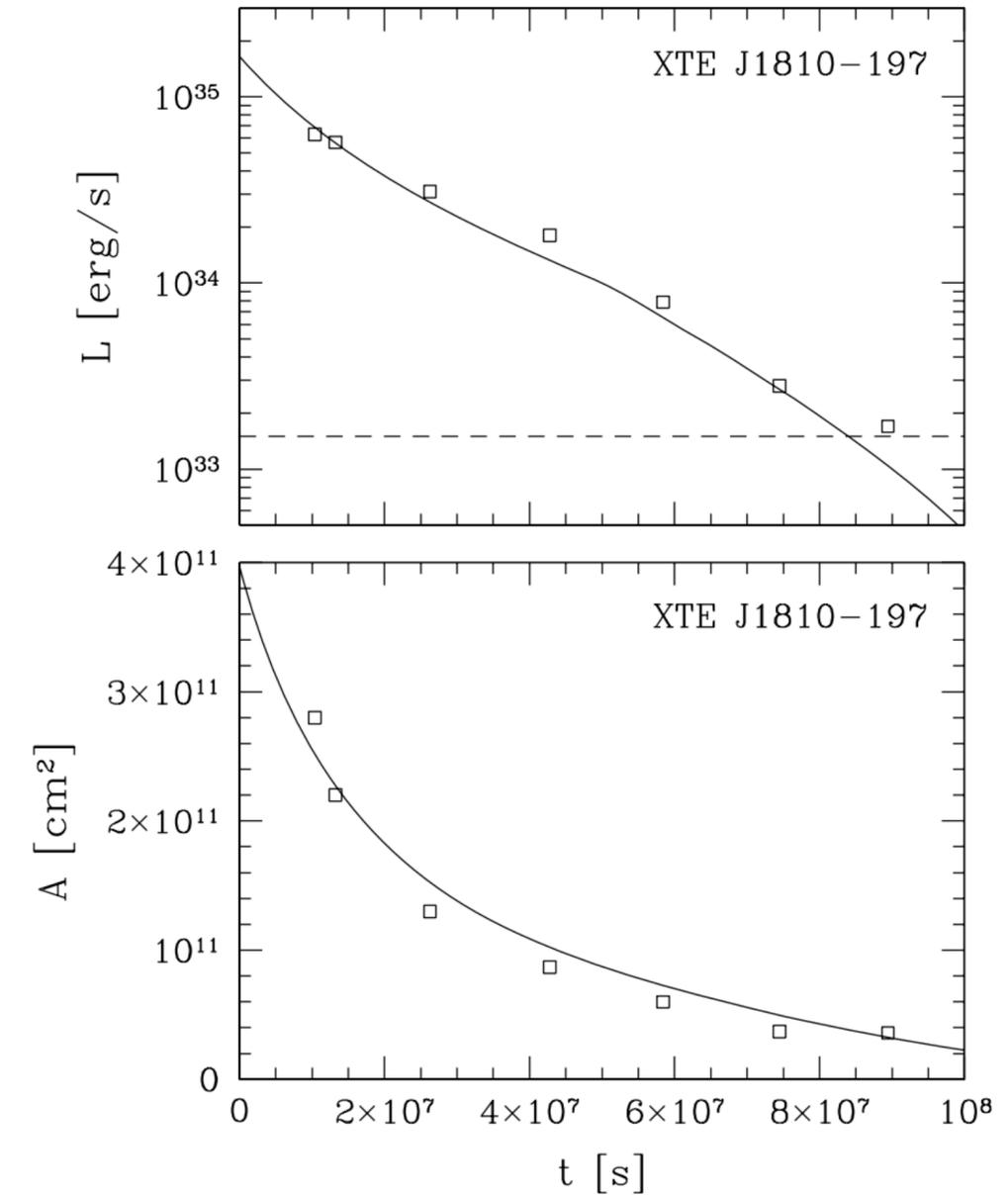
- Persistent magnetar activity can be described by twisted magnetic fields



X-ray magnetars in outburst (Coti Zelati+ 2018)



Pulse fit with twisted field (Igoshev+ 2021)



Long-term twist evolution (Beloborodov, 2013)

Ingredient 2: Twist-induced radio emission

- Persistent magnetar activity can be described by twisted magnetic fields

Rotation: $\rho_{\text{GJ}} \approx \frac{2B}{cP}$

Twist: $\rho_{\psi} \approx \frac{B \sin^2(\theta_{\text{fp}}) \Psi}{4\pi R_{\text{NS}}}$

- Twists require charge similarly to rotation

Analogous to preexisting “low-twist” model for FRBs proposed by Wadiasingh & Timokhin 2019, Wadiasingh+2020

Crustal density and temperature gradients also drive magnetic field evolution

Alternative: Thermoelectrically induced small twists

$$\frac{\delta \vec{B}}{\delta t} = -c \nabla \times \left(\frac{(\nabla \times \vec{B}) \times \vec{B}}{4\pi q n_e} + \frac{c \nabla \times \vec{B}}{4\pi \sigma_{\text{con}}} - \frac{S_e}{q} \nabla T \right)$$

$d\vec{B}/dt$ orthogonal to both gradients > twisted component

Positive feedback loop: $dT/dx \neq 0 > \text{Twist} > \text{Particle acceleration} > \text{Hotspot heating} > dT/dx \neq 0$

Ingredient 2: Twist-induced radio emission

- Persistent magnetar activity can be described by twisted magnetic fields

$$\rho_{\text{GJ}} \approx \frac{2B}{cP}$$

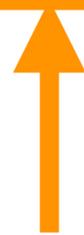
$$\rho_{\psi} \approx \frac{B \sin^2(\theta_{\text{fp}}) \Psi}{4\pi R_{\text{NS}}}$$

Analogous to preexisting “low-twist” model for FRBs proposed by Wadiasingh & Timokhin 2019, Wadiasingh+2020

- Twists require charge similarly to rotation

$$\Psi_{\text{crit}} \equiv \Psi(\rho_{\psi} = \rho_{\text{GJ}})$$

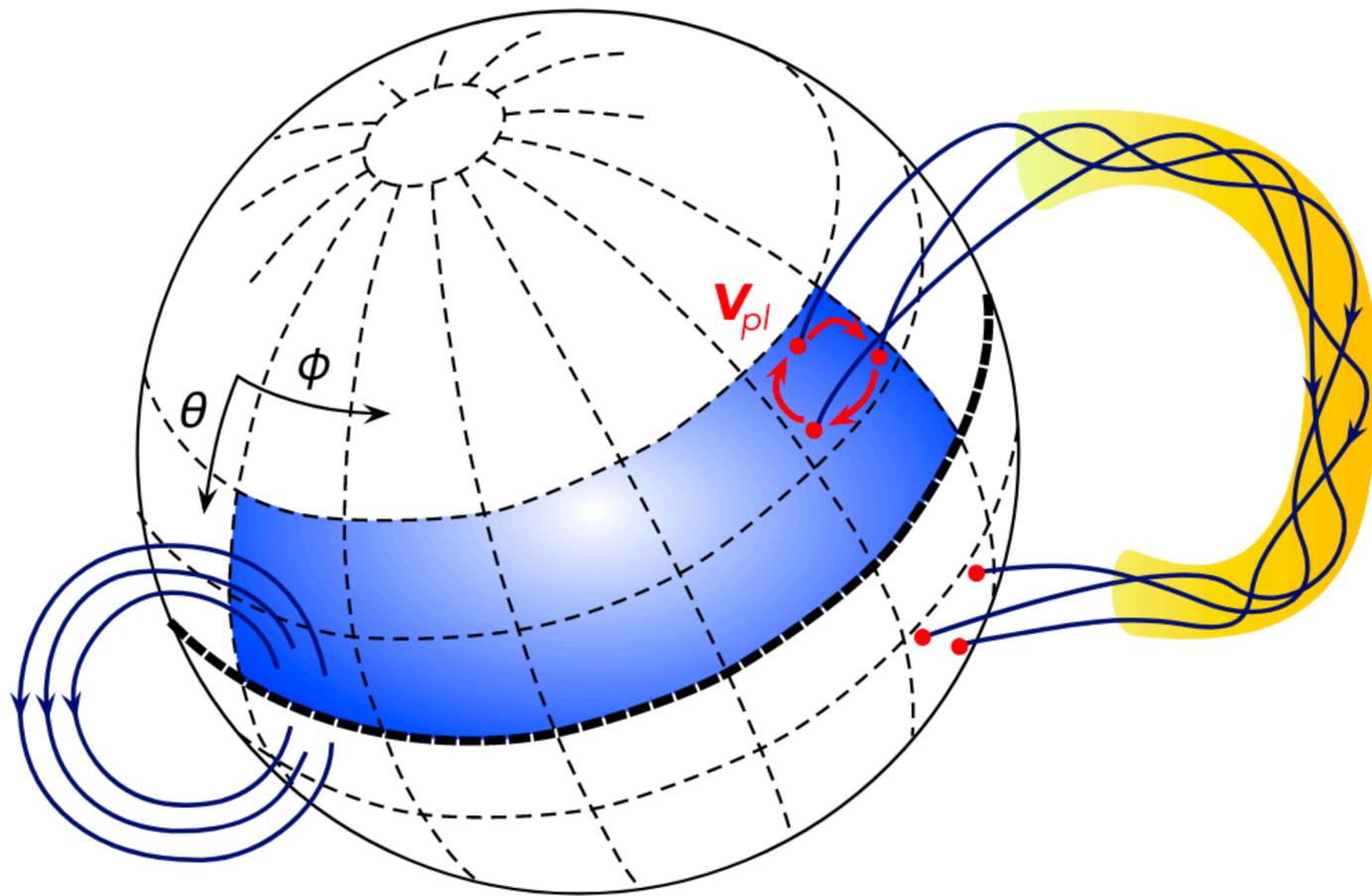
Broadband radio emission
(e.g. Timokhin 2010, Timokhin & Arons 2013, Philippov et al., 2020, Benáček+2024,2025)



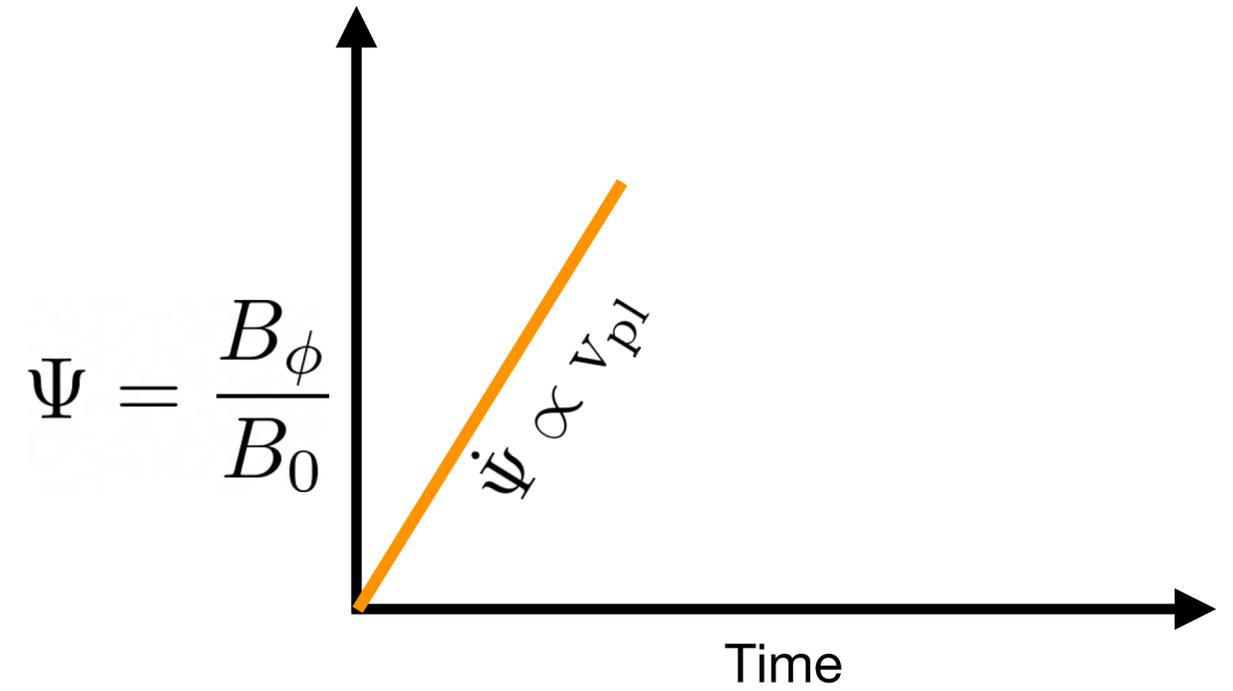
If $\rho_{\psi} > \rho_{\text{GJ}}$  Twist dissipation  Acceleration gaps  Pair production

Model: Five steps

(1) Plastic flow twists field lines in $\sim\text{km}^2$ patches



Lander 2023



Model: Five steps

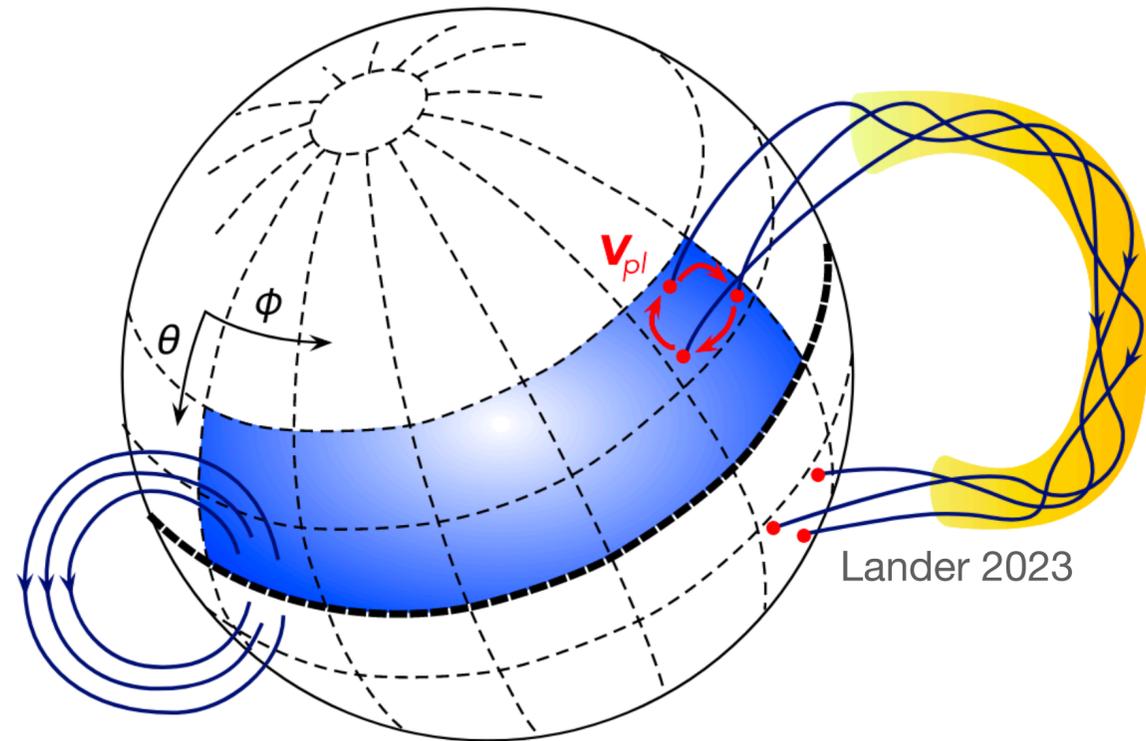
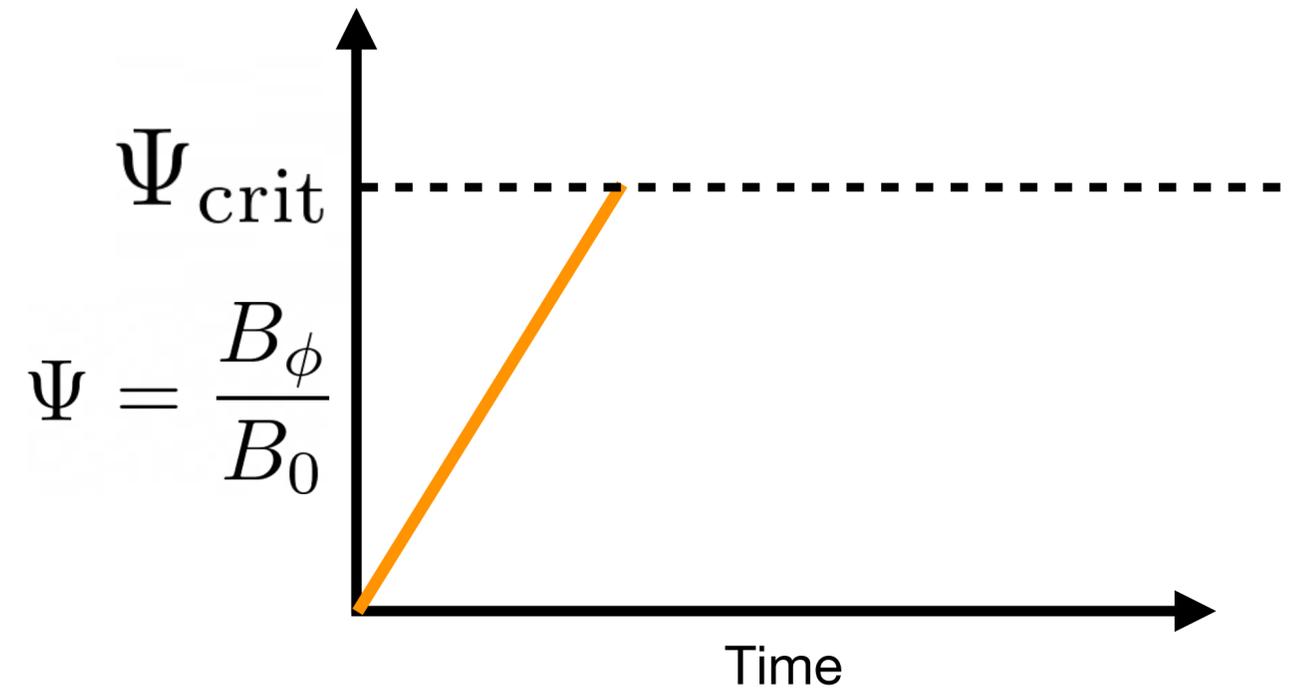
(1) Plastic flow twists field lines in $\sim\text{km}^2$ patches

(2) Twist current requirements exceed GJ current

$$\rho_\psi > \rho_{GJ}$$

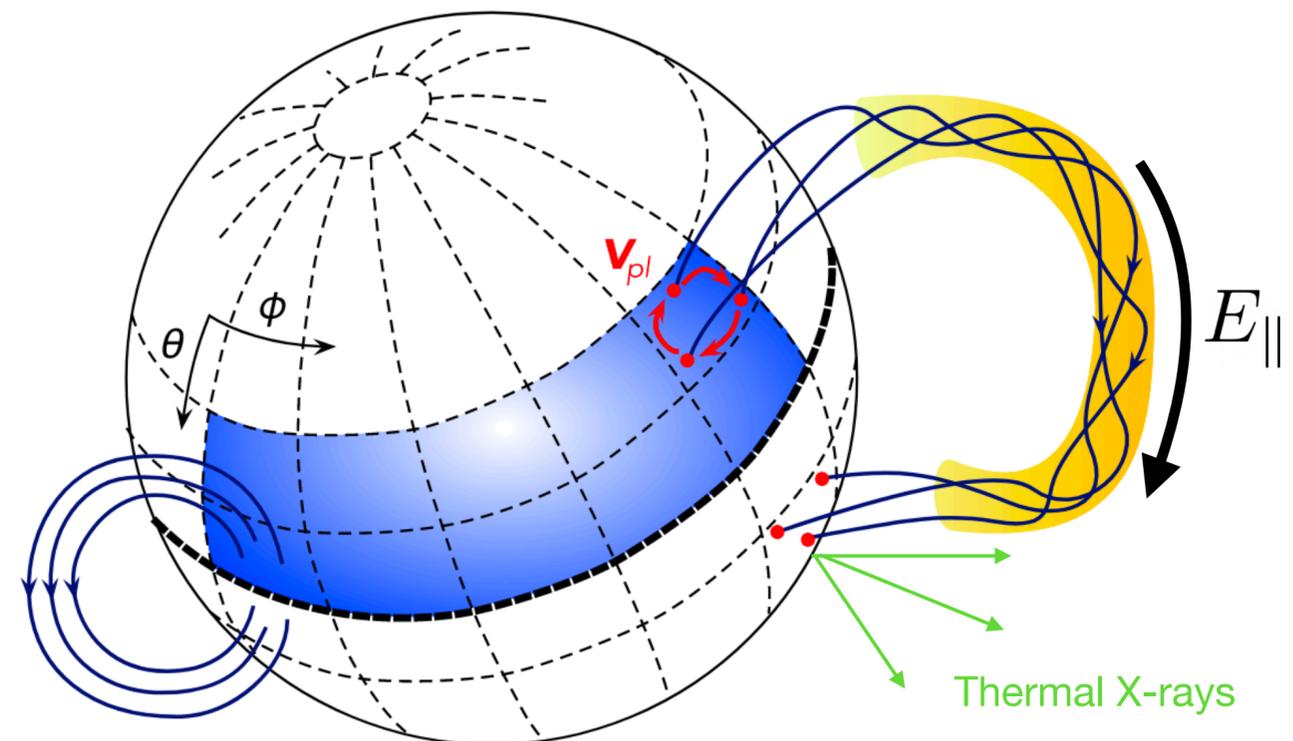
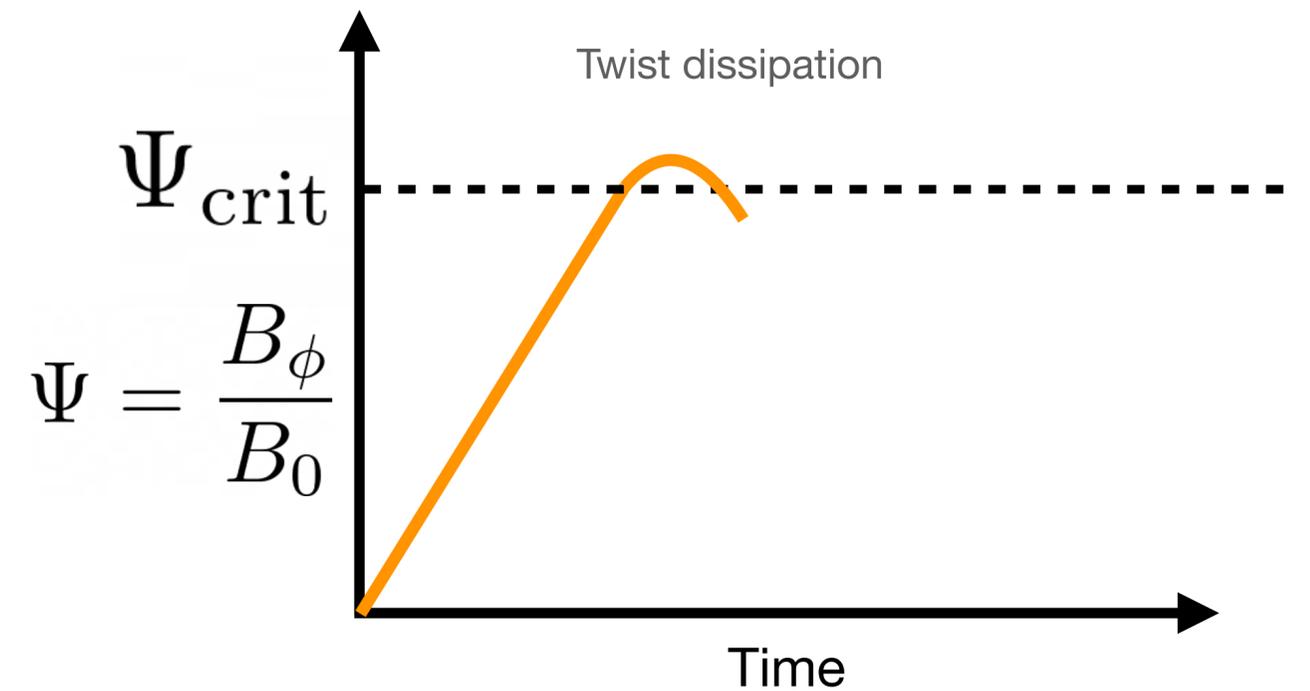
$$\Psi_{\text{crit}} \approx 10^{-4} P_{\text{NS},3}^{-1} \theta_{\text{fp},-1}$$

Analogous to preexisting “low-twist” model for FRBs proposed by Wadiasingh & Timokhin 2019, Wadiasingh+2020



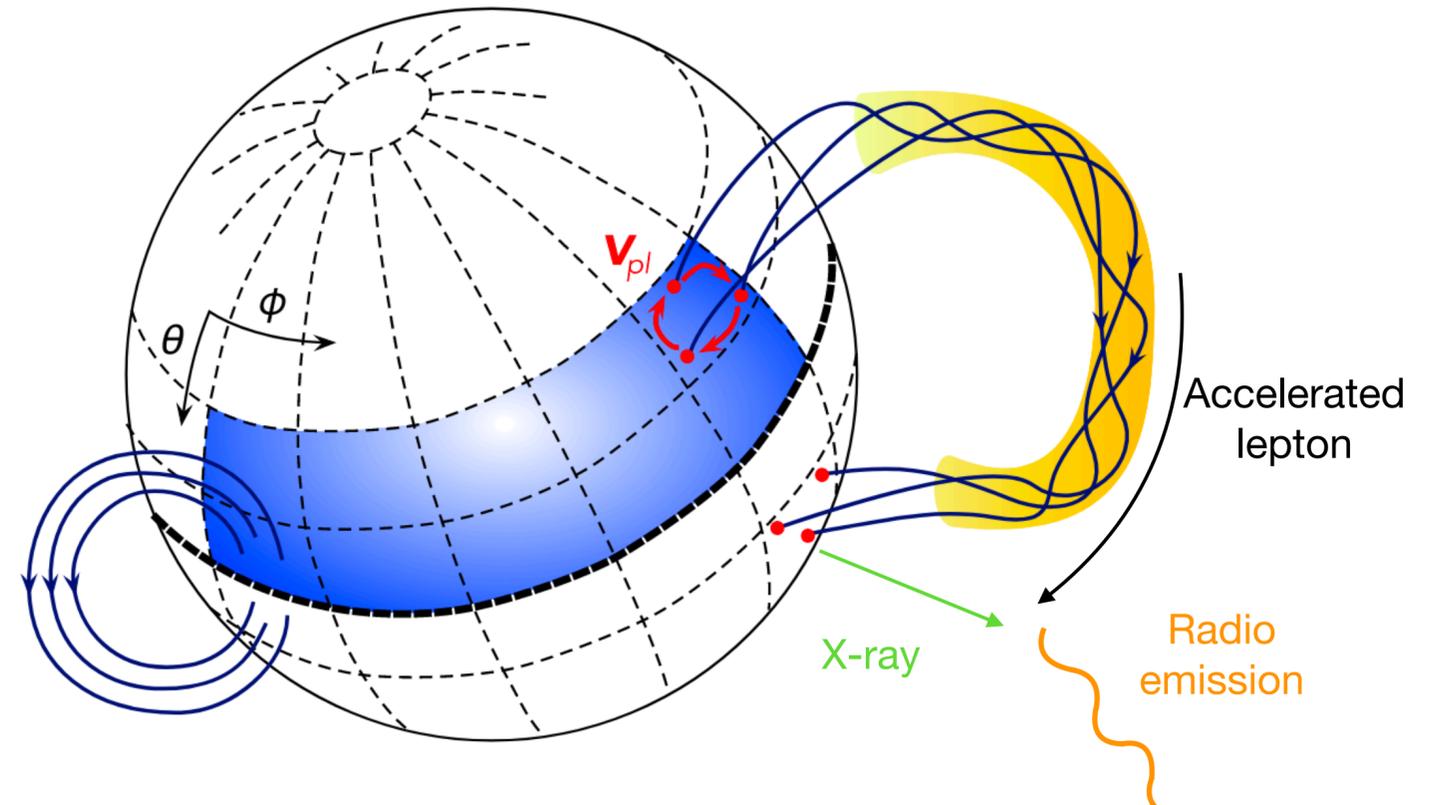
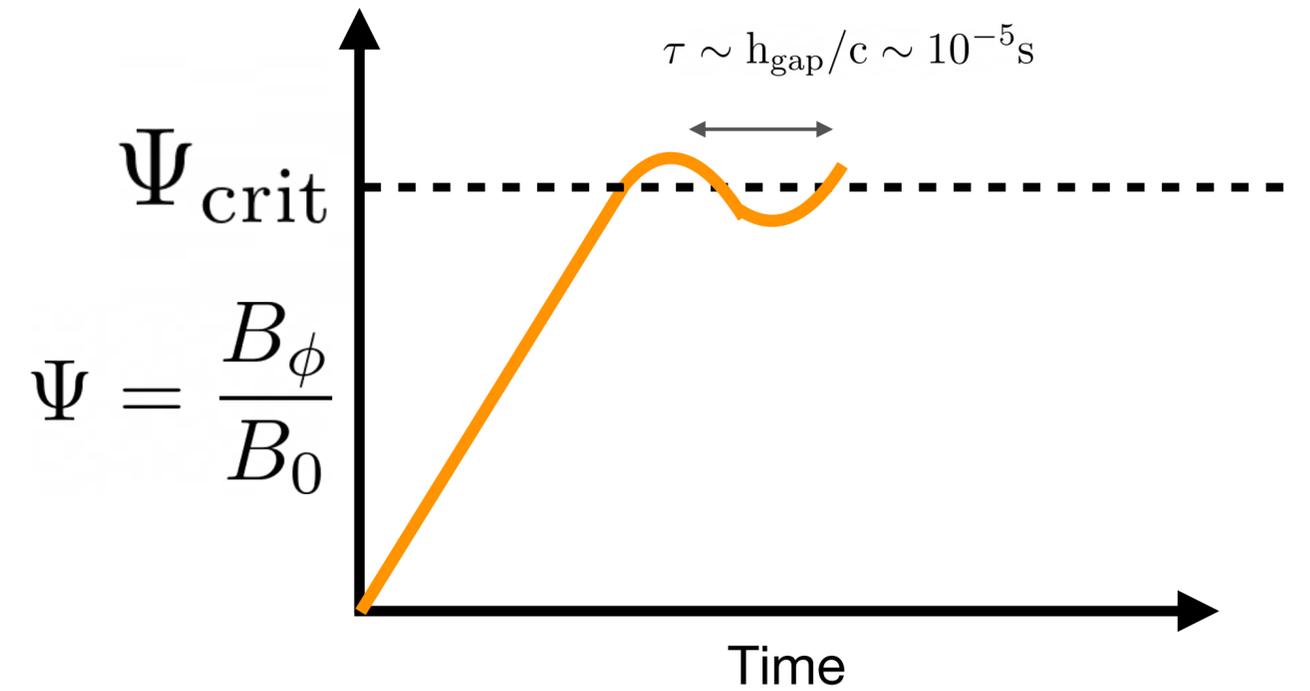
Model: Five steps

- (1) Plastic flow twists field lines in $\sim\text{km}^2$ patches
- (2) Twist current requirements exceed GJ current
- (3) Voltage increases rapidly accelerating particles
 - (a) Particles bombard surface producing thermal UV/X-ray counterpart



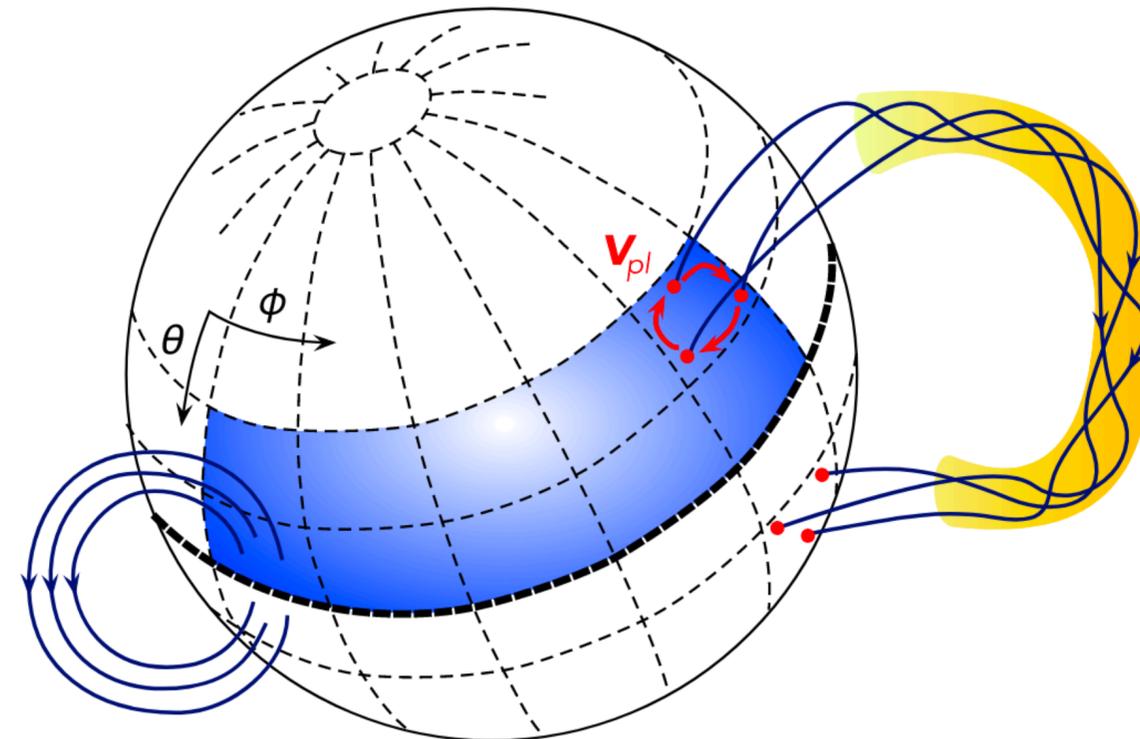
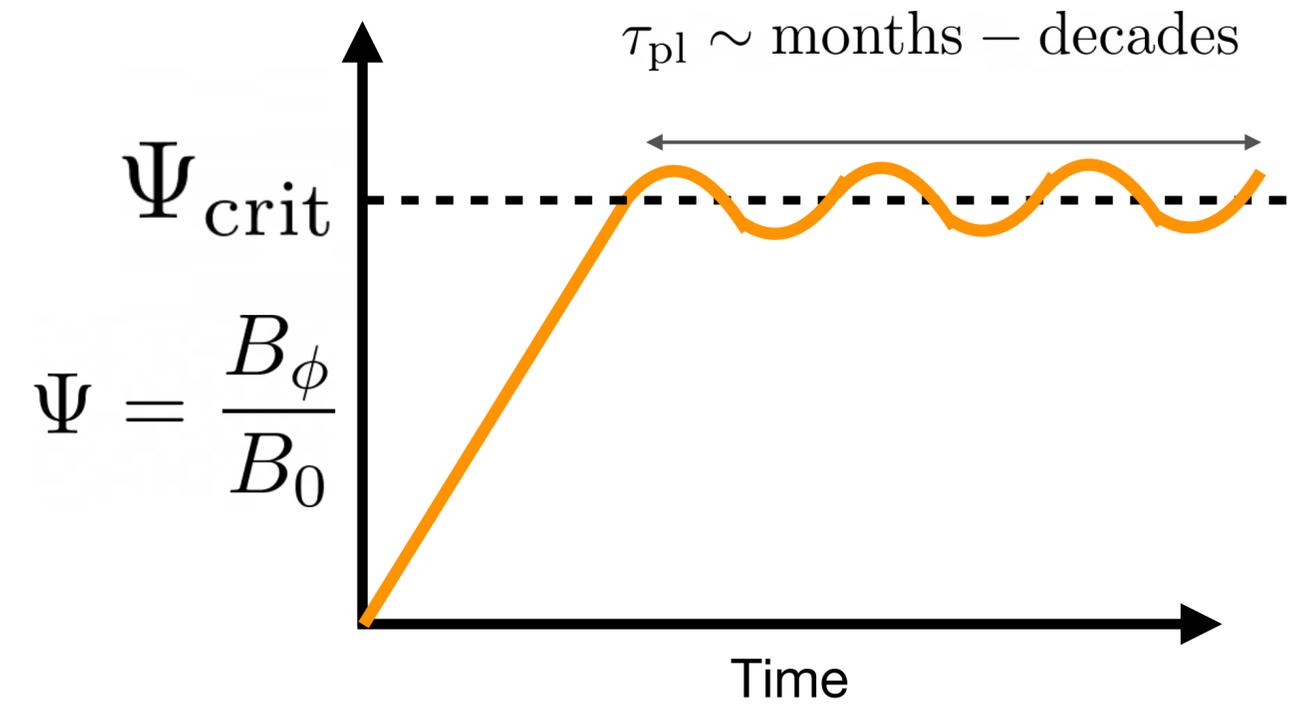
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- (4) Particles produce pairs via **RICS or curvature** photons, screen E-field, emitting radio waves

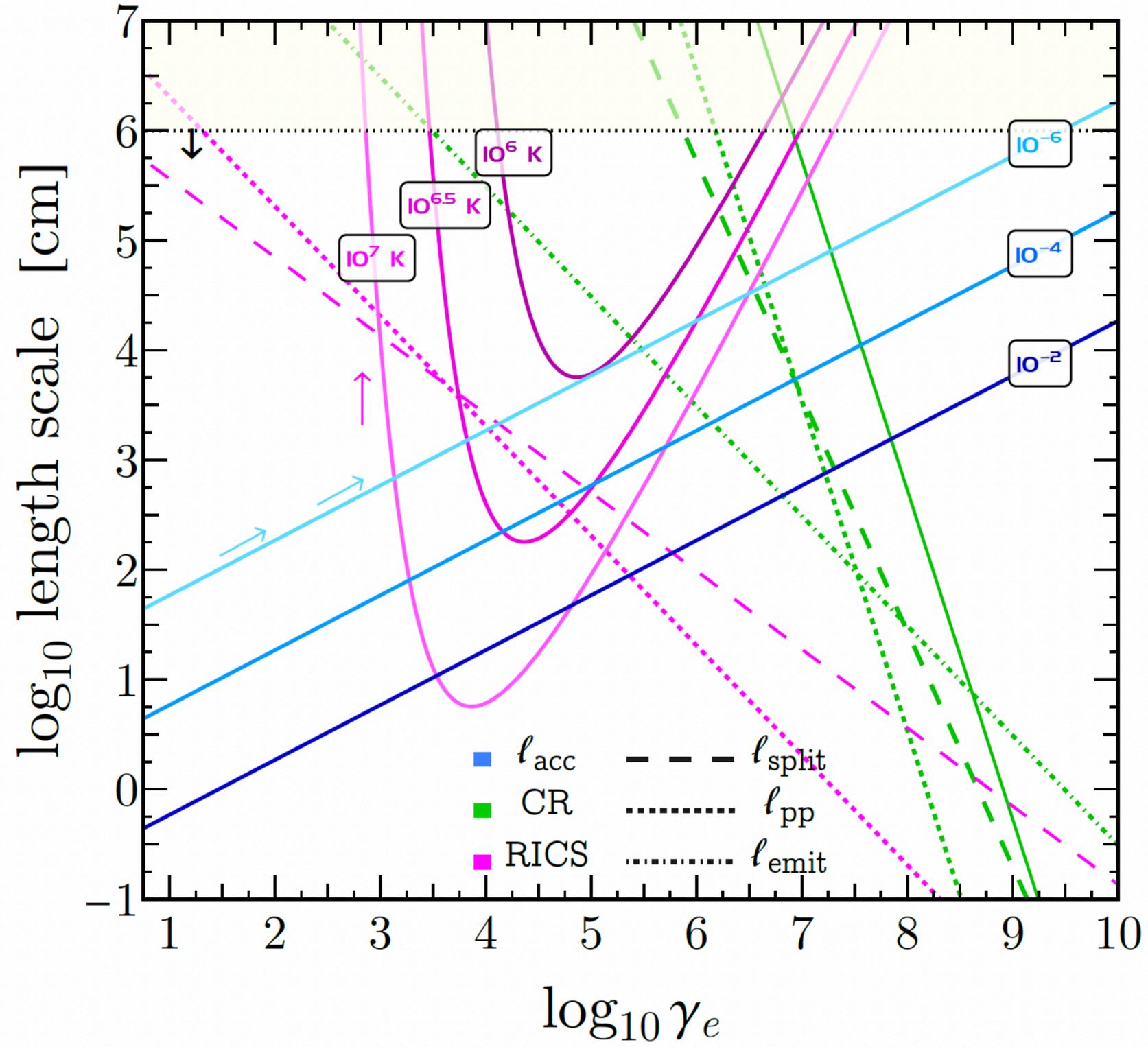


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 - (a) Particles bombard surface producing thermal UV/X-ray counterpart
- (4) Particles produce pairs via **RICS** or **curvature** photons, screen E-field, emitting radio waves
- (5) Twist oscillates around critical value, with stable dissipation as long as plastic motion continues



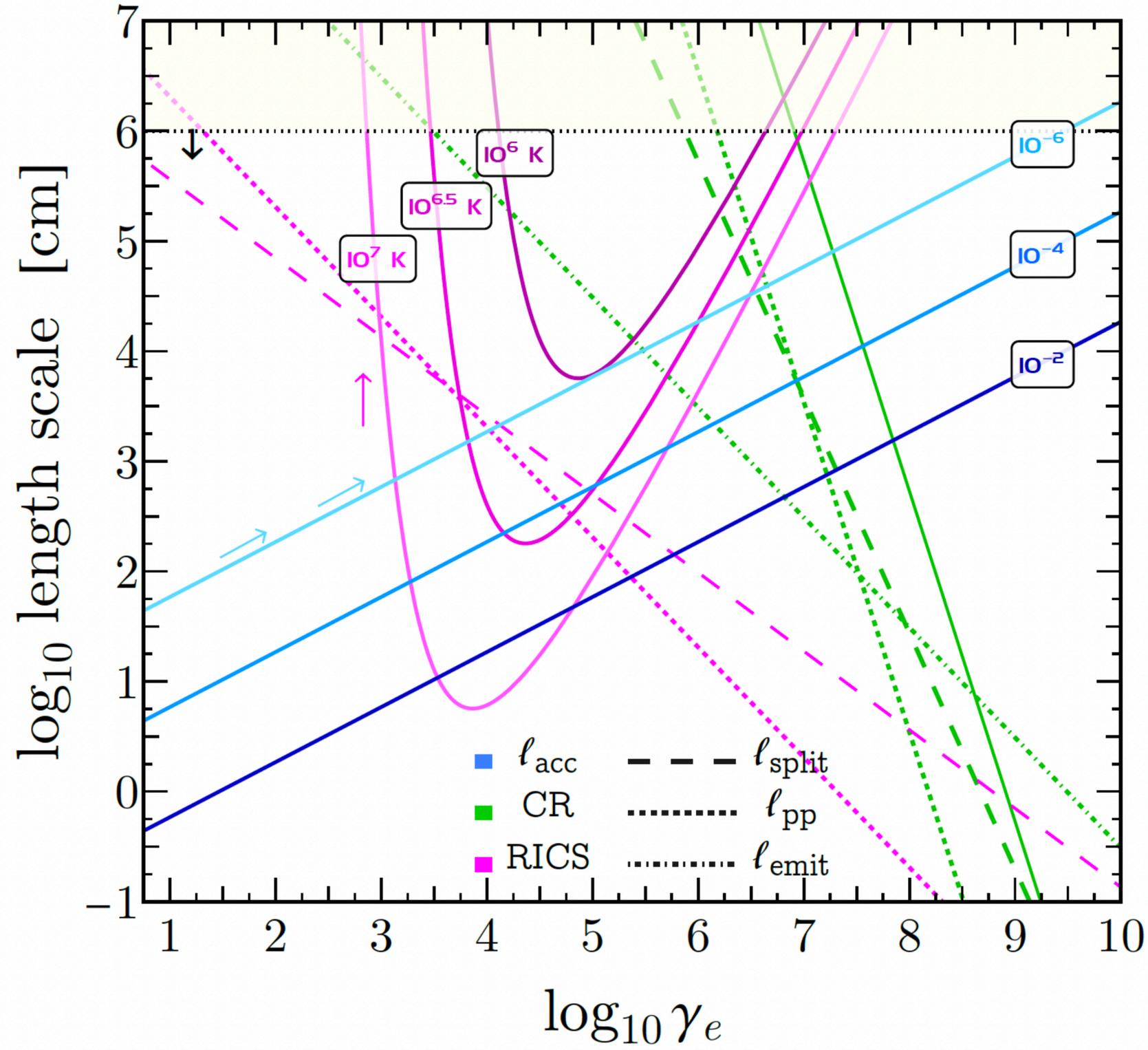
Microphysics



Particles accelerated
as they cross gap height

Microphysics

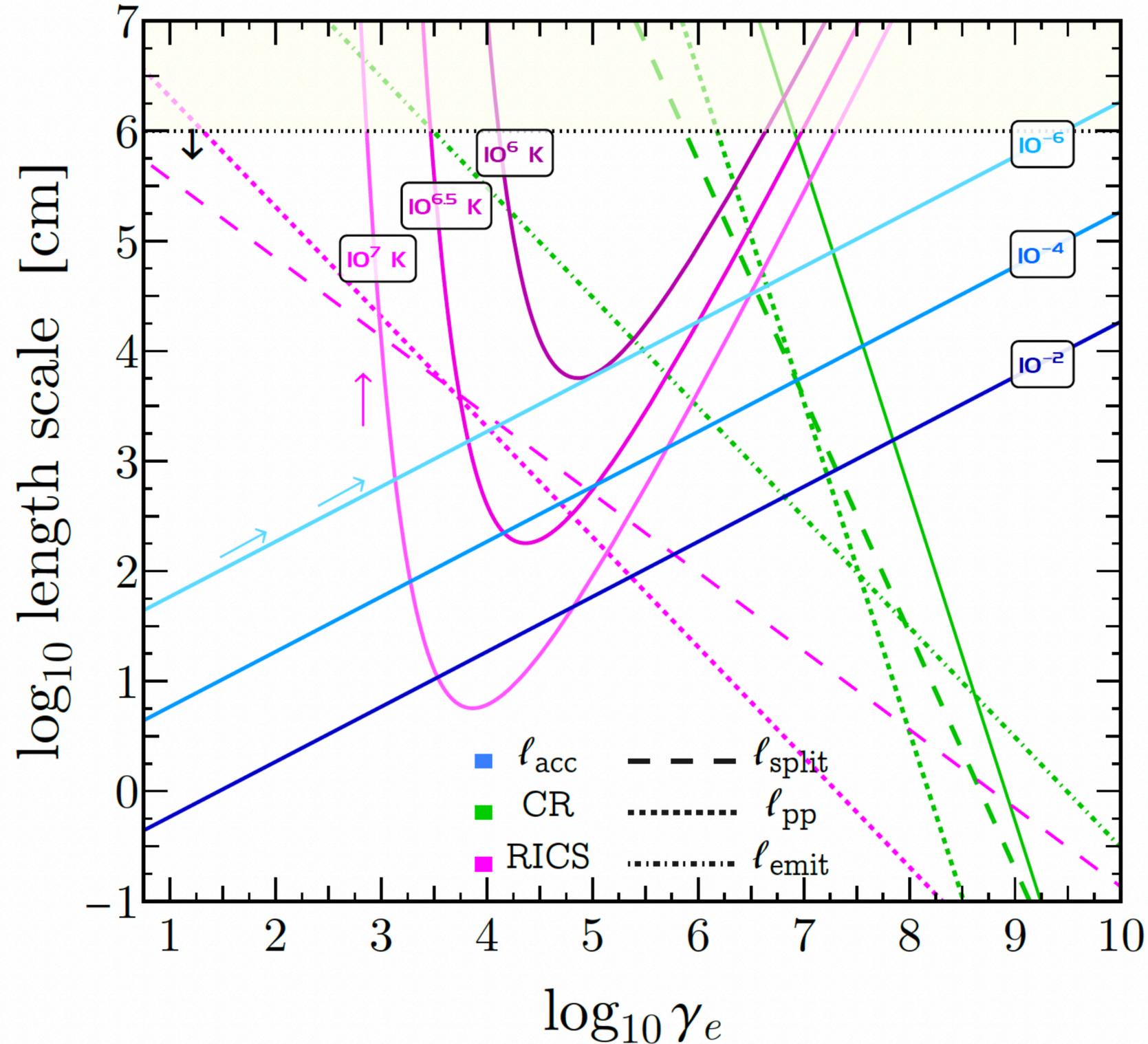
For high temperatures they produce gamma-rays via RICS of thermal photons, which produces pairs



Particles accelerated as they cross gap height

Microphysics

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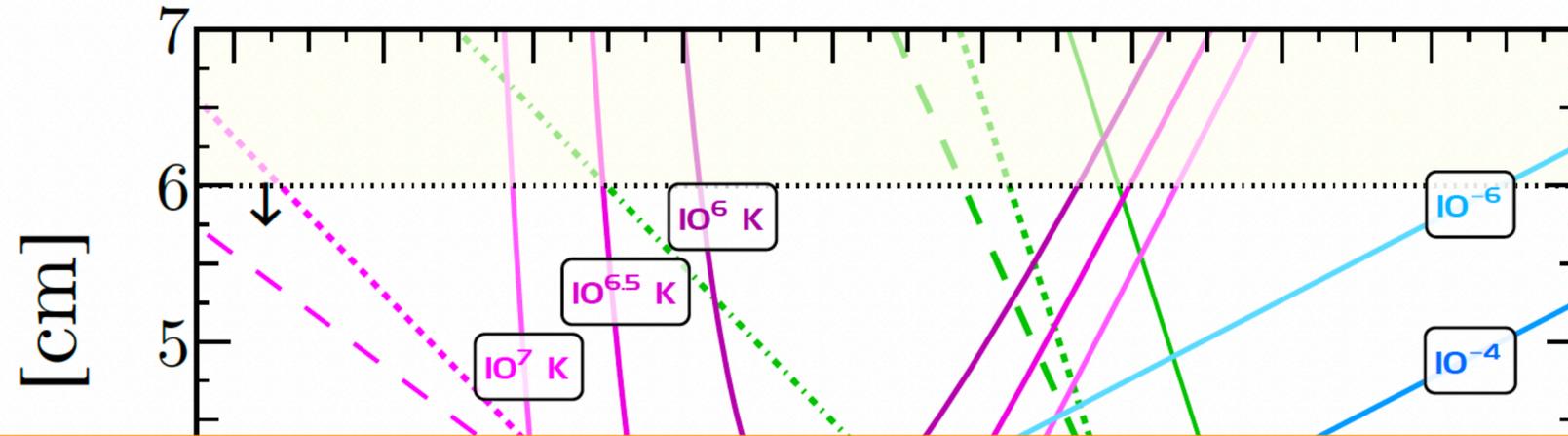


Particles accelerated as they cross gap height

If they 'miss' RICS tail, they produce gamma-ray curvature photons which produces pairs

Microphysics

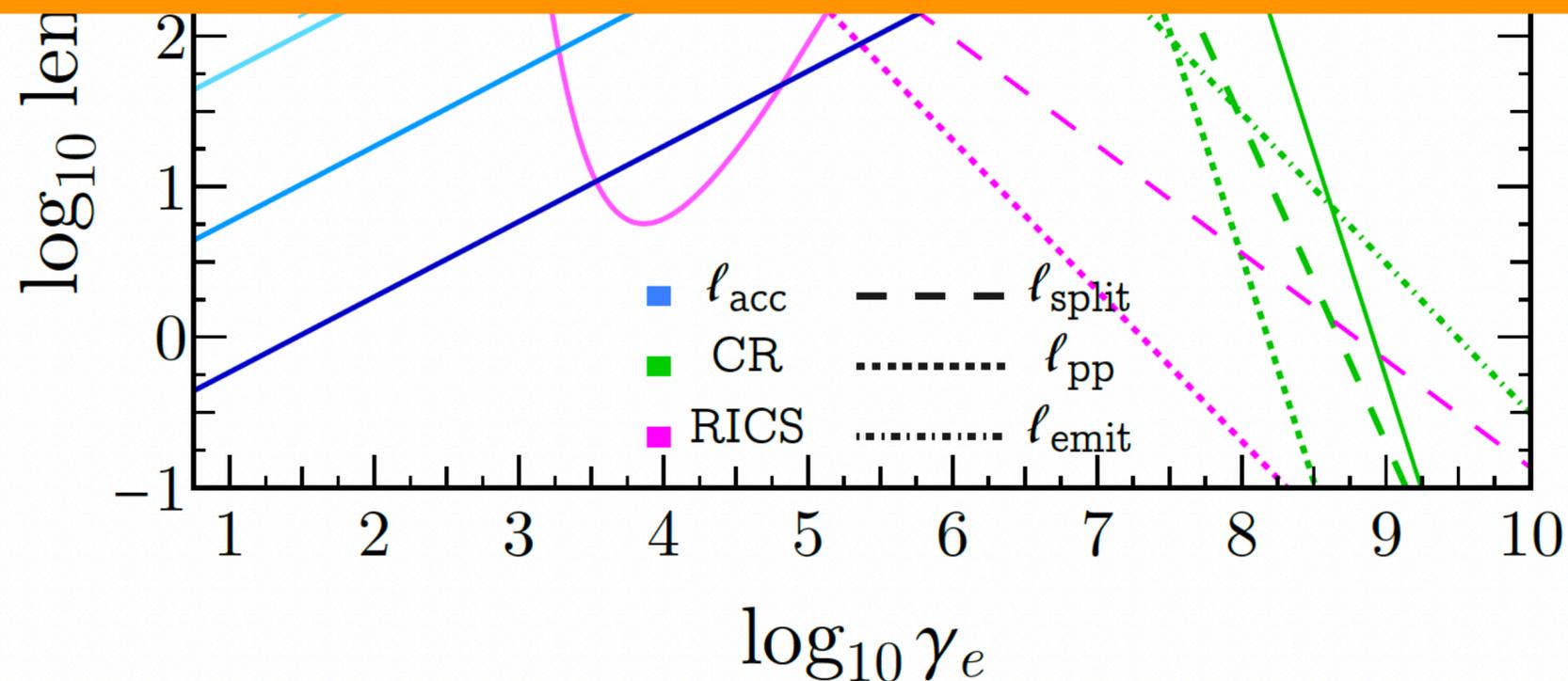
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Particles accelerated as they cross gap height

Take Home Message

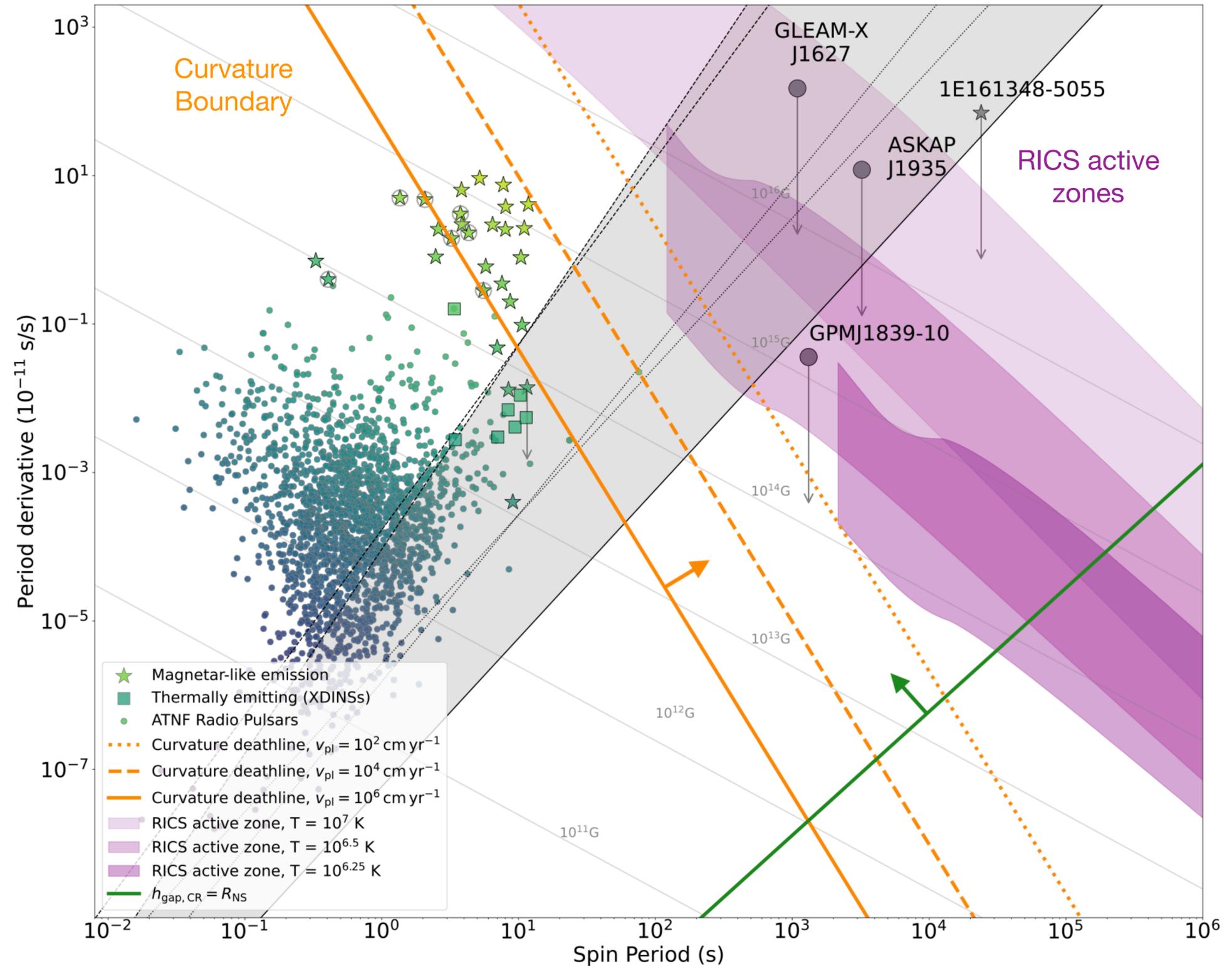
Either curvature photons or upscattered photons will produce pairs and create radio emission



If they 'miss' RICS tail, they produce gamma-ray curvature photons which produces pairs

Model Predictions

- Microphysics defines allowed parameter space



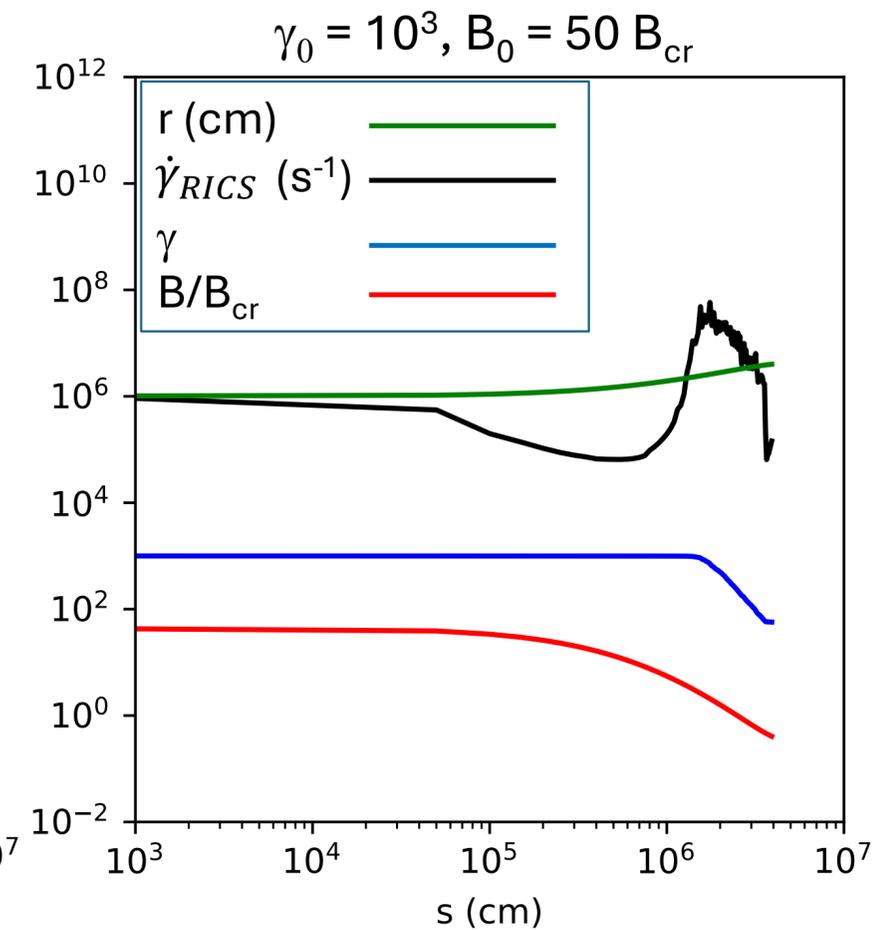
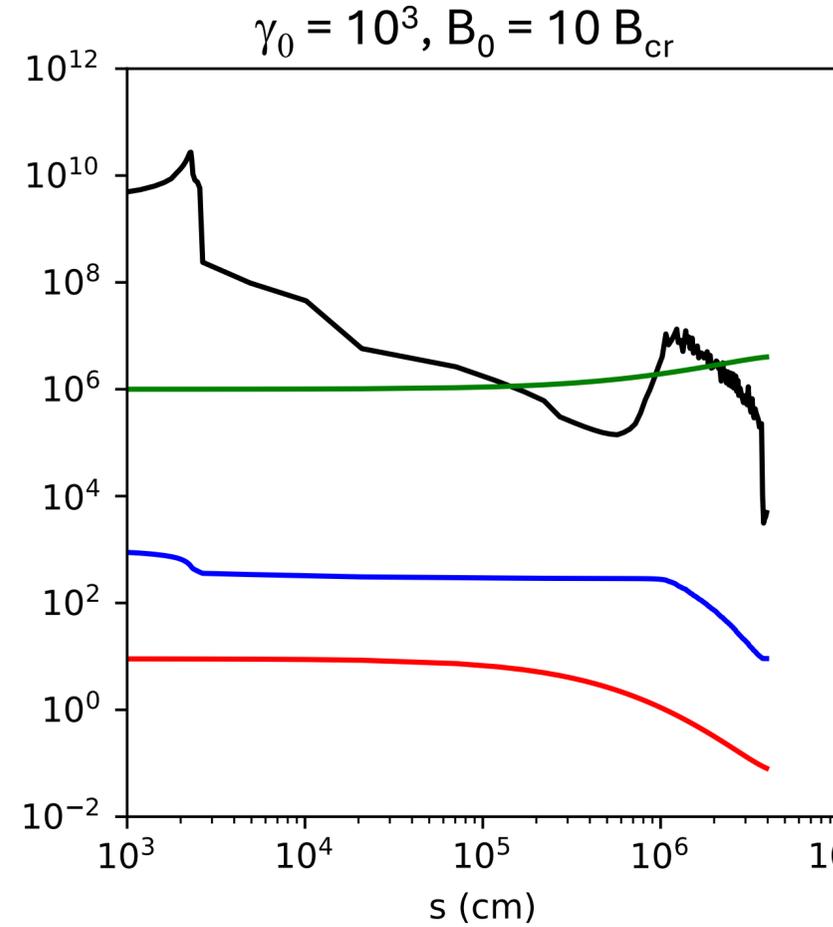
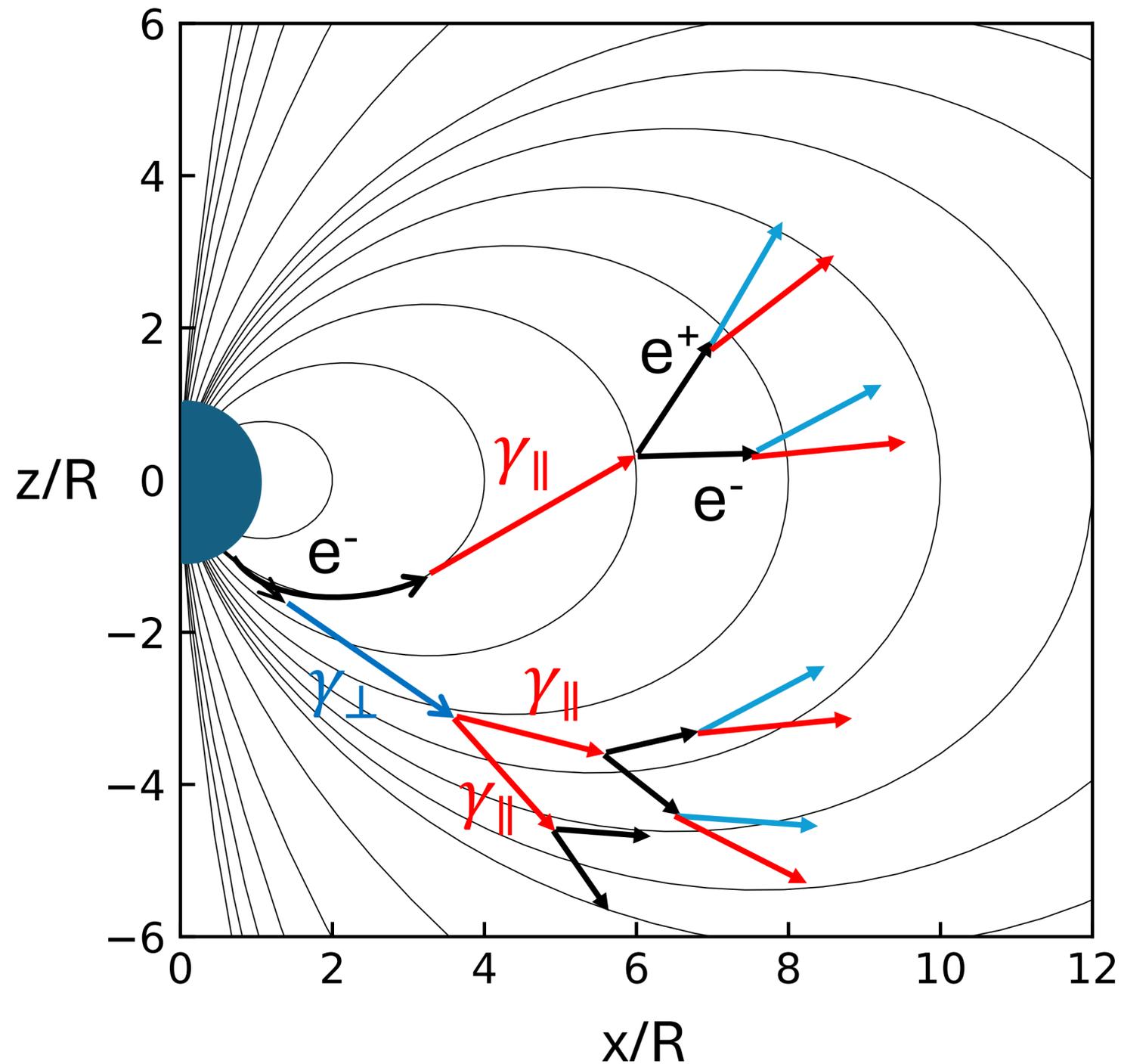
Predictions

- Minimum period depending on v_{pl} or T

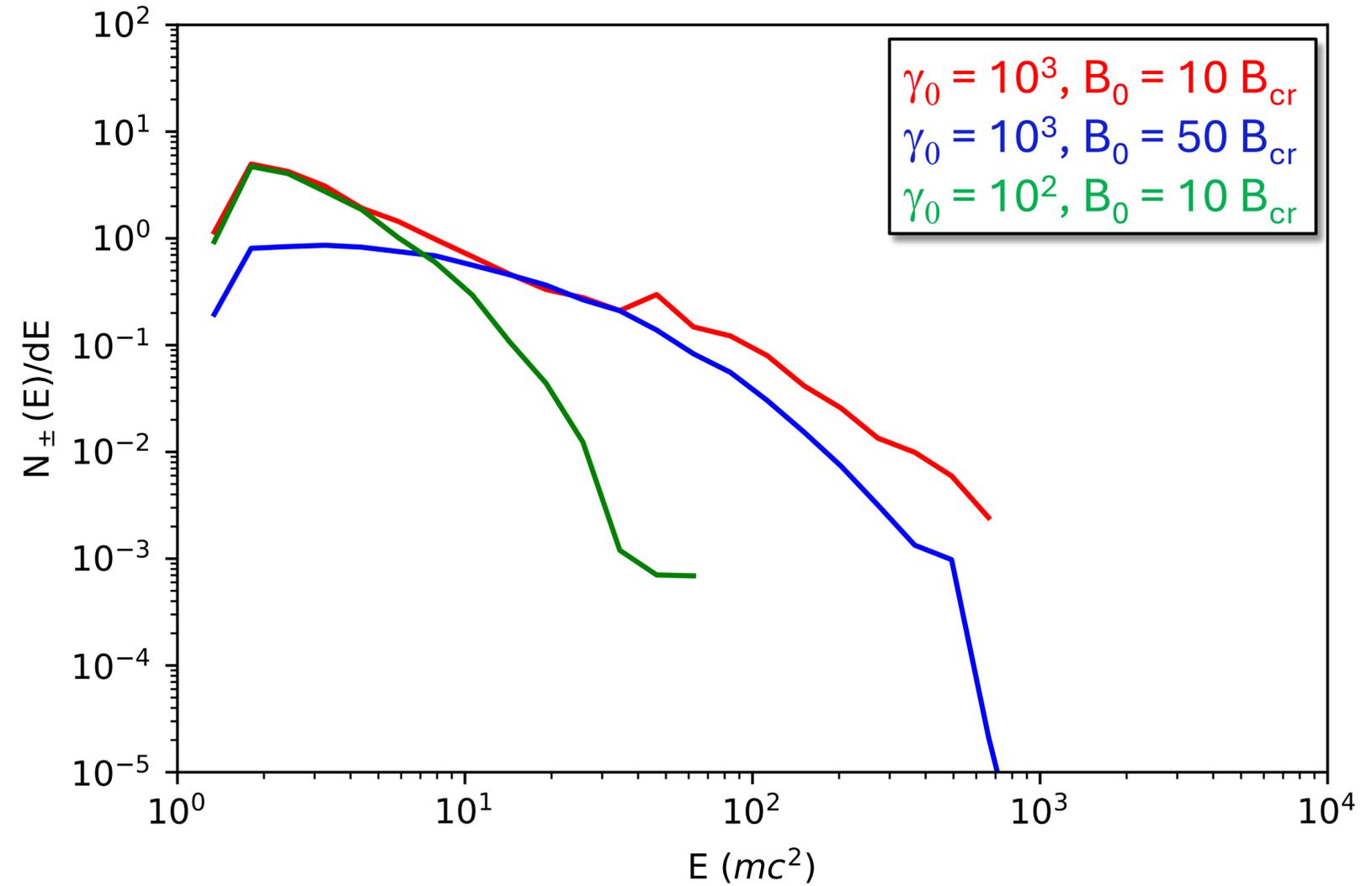
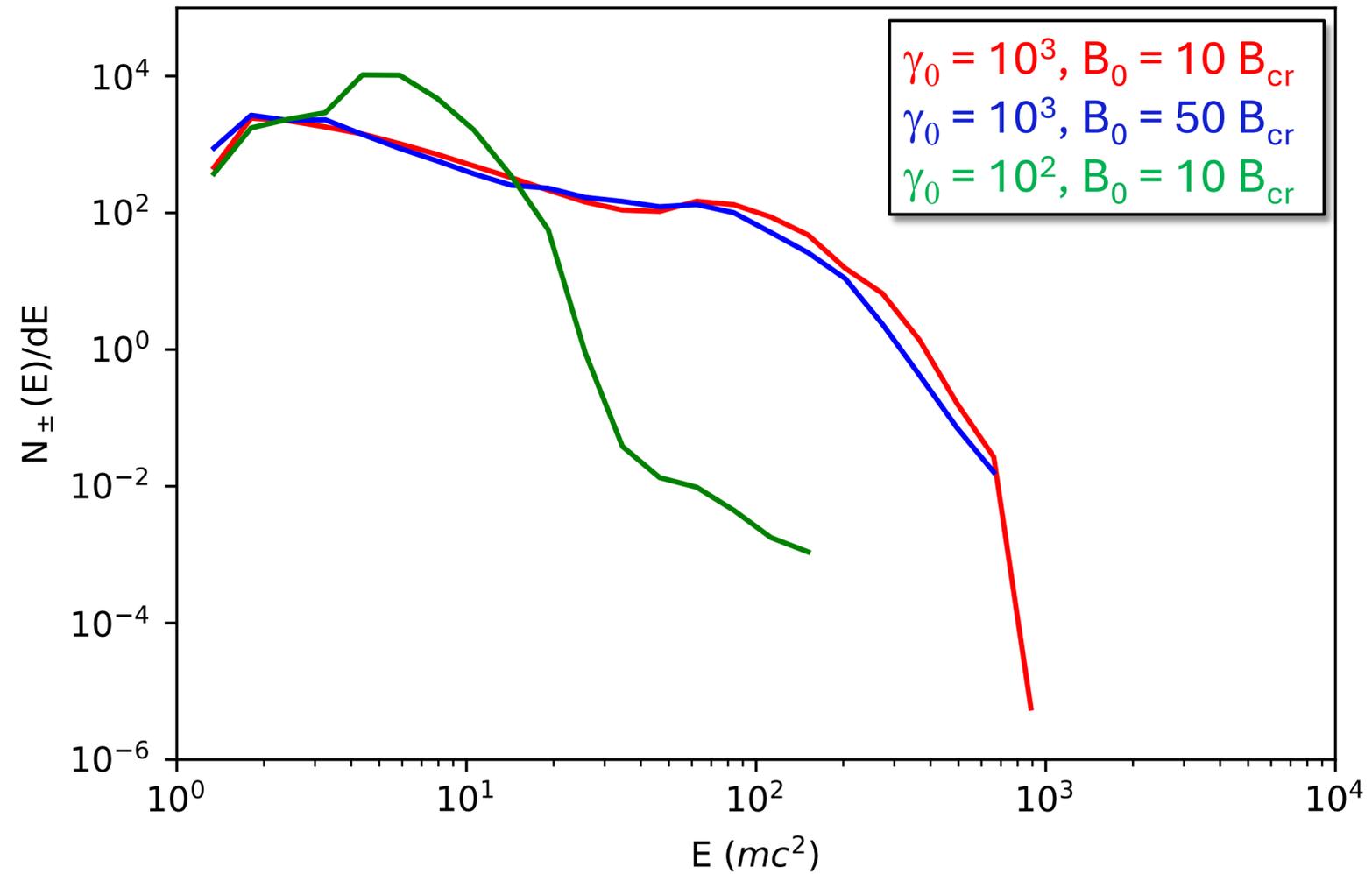
$$P_{\text{RICS}} \gtrsim 120 (T/10^{6.5}\text{K})^{-5} \text{ sec}$$

$$P_{\text{curv}} \gtrsim 150 (v_{\text{pl}}/10^3 \text{ cm yr}^{-1})^{-7/6} \text{ sec}$$

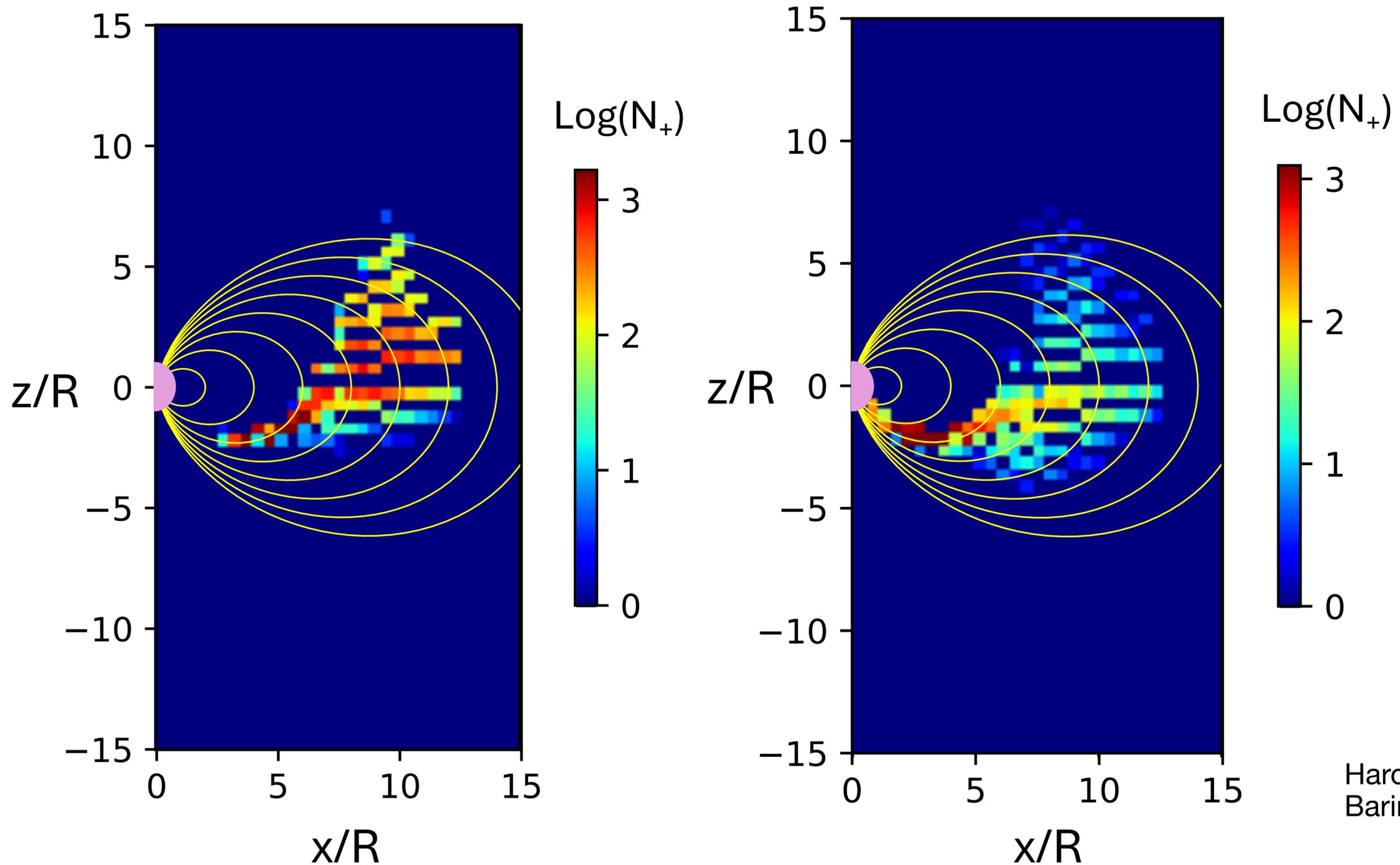
More accurate magnetar RICS cascades with photon splitting



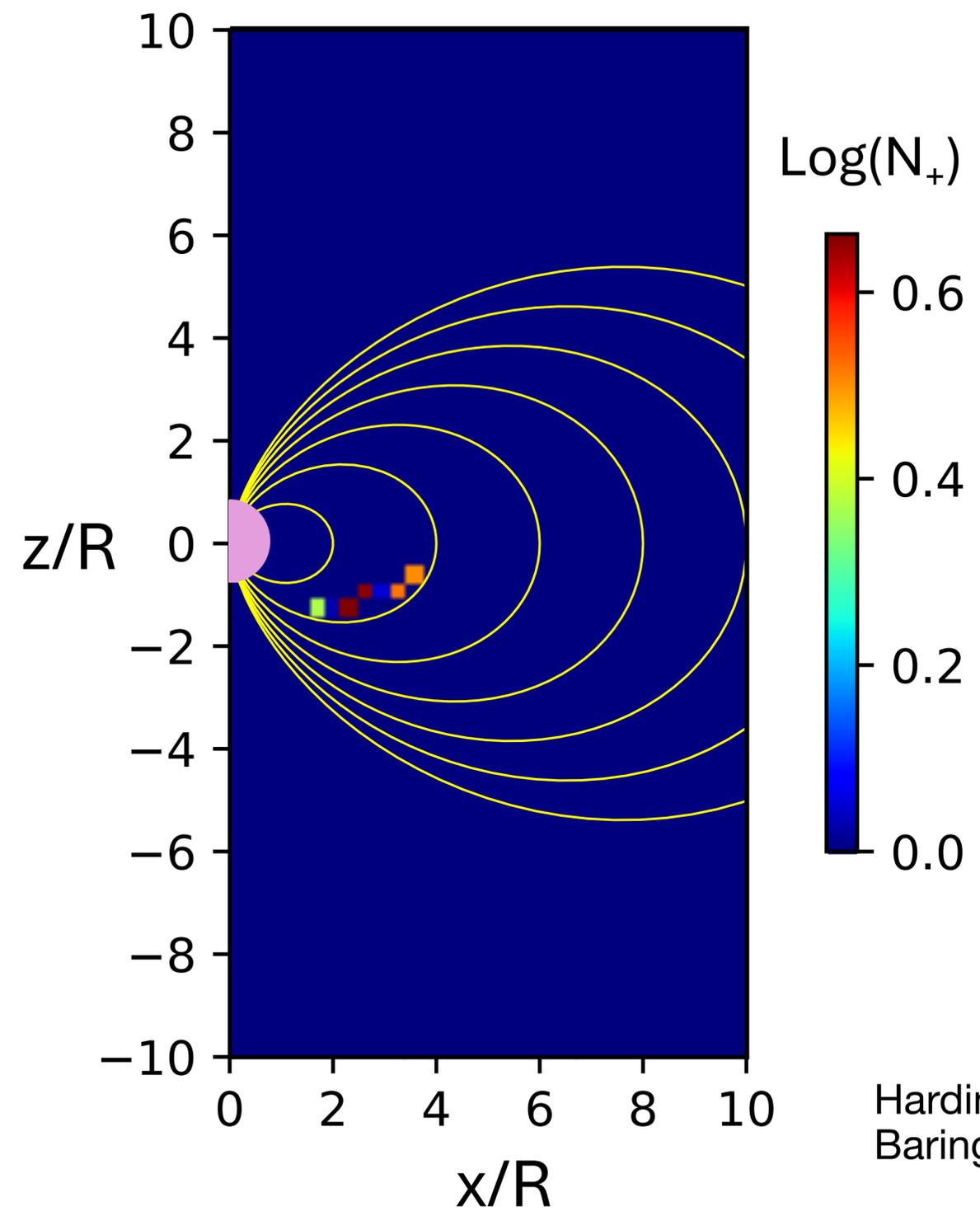
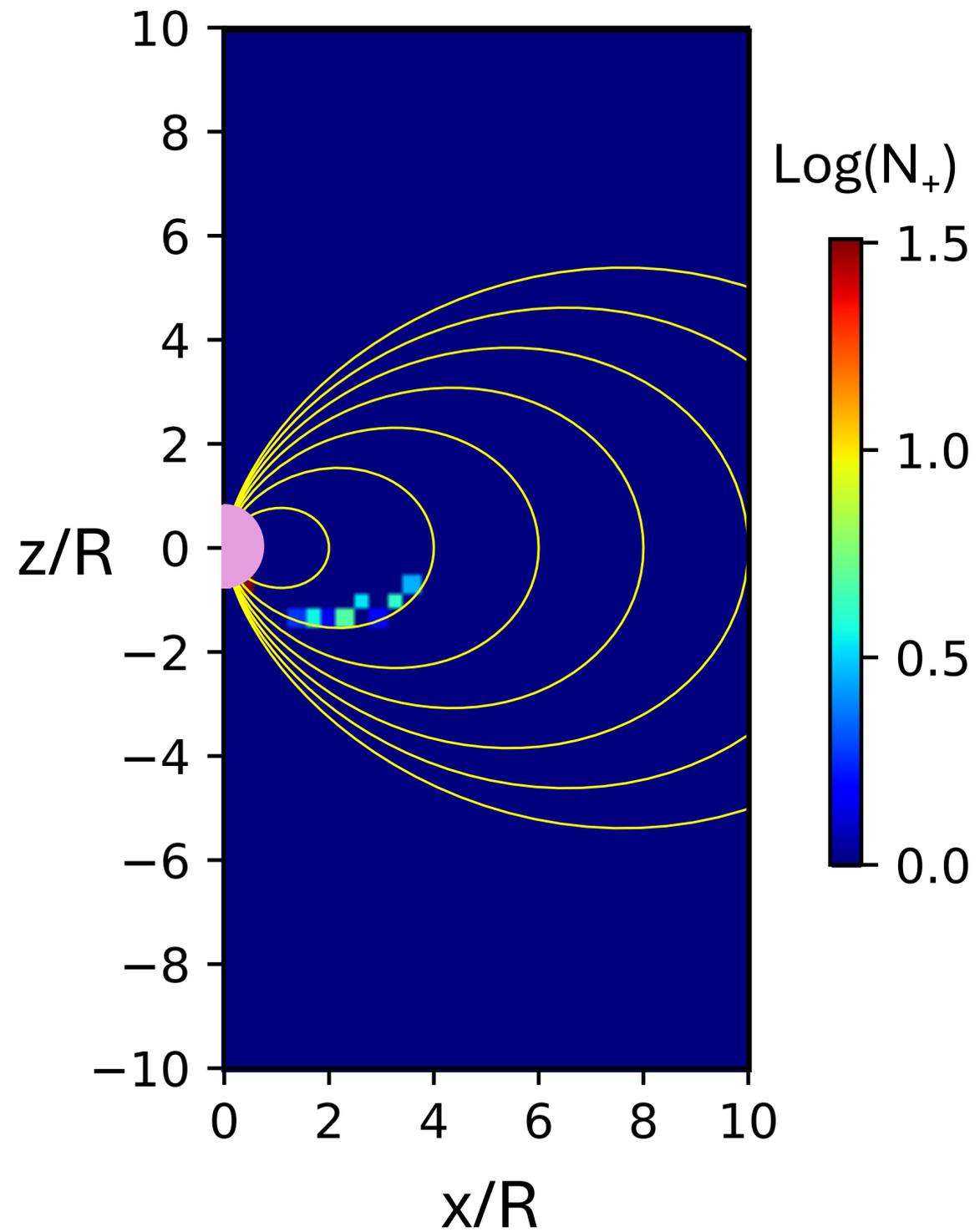
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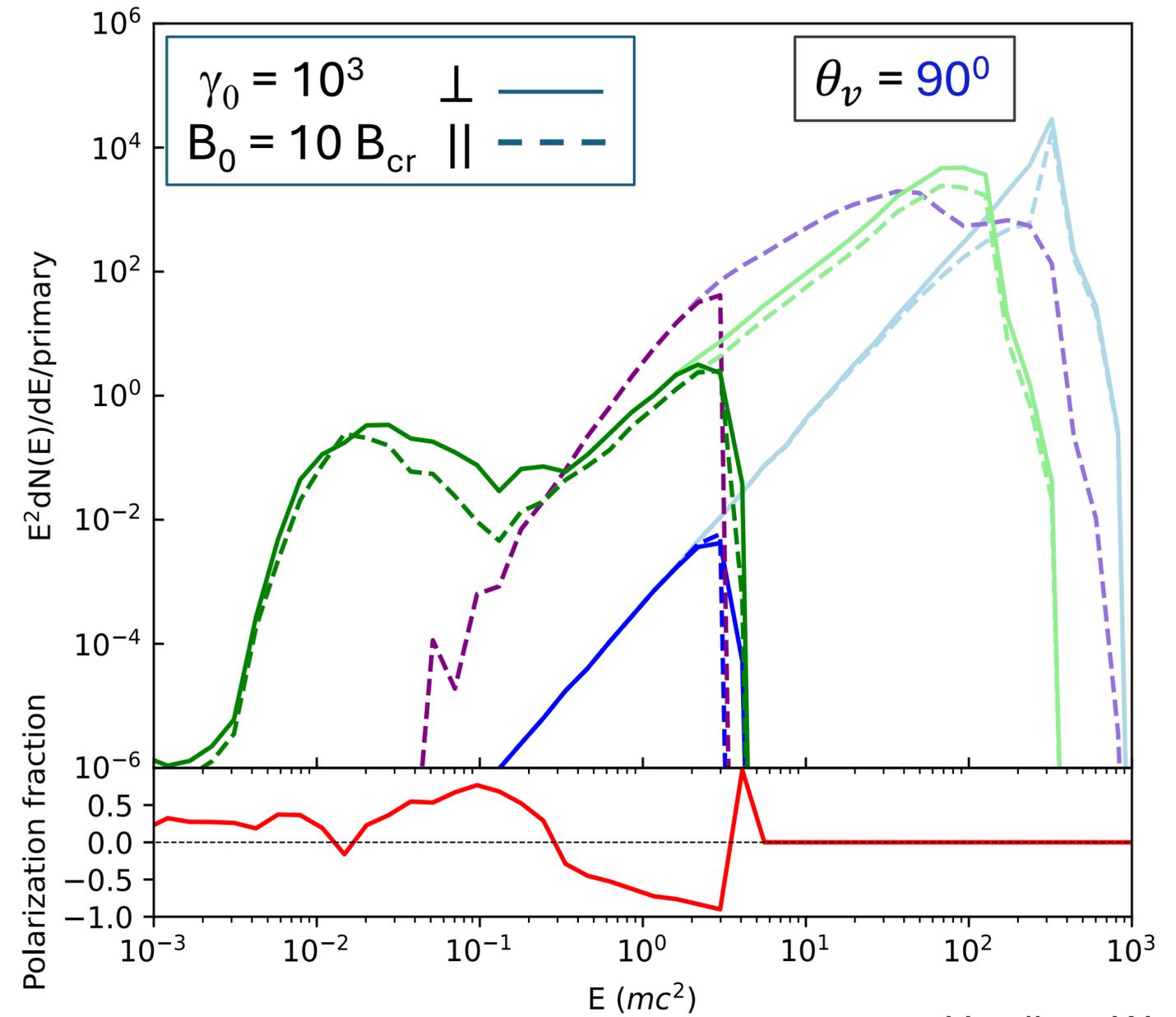
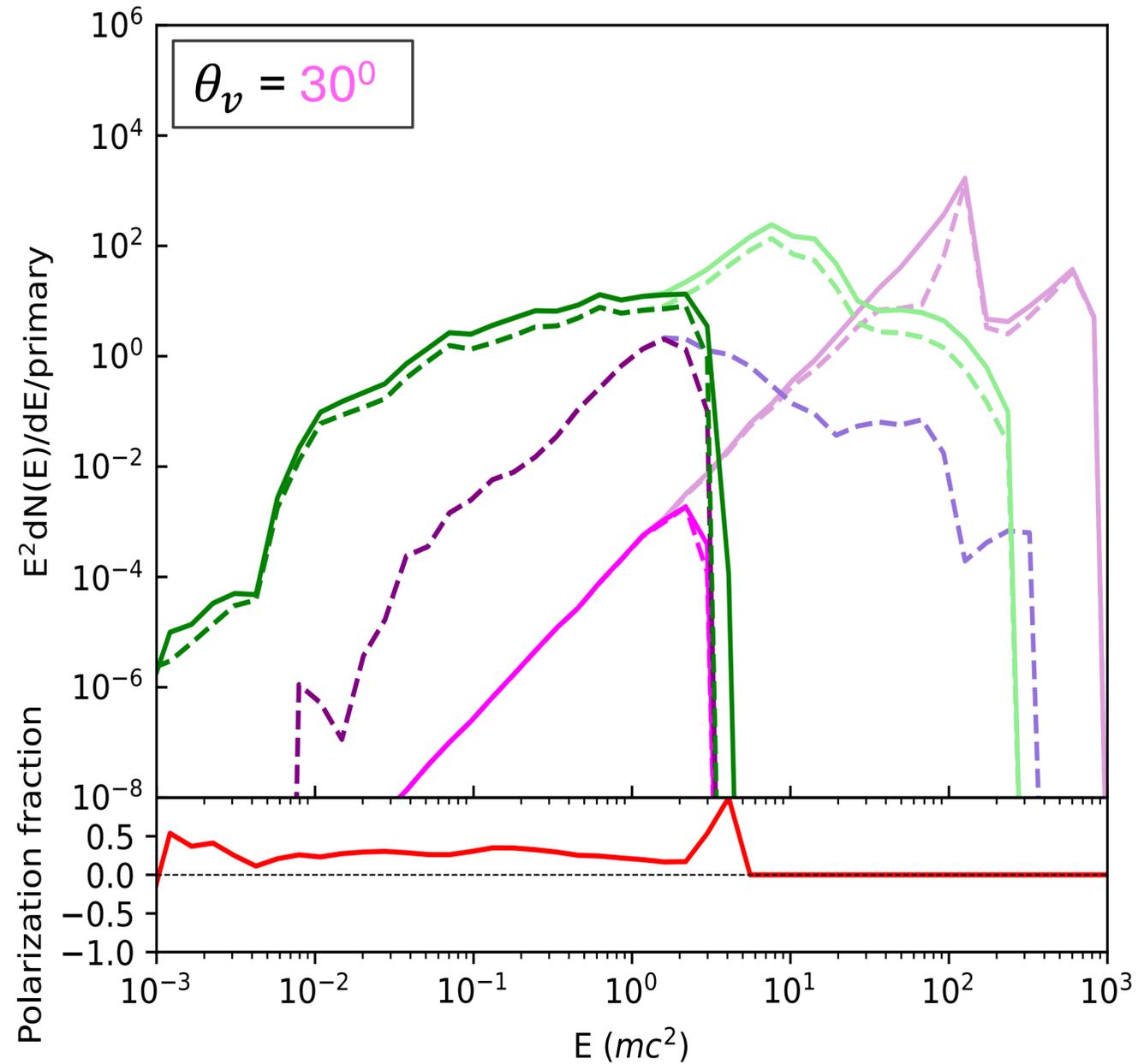
More accurate magnetar RICS cascades with photon splitting



More accurate Magnetar RICS cascades with photon splitting

	No RICS losses			RICS losses		
$r_{\max} = 4$	γ_0			γ_0		
B_0/B_{cr}	10	10^2	10^3	10	10^2	10^3
10	3	4	5	3	6	4
	9×10^{-8}	4.8×10^4	2.7×10^4	9×10^{-8}	15	62
50		5	5		4	5
		7.8×10^3	7.8×10^3		5	20
$r_{\max} = 6$						
10		6	6			
		4×10^4	2.4×10^4			

More accurate Magnetar RICS cascades with photon splitting



Observational Outlook

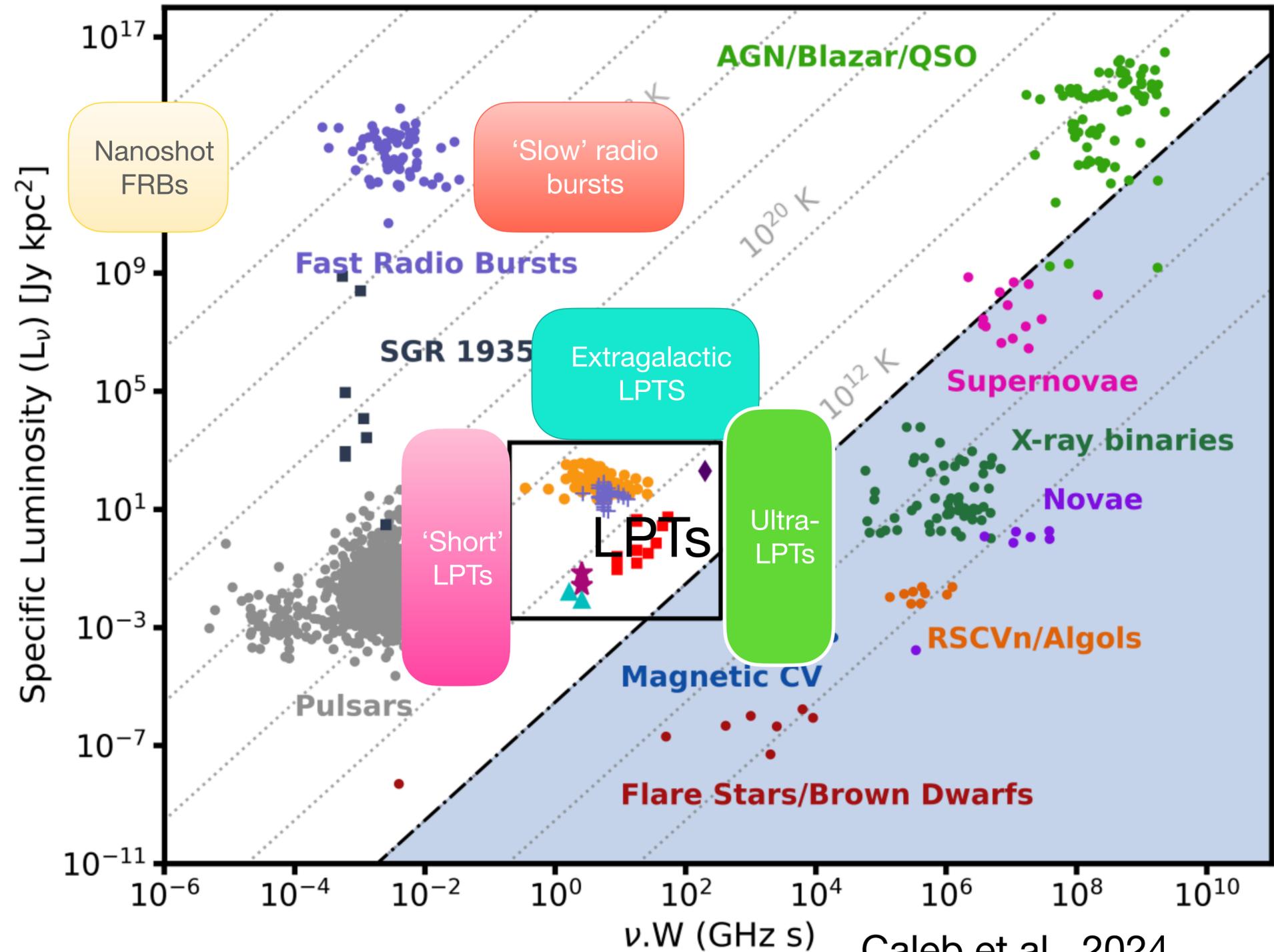
Many more LPTs will be found given current discovery rate

Extra-galactic LPTs will be uncovered

Sensitive X-ray/optical observations of nearby and high latitude sources crucial

New sources will bridge the gap between pulsars and LPTs

see e.g., Wang et al. 2025, 41sec pulsar appears to not rotationally powered)



Mechanisms for Producing ULPMs - enhanced spin-down

Phenomenological evidence for enhanced spin-down

Enhanced spin-down associated with GFs and strong bursting behavior

- SGR 1900+14: $x_p \equiv \frac{\Delta P}{P} \sim 10^{-4}$ after 1998 GF
- SGR 1806-20: Increased \dot{P} since 2004 GF. Up to 2012, P increased by extra 2% compared to pre-GF extrapolation (Younes et al. 15).
- Kinematic age constraints of these magnetars suggest further \dot{P} enhancements in their past (Tendulkar et al. 12)

Simplest phenomenological model⁶ in ~ 100 days (Archibald et al. 13)

If $x_p = \text{const}$ then $P_f = P_0 \exp(N_p x_p) \rightarrow P_f \gg P_0$ for $N_p > x_p^{-1}$

- With $E_{GF} \sim 4 \times 10^{44} \text{ erg}$ and $x_p \sim 10^{-4}$, a significant increase of P requires a magnetic energy reservoir of $> 4 \times 10^{48} \text{ erg}$ or internal field $B_{int} > 5 \times 10^{15} \text{ G}$
- Compare to SGR 1900+14: $B_{dip} = 7 \times 10^{14} \text{ G}$ and recall that $B_{int} \sim 10 B_{dip}$ inferred from X-rays

Physical mechanisms for enhanced spin-down

Charged particle winds

- Mass-loaded charged wind with $L_{pw} > L_{dip}$ opens up B lines beyond

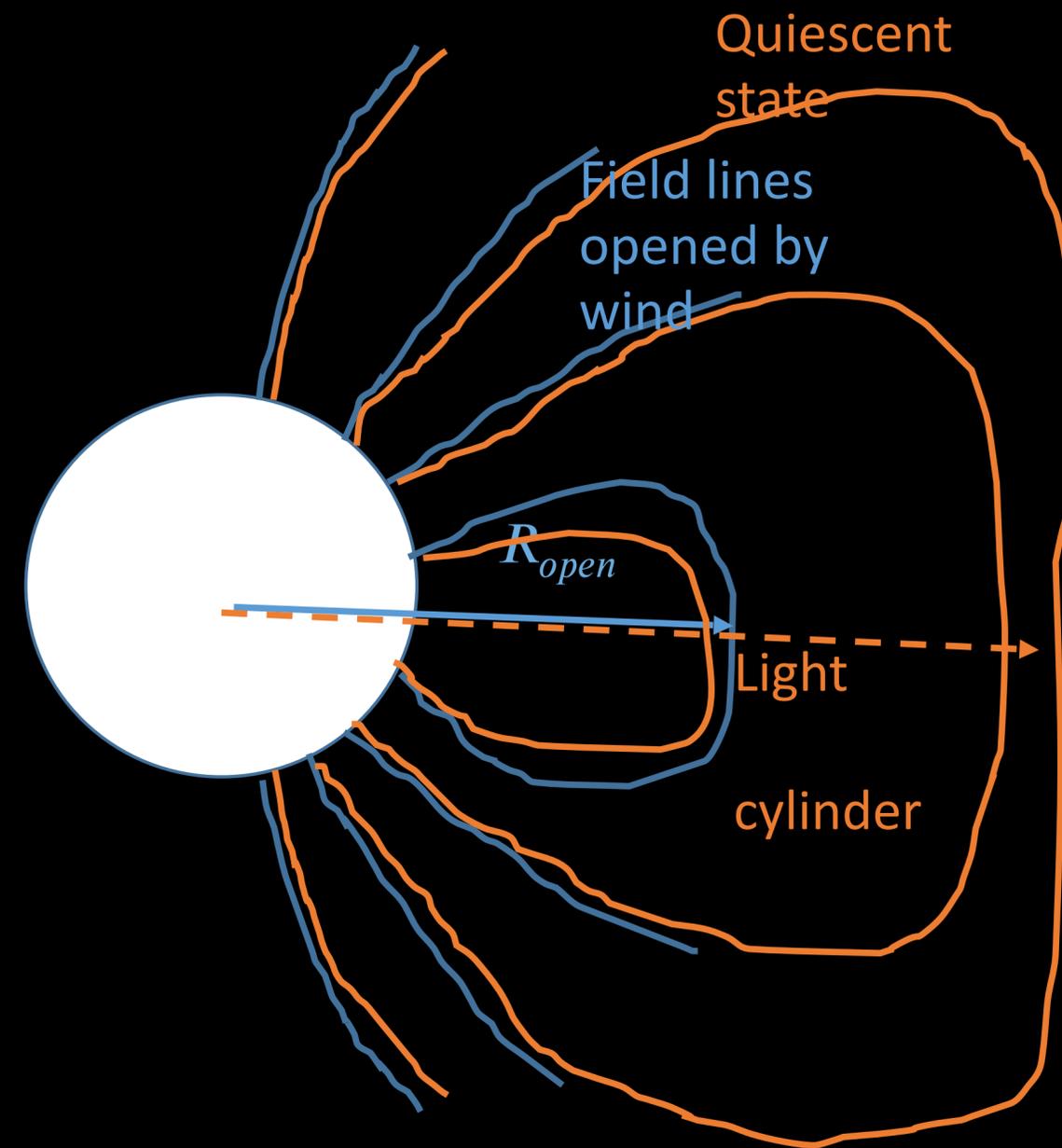
$$R_{open} \sim R_{NS} \left(\frac{B_{dip}^2 R_{NS}^2 c}{L_{pw}} \right)^{1/4} \quad (\text{Thompson \& Blaes 98, Harding et al. 00})$$

- Spindown scales as open flux squared \rightarrow Enhanced spindown $\dot{P} \propto P$

$$\tau = \frac{IcR_{open}^2}{B_{dip}^2 R_{NS}^6} = \frac{Ic^{3/2}}{B_{dip} R_{NS}^3 L_{pw}^{1/2}} = 5 \times 10^7 B_{dip,15}^{-1} L_{pw,40}^{-1/2} \text{ s}$$

- $P_f = P_0 \exp\left(-\frac{t}{\tau}\right)$ with

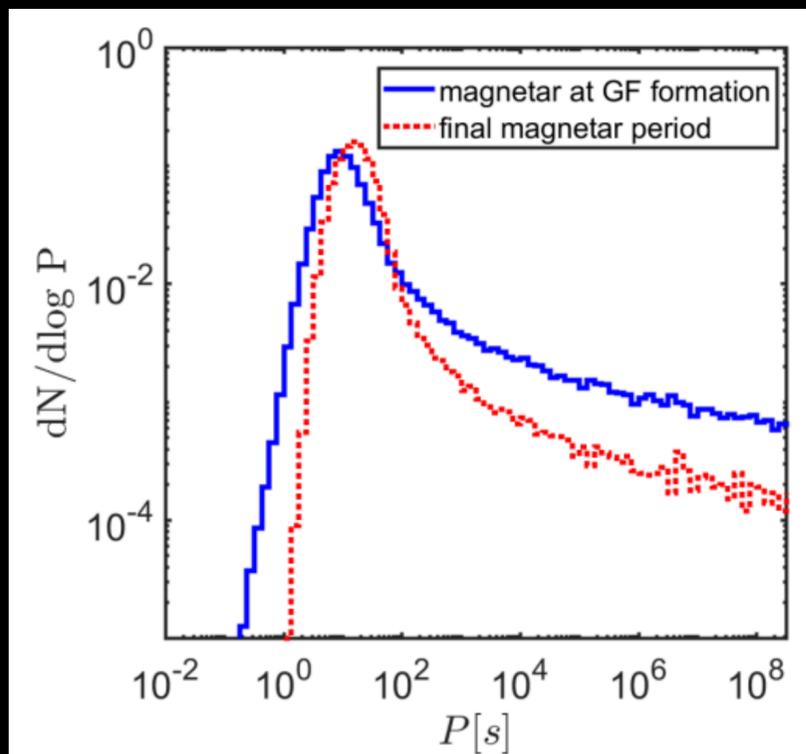
$$P_f = P_0 \exp\left[\frac{E_B \Delta t_{pw}}{E_f \tau}\right] = P_0 \exp\left[0.7 \frac{B_{int,16}^2 B_{dip,15} E_{pw,42}^{1/2} \Delta t_{pw,2}^{1/2}}{E_{f,44}}\right]$$



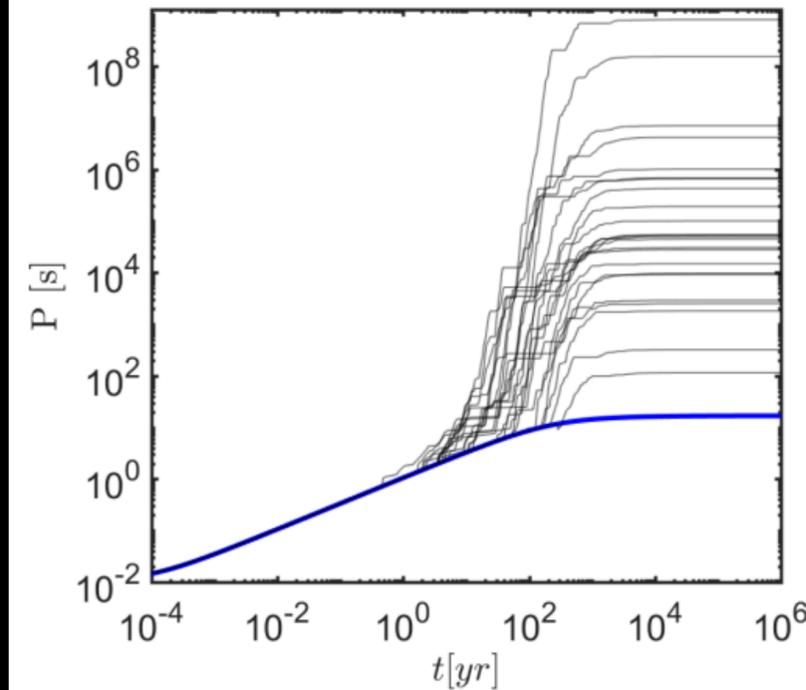
Physical mechanisms for enhanced spin-down

Charged particle winds

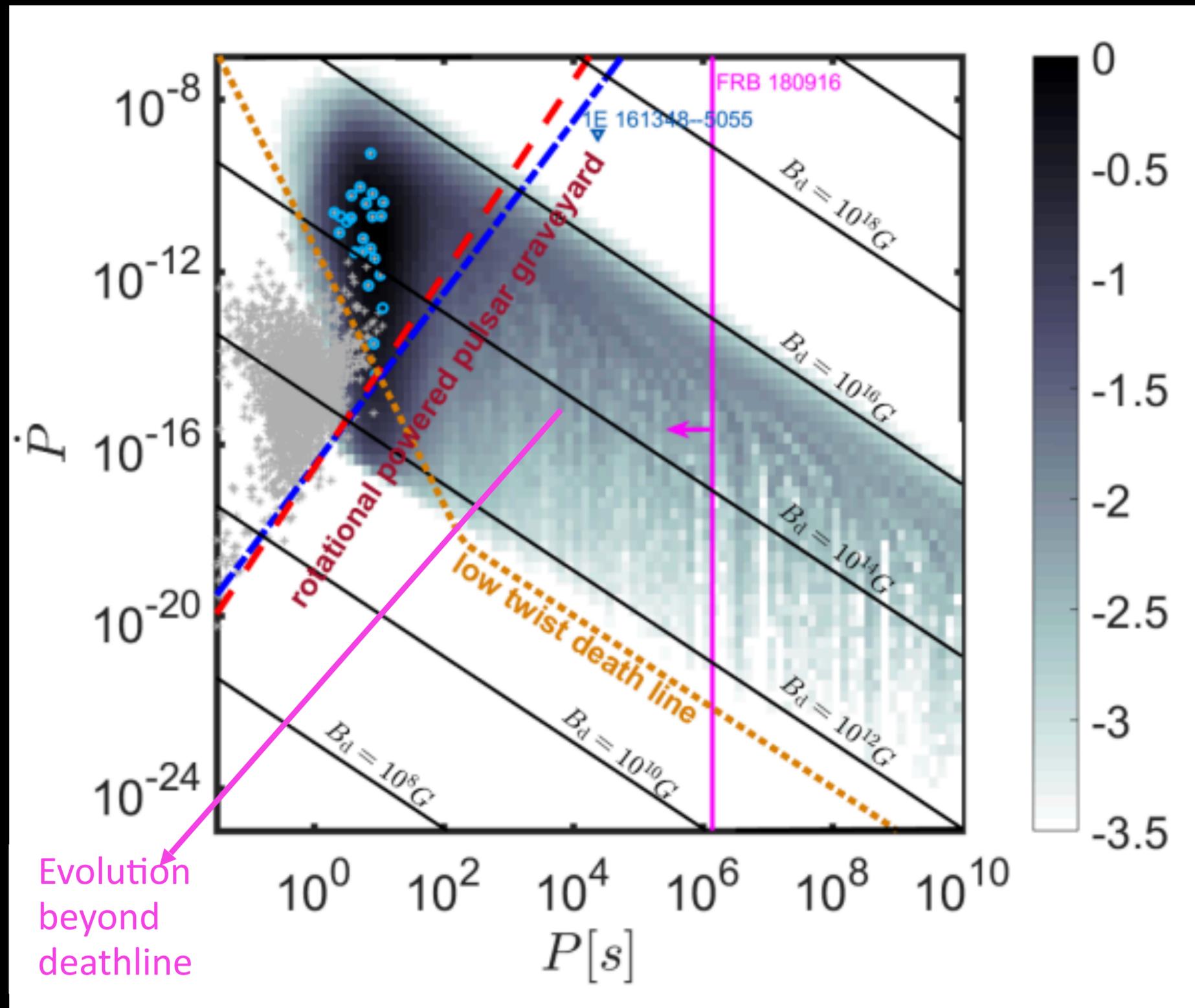
- Monte Carlo proof of concept:



Flat P distribution at large P



Example P evolutions

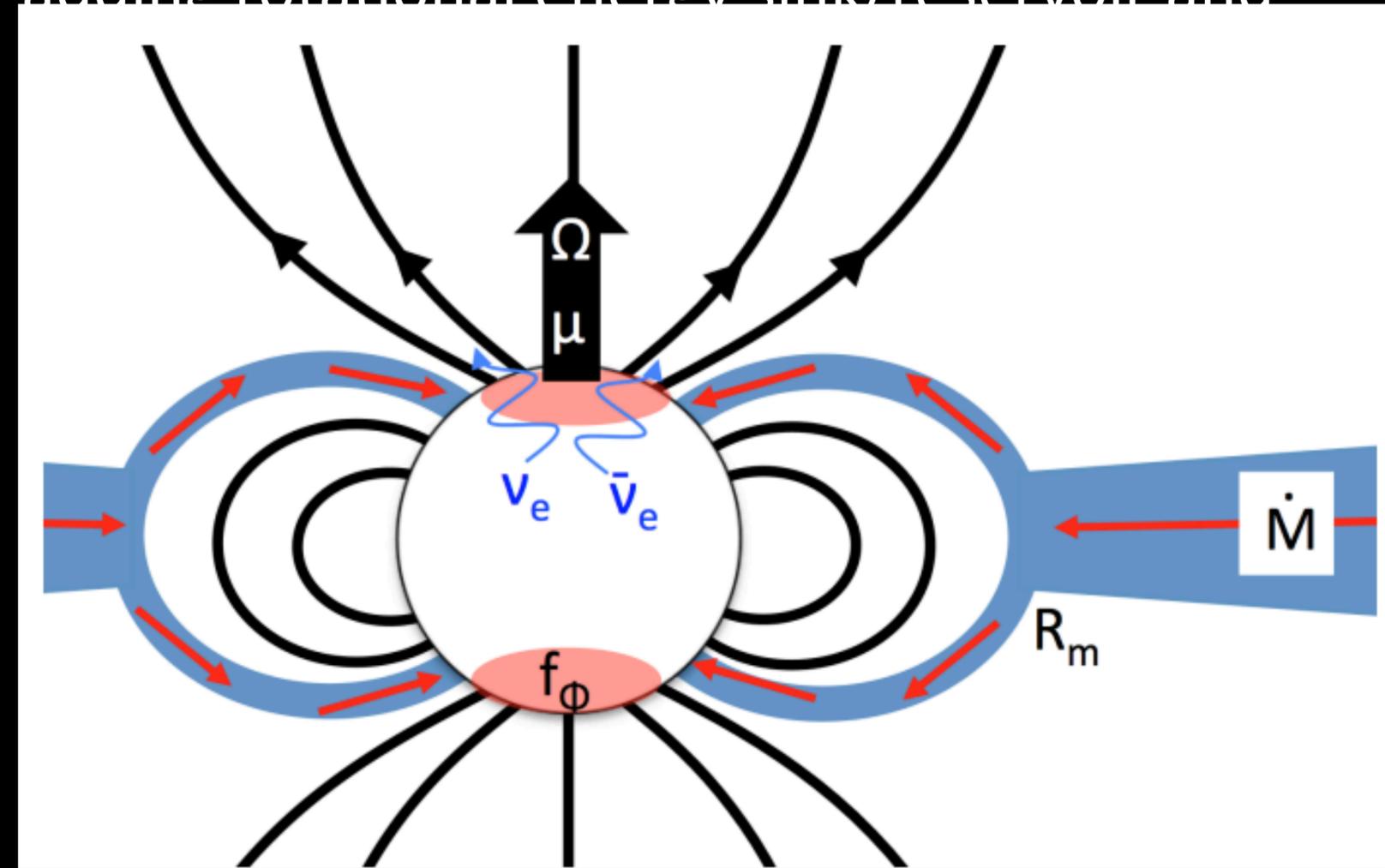
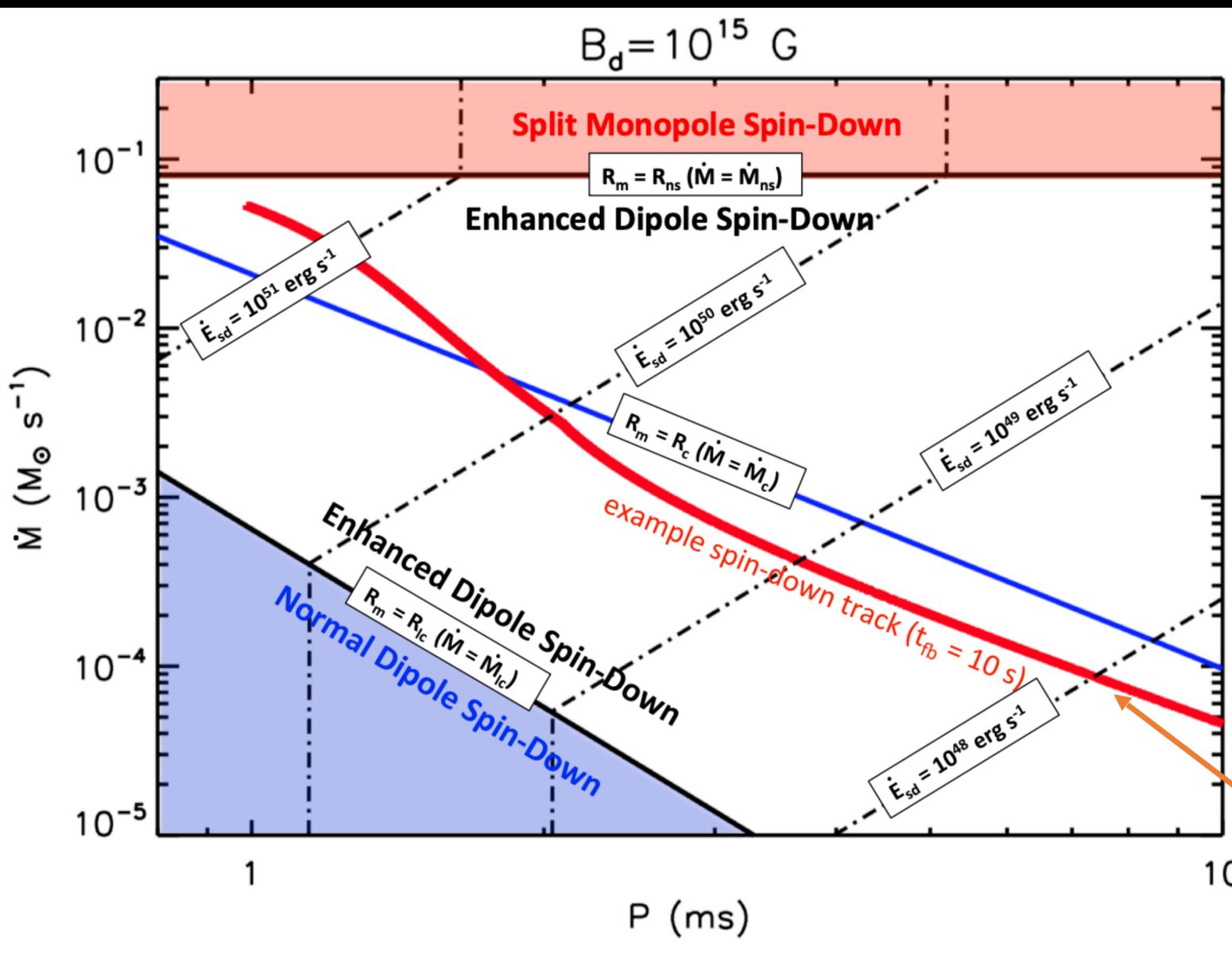


Evolution beyond deathline

Physical mechanisms for enhanced spin-down

Fallback accretion

- RCW103 – sub-energetic SN remnant: consistent with more fallback (Braun et al. 2019)
- Fallback accretion alters magnetar evolution by adding rotational energy sink/reservoir and enhancing spindown by opening up field lines



Rough equilibrium between co-rotation and Alfvén radius

Physical mechanisms for enhanced spin-down

Fallback accretion

- P exponentially increases until $R_m \sim R_c$ and evolves as $t^{3\zeta/7}$ afterwards, where $\dot{M} \propto t^{-\zeta}$
- Large ζ expected for high \dot{M} RIAFs
- ζ cannot be too large to avoid early disk disruption
- Maximum period set by time it takes magnetic field to decay (relative to initial fallback time)

